

1 **The instability theory of drumlin formation and its explanation of their**
2 **varied composition and internal structure**

3
4 Chris R. Stokes¹, Andrew C. Fowler^{2, 3}, Chris D. Clark⁴, Richard C.A. Hindmarsh⁵, Matteo
5 Spagnolo⁶

6 ¹*Department of Geography, Durham University, Durham, DH1 3LE, UK (c.r.stokes@durham.ac.uk)*

7 ²*MACSI, Department of Mathematics and Statistics, University of Limerick, Limerick, Republic of Ireland*

8 ³*Oxford Centre for Industrial and Applied Mathematics (OCIAM), Mathematical Institute, 24-29 St. Giles',*
9 *Oxford, OX1 3LB, UK*

10 ⁴*Department of Geography, University of Sheffield, Sheffield, S10 2TN, UK*

11 ⁵*British Antarctic Survey, Madingley Road, Cambridge, CB3 0ET, UK*

12 ⁶*School of Geosciences, University of Aberdeen, UK*

13
14 *Corresponding author: Tel. +44 (0)191 334 1955; Fax. +44 (0)191 334 1801

15 E-mail address: c.r.stokes@durham.ac.uk (C R Stokes)

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18
19 **Abstract:**

20 Despite their importance in understanding glaciological processes and constraining large-
21 scale flow patterns in palaeo-glaciology, there is little consensus as to how drumlins are
22 formed. Attempts to solve the 'drumlin problem' often fail to address how they are created
23 from an initially flat surface in the absence of obvious cores or obstacles. This is a key

24 strength of the instability theory, which has been described in a suite of physically-based
25 mathematical models and proposes that the coupled flow of ice and till causes spontaneous
26 formation of relief in the till surface. Encouragingly, model predictions of bedform height
27 and length are consistent with observations and, furthermore, the theory has been applied to a
28 range of subglacial bedforms and not just drumlins. However, it has yet to confront the
29 myriad observations relating to the composition and internal structure of drumlins and this
30 could be seen as a major deficiency. This paper is a first attempt to assess whether the
31 instability theory is compatible with the incredible diversity of sediments and structures
32 found within drumlins. We summarise the underlying principles of the theory and then
33 describe and attempt to explain the main types of drumlin composition (e.g. bedrock, till,
34 glacio-fluvial sediments, and combinations thereof). Contrary to a view which suggests that
35 the presence of some sedimentary sequences (e.g. horizontally stratified cores) is inconsistent
36 with the theory, we suggest that one would actually expect a diverse range of constituents
37 depending on the inheritance of sediments that pre-date drumlin formation, the duration and
38 variability of ice flow, and the balance between erosion and deposition (till continuity) at the
39 ice-bed interface. We conclude that the instability theory is compatible with (and potentially
40 strengthened by) what is known about drumlin composition and, as such, offers the most
41 complete and promising solution to the drumlin problem to date.

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43

44 **1. Introduction**

45 Drumlins are one of the most widely studied landforms on Earth, with >1300 contributions
46 (papers, abstracts and theses) in the literature and >400 scientific papers since 1980 (Clark *et*
47 *al.*, 2009). Their importance stems from their relevance to both glaciology and palaeo-

48 glaciology. In glaciology, they are important because they are formed at the ice-bed interface
49 and may exert a modulating effect on ice flow (Schoof, 2002a). However, our understanding
50 of the subglacial processes that occur at this interface is incomplete, largely because of the
51 inaccessibility of this environment under extant glaciers and ice sheets, which limits our
52 observations to geophysical surveys or borehole sampling that cover relatively small areas
53 (e.g. King *et al.*, 2007; 2009; Tulaczyk *et al.* 2000). Thus, investigation of drumlins on
54 former ice sheet beds has the potential to uncover important new insights regarding the
55 mechanisms and feedbacks that act to sustain and/or inhibit ice flow and, importantly,
56 formulate and test models of subglacial processes at the ice-bed interface (e.g. Schoof, 2007a,
57 b; Fowler, 2000; 2009a; Hindmarsh, 1998a; 1998b; 1999). Ultimately, the success of such
58 models to account for drumlin formation will improve our ability to predict the rate at which
59 ice and sediment is transported from continents to the oceans, with important implications for
60 future ice sheet stability (e.g. Schoof, 2002; 2004).

61 In palaeo-glaciology, drumlins also record key information relating to ice sheet flow history,
62 e.g. ice flow direction and changes through time (cf. Boulton & Clark, 1990; Clark, 1993;
63 Kleman and Borgström., 1996; Kleman *et al.*, 1997, 2006), and even ice velocity (cf. Hart,
64 1999; Stokes and Clark, 2002). Thus, they are a vital ingredient for glacial inversion
65 techniques that use the geological record of former ice sheet beds to reconstruct their time-
66 dependent behaviour (see Kleman and Borgström, 1996; Kleman *et al.*, 2006). It could be
67 argued, however, that their use is yet to fulfil its true potential. If, for example, we knew the
68 specific conditions under which drumlins of different shapes and sizes developed (e.g.
69 specific ranges of ice thickness, velocity, effective pressures, etc.), then their importance to
70 palaeoglaciology would be considerably magnified.

71 With the above considerations in mind, the quest for a physically-based model of drumlin
72 formation takes on huge importance and yet, despite this, their origin remains enigmatic and

73 controversial. Numerous hypotheses of drumlin formation have been espoused and include
74 accretion around obstacles (e.g. Fairchild, 1929), dilatant till behaviour (e.g. Smalley and
75 Unwin, 1968), catastrophic meltwater floods (e.g. Shaw, 1983), deformation of till around
76 more competent cores (e.g. Boulton, 1987), lee-side cavity infillings (e.g. Dardis, 1985) and
77 an instability at the ice-till interface (e.g. Hindmarsh, 1998a). As noted by Clark (2010),
78 however, most ideas/hypotheses of drumlin formation fail to address how bumps (drumlins)
79 are created from a flat surface in the absence of obvious obstacles or cores. This appears to be
80 a critical aspect of the ‘drumlin problem’ because although some drumlins possess an
81 obvious core (e.g. of bedrock or ‘stiffer’ material), there are numerous reports in the literature
82 of those that do not (see reviews in Patterson and Hooke, 1995; Stokes *et al.*, 2011).
83 Furthermore, most of ideas regarding drumlin genesis are restricted to qualitative
84 descriptions/explanations and very few have progressed to physically-based mathematical
85 models that are capable of making predictions that can be tested against observations.

86 One theory of drumlin formation that does address relief amplification from an initially
87 featureless surface is the instability theory and, significantly, the last decade or so has seen it
88 described in numerical models of ice flowing over a layer of deforming sediment (e.g.
89 Hindmarsh, 1998a, b; 1999; Schoof, 2007a, b; Fowler, 2000; 2009a; 2010a). It proposes that
90 the coupled flow of ice and till causes the spontaneous formation of relief in the till surface,
91 whereby local highs at the bed will accumulate till by deposition, and lows will be
92 preferentially eroded. This leads to the creation of pattern and structure in the bed that is
93 manifest in a wide range of features termed subglacial bedforms. Indeed, a further appeal of
94 this theory is that it has the potential to provide a unifying explanation for the production of a
95 continuum of subglacial bedforms (cf. Aario, 1977; Rose, 1987) and not only drumlins,
96 having been applied to ribbed moraine (Dunlop *et al.*, 2008; Chapwanya *et al.*, 2011) and
97 recently adapted to address the formation of mega-scale glacial lineations (Fowler, 2010b).

98 Although models of the instability theory are yet to generate fully three-dimensional drumlins
99 (see section 2.2), predictions of bedform height and length in two-dimensional treatments are
100 consistent with observations (e.g. Fowler, 2000; 2009), which is encouraging (see discussion
101 in Clark, 2010). Perhaps more serious, however, is that the theory has yet to confront the
102 multitude of observations relating to the composition and internal structure of drumlins. This
103 was recently highlighted by Hiemstra *et al.* (2011) who noted that “theoretical studies of flow
104 instability have yet to provide a solution for the sedimentological and structural-architectural
105 variability in drumlins as recorded in the field”. For some, this might be viewed as a
106 deficiency: Hart (2005: p. 194), for example, notes that “any model of drumlin formation
107 needs to be related to the sedimentology and structural geology of the drumlins themselves”
108 and Schoof (2007a) questions whether the instability theory can be reconciled with
109 observations of drumlins with stratified cores of glaciofluvial material (e.g. Easterbook, 1986;
110 Sharpe, 1987). Indeed, the oft-cited complexity of drumlin composition (e.g. Menzies, 1979;
111 Patterson and Hooke, 1995) has frequently been seen as a major obstacle for a unifying
112 theory of their formation, although this pessimism may be misplaced (see Stokes *et al.*, 2011).

113 Whilst the instability theory is principally concerned with the evolution of the ice-till
114 interface, it is important that it can explain observed sedimentary sequences within drumlins
115 (at least qualitatively, but with further progress one anticipates quantitatively). If it is unable
116 to accommodate common sedimentary architectures that are produced or, more commonly,
117 inherited and preserved in a drumlin, then its credibility is damaged. With this in mind, this
118 paper is the first attempt to assess the compatibility of the instability theory of drumlin
119 formation with observations of their composition and internal structure. The underlying
120 principles of the theory are introduced and we then outline a new framework for considering
121 the composition of drumlins, before summarising the main types of drumlin composition
122 reported in the literature (cf. Menzies, 1979; Patterson and Hooke, 1995; Stokes *et al.*, 2011).

123 Given the conceptual basis of the instability theory, we then consider what kinds of sediments
124 and structures might be expected to occur based on firm physical principles, i.e. we consider
125 the formation of drumlins in its physical context and use what we know or can reasonably
126 infer about this context to provide an explanation of what has been observed in the field.

127

128

129 **2. The Instability Theory for Drumlin Formation**

130 *2.1. Underlying Principles*

131 A system can be described as ‘unstable’ when positive feedbacks act to amplify small
132 disturbances, such that small ‘natural’ variations (perturbations) become larger. A simple
133 illustration of this process can be seen on a flat sand surface (e.g. beach) where a small
134 perturbation (e.g. subtle change in sand thickness) encourages local sediment accretion and
135 the growth of a sand ripple. Such instabilities generally grow at different wavelengths and
136 exponentially, at least initially, and tend to grow fastest at a preferred wavelength. This
137 wavelength of maximum growth rate is determined by the physical operation of the system
138 and, significantly, because one wavelength tends to emerge as dominant, the result is often a
139 pattern of similarly sized and spaced ripples (bedforms) in a field. Such relief amplification
140 from an unstable interface is considered a fundamental mechanism for creating
141 bedforms/waveforms into recognisable patterns, e.g. dunes and ripples in aeolian and fluvial
142 landscapes (e.g. Prigozhin, 1999; Fowler, 2011), which resemble subglacial bedforms, see
143 Figure 1. As noted, the regularity of relief amplification (i.e. the spacing of bedforms at a
144 dominant wavelength) arises because an instability in the system determines that one
145 wavelength will usually grow more quickly than others and that patterning will further
146 develop from bedform interactions, e.g. migration, merging, lateral linking, and

147 cannibalisation might push the system towards fewer larger more widely spaced bedforms
148 (Kocurek *et al.*, 2010).

149 Clark (2010) provides a detailed review of the development of ideas relating to the instability
150 theory for drumlin formation, which can be traced back to the seminal work of Smalley and
151 Unwin (1968), who argued that drumlins might be the product of the flow of sediments
152 beneath ice sheets. However, the notion that subglacial bedforms or, more specifically, wave-
153 forms (instabilities), might arise spontaneously from fluid dynamics at the ice-bed interface
154 was first explored by Hindmarsh (1996; 1998a, b), following the ‘paradigm shift’ in
155 glaciology in the 1980s (Boulton, 1986), which recognised the importance of the coupled
156 flow of ice and till and its potential to create bedforms (e.g. Boulton, 1987; Boulton and
157 Hindmarsh, 1987). Further analytical developments were made by Fowler (2000, 2009,
158 2010a) and Schoof (2002a, b; 2007a, b) who confirmed the likelihood of instabilities and,
159 crucially, found that they were largely independent of whether a plastic or viscous till
160 rheology is used, including the highly nonlinear shear-thinning ones typically thought most
161 appropriate for the description of 'nearly plastic' sediment (e.g. Schoof, 2007a).

162 The basic ingredients and underlying principles of the instability mechanism are shown in
163 Figure 2. The base of the ice is assumed to be at the melting point, and producing sufficient
164 water through basal melting that the till is unfrozen and water saturated. It is then assumed
165 that the till will deform if subjected to a sufficiently high shear stress. The model then
166 considers the flow of ice to be Newtonian viscous, that there is a sliding law relating basal
167 shear stress (τ) to basal velocity (u) and basal effective pressure (N : overburden pressure
168 minus till pore water pressure); and similarly that sediment flux (q) is a function of shear
169 stress and till effective pressure. It is worth noting that the work of Dunlop *et al.* (2008) used
170 a non-linearly viscous model of ice flow, with little qualitative effect on the results.

171 It is also assumed that, as a granular material, till will only deform if $\tau > \mu N$, where μ is a
172 coefficient of friction; and it is assumed that τ increases with u and N , while q increases with
173 τ but decreases as N increases. In particular, because the effective pressure in the till increases
174 with depth below the ice-till interface, deformation of the till will be limited to a thin mobile
175 layer whose thickness is expected to lie in the range of tens of centimetres to metres. It is then
176 found that the flow of ice over a level substrate is unstable and, for most reasonable choices
177 of sliding law, bedforms grow and equilibrate at finite amplitude (e.g. Hindmarsh, 1999;
178 Fowler, 2009), the height and length of which are consistent with observations, i.e. 10s of
179 metres (cf. Fowler, 2009). The instability occurs because when ice moves over a shallow
180 bump in the interface, it generates a higher compressive stress on the bump's upstream side
181 than in its lee. If, in addition, the effective sediment viscosity is low compared with that of ice,
182 interfacial velocity remains approximately constant, and this then implies that more sediment
183 flows into the bump than out of it, causing it to grow (Schoof, 2007a). The preceding
184 discussion represents the physical context to the theory and previous work provides further
185 details and justification for these underlying assumptions (e.g. Hindmarsh 1998a; Fowler,
186 2000).

187

188

189 *2.2. Recent developments and current limitations*

190 The theory put forward initially by Hindmarsh (1998a) and Fowler (2000) provides an
191 explicit theoretical mechanism for an instability in the flow of ice over deformable sediments,
192 which can generate a pattern of bedforms from an initially planar surface. It is important,
193 however, to outline the current limitations of the theory, which arise partly from the difficulty
194 of the problem (a physically-based model of drumlin formation is clearly not trivial), but

195 perhaps also from the view that makes genetic distinctions between different ‘types’ of
196 bedforms (e.g. ribbed moraine, drumlins, mega-scale glacial lineations). An alternative view,
197 implicit within the instability theory, is that they sit along a ‘bedform continuum’ (cf. Aario,
198 1977; Rose, 1987), such that in two dimensions there may be little physical distinction
199 between ribbed moraines with a short wavelength in the along flow direction and drumlins
200 with a longer wave-length.

201 In relation to this, the original theory (Hindmarsh, 1998) was a two-dimensional one, and was
202 then used by Dunlop *et al.* (2008) to explain ribbed moraine, but a distinction has often been
203 made between drumlins and ribbed moraine, with the implication being that they are different
204 bedforms, and hence may have different origins. Furthermore, it has been noted that the
205 instability theory has, thus far, failed to generate fully three dimensional drumlins (see
206 discussions in Schoof, 2007a; Pelletier, 2008; Clark, 2010).

207 Clearly, the fact that three-dimensional bedforms have not yet been predicted by the theory
208 represents a significant challenge to it, but recent developments look promising in this regard.
209 Finite amplitude calculations have been undertaken (Fowler 2009) and three-dimensional
210 modelling of evolving ribbed moraine (Chapwanya *et al.* 2011) have generated ‘drumlin-like’
211 culminations (see ‘terrain’ shaded red in Fig. 1f), although the model failed to produce the
212 expected evolution from drumlinised ribbed moraine to just drumlins (Clark, 2010). The
213 differences between these two implementations lies in the assumption by Fowler (2009) that
214 the water pressure in the deforming till layer is at hydrostatic equilibrium with the subglacial
215 stream system, while Chapwanya *et al.* (2011) assumed a slowly relaxing hydrostatic
216 disequilibrium. In particular, the slow relaxation was due to an assumed till permeability of
217 10^{-15} m^2 , comparable with silt. For a sandy till, or if water flows off drumlins by interfacial
218 rivulets, the present low value may be unwarranted. A working hypothesis relevant in the
219 context of this paper would be that drumlins may be stationary if the till is well-drained, but

220 that they move if the till is poorly drained (we shall discuss both possibilities below). Thus, it
221 would seem that the treatment of subglacial water is vital, which is reminiscent of the views
222 of Shaw and co-workers (e.g. Shaw, 1983; Shaw and Kvill, 1984), although for very different
223 reasons. Indeed, the theory has been modified to allow an active subglacial hydraulic system
224 (Fowler 2010), and this development led to the discovery of a rilling instability, which
225 generated features whose lateral spacing is consistent with observations of mega-scale glacial
226 lineations.

227 In summary, it is important to stress that none of the recent implementations of the theory are
228 the last word on the subject, since none of them properly solve the coupled ice-till-water flow
229 problem posed by Fowler (2010b). Thus, we would argue that whilst the theory has yet to
230 produce fully three-dimensional drumlins (they emerge as bumps of finite amplitude in 2D
231 models), it is not yet developed to a state where it could be rejected on this basis.

232

233 *2.3. A new framework for explaining drumlin composition and internal structure*

234 In relation to drumlin composition and internal structure, it is important to stress that the
235 instability theory, thus far, has purposely ignored realistic complications which are
236 nonetheless immaterial to the development of a wavy interface. For example, till is modelled
237 as a homogeneous material and not at the grain-to-grain scale. As such, the model is
238 incapable (in its present form) of making predictions of the kinds of detailed micro-scale
239 sedimentary features that might be generated and observed. This is not necessarily a problem
240 because, as noted by Menzies (1979, p. 350), “it is critical that if any unifying drumlin theory
241 is to be developed it must not be created around unique or special conditions either within the
242 ice mass or drumlin material”. It is important, however, that the theory is not intrinsically
243 contradicted by observations of the commonly observed contents found inside drumlins (e.g.

244 till, stratified glaciofluvial material, etc.) and nor by the structure or architecture of these
245 contents (e.g. conformable/unconformable with the drumlin surface). However, before
246 confronting the theory with observations, we must first consider a new framework which
247 allows for the vertical movement of the unstable (wavy) interface.

248 To date, the instability theory provides a framework in which the bed elevation (s) is
249 described by a partial differential equation, known as the Exner equation, which takes the
250 form (in two dimensions),

251

$$252 \quad s_t + q_x = 0 \quad (1)$$

253

254 In which (t) represents time, (x) represents downstream distance, (q) represents the till flux,
255 and the subscripts denote partial derivatives. Additional assumptions are that ice flow is
256 continuous and constant, the sediment is constantly saturated, the subglacial hydraulic regime
257 is uniform and constant in time, and the sediment supply upstream is constant, and equal to
258 that downstream. These assumptions are made not because we believe them to be true, but
259 because a theory, any theory, has to make some assumptions, and these are the simplest that
260 we can make in the context of drumlin formation. However, in our present intention of
261 addressing how the theory might explain observations of drumlin composition, we have to
262 allow for relaxation of these assumptions, and we now discuss these in turn.

263 In consideration of any specific segment of an ice flow line, it is unlikely that sediment influx
264 and efflux will be in balance. The generalisation of equation (1) to describe such situations is
265 the modified equation:

266

267
$$s_t + q_x = -E \quad (2)$$

268

269 where E represents the net rate of erosion of sediments. We conceive of the actively
270 deforming till layer as having a (fixed) thickness controlled by the depth at which the yield
271 stress is reached as effective pressure increases. Net removal of the sediment above will thus
272 cause entrainment of the sediment below, and thus an effective erosion of the till bed. We
273 thus distinguish two cases: $E > 0$, an erosional environment, and $E < 0$, a depositional
274 environment. Both will play a part in our interpretations of observed drumlin stratigraphy.

275 The theory considers the smooth, continual evolution of the bed under constant external
276 conditions. In reality, since nothing happens during dormant periods, the model can also
277 describe the more likely scenarios where evolution occurs in discrete periods, due to distinct
278 drumlin-forming events. One reason for supposing this is that, as a granular material, till will
279 not deform at all unless the effective pressure (overburden minus hydraulic pressure) N is less
280 than τ / μ , where τ is shear stress and μ is a coefficient of friction. In practice, this means $N <$
281 1 bar. Such low effective pressures are known to occur under ice streams (cf. Kamb, 1991),
282 but may not be common where channelized drainage occurs (as recorded by eskers formed in
283 R othlisberger channels), which typically have much higher effective pressures. Moreover, till
284 deformation implies water saturated sediments, which requires not only that the basal
285 temperature be at the melting point, but also that there is net production of water. So it seems
286 natural to suppose that as an ice sheet evolves, basal conditions change so that drumlins are
287 not built continuously, but episodically, and probably quite rapidly (i.e. few decades: cf.
288 Smith *et al*, 2007), which is a further prediction of the instability theory (see Fowler, 2009).

289 As stated above, the instability theory is still in a state of development and this section
290 indicates our best present understanding. We introduce the minimum ingredients and find

291 what they explain: principally the size and wavelength of subglacial bedforms. To explain
292 further features (their composition and internal structure), we introduce further plausible
293 ingredients. In this preliminary discussion of drumlin composition, we are simply exploring
294 the most likely possibilities that we construe will emerge as the theory is developed.

295

296

297 **3. Observations of Drumlin Composition and Internal Structure**

298 Within the vast drumlin literature, numerous papers (>200) report on their composition and
299 internal structure and it is true to say that they are composed of a range of different sediments,
300 exhibit a variety of different structures (e.g. horizontally stratified versus conformable with
301 landform surface), and show evidence of a variety of styles and extent of deformation (see
302 reviews in Menzies, 1979; Patterson and Hooke, 1995; Hart, 1997; Stokes *et al.*, 2011).
303 Perhaps unfortunately, this diversity has led to a large range of explanations/hypotheses of
304 drumlin formation and it has famously been noted that “there are almost as many theories of
305 drumlin formation as there are drumlins” (Sugden and John, 1976, p. 239). Indeed, although
306 drumlin morphology is also variable (though recently shown to have unimodal distributions
307 of length, width, height and shape: Clark *et al.*, 2009; Spagnolo *et al.*, 2010; 2011; 2012), it is
308 likely that had drumlins only ever been observed to contain the same contents, the “drumlin
309 problem” (Menzies and Rose, 1987; p. 7) would not be so much of a problem.

310 Based on a systematic review of the literature and in an attempt to reduce the oft-cited
311 complexity of drumlin composition, Stokes *et al.* (2011) have recently suggested that there
312 are, essentially, just five basic types, albeit with subtle variants, shown in Figure 3. These are:

313 i. Mainly bedrock

314 ii. Part bedrock/part till

- 315 iii. Mainly till
- 316 iv. Part till/part sorted sediments
- 317 v. Mainly sorted sediments

318 Whilst acknowledging the inherent limitations of such a classification, Stokes *et al.* (2011)
319 argue that explaining these basic types provides a more realistic goal for theories or
320 numerical models of drumlin formation to address. They go on to suggest (as others have
321 done, e.g. Dionne, 1987) that the first type (type 1, Fig. 3), purely bedrock forms, could be
322 viewed as genetically different from drumlins formed of unconsolidated sediment (whaleback
323 may be a more appropriate term, cf. Evans, 1996).

324 Stokes *et al.* (2011) also postulate that because of the unimodal distribution of drumlin
325 dimensions (which suggests a single population of landforms, rather than different types:
326 Clark *et al.*, 2009) and because the other four types of drumlin content can often occur within
327 the same drumlin field (e.g. Hill, 1971), and sometimes in a continuum (e.g. Boyce and Eyles,
328 1991), they are probably genetically related, i.e. their differing contents should not be seen as
329 an obstacle to a unifying theory of drumlin formation. The challenge for the instability theory
330 therefore, is whether it can explain all of the remaining four types of drumlin (listed above).
331 The next section addresses this issue and takes each type of drumlin in turn (excluding purely
332 bedrock forms) and assesses whether the physical principles and mechanisms that underlie
333 the instability theory can explain/predict such observations.

334

335

336 **4. Qualitative and Quantitative Explanations of Drumlin Composition and Architecture** 337 **Predicted by the Instability Theory**

338 *4.1. Drumlins composed of part bedrock/part till*

339 *4.1.1. Observations*

340 Drumlins composed of part bedrock/part till (often called ‘rock-cored drumlins: type 2, Fig. 3)
341 are commonly reported in the literature (e.g. Crosby, 1934; Dionne, 1987; Möller, 1987;
342 Boyce and Eyles, 1991). In their inventory, Stokes *et al.* (2011) list 28 papers that describe
343 this type of drumlin but point out that they are probably far more common than is reported,
344 compared to other types, because of the bias towards sampling drumlins away from regions
345 underlain by crystalline bedrock (see their Table 1 and Figure 17) and partly because they are
346 sometimes called crag-and-tails. The bedrock ‘core’ can be located at the stoss (e.g. Glückert
347 1973), middle (Tavast, 2001), or lee side of the drumlin (Tavast, 2001), although it is most
348 common to be positioned at the stoss side (Stokes *et al.*, 2011). To distinguish these features
349 from ‘crag and tails’, Dionne (1987) suggested that till should account for at least 25% of the
350 entire drumlin volume and cover at least portion of the stoss end.

351

352 *4.1.2. Model Explanation/Prediction*

353 From the point of view of the instability theory, part bedrock/part till drumlins are relatively
354 easy to explain because an instability will form a drumlin due to any small perturbation in the
355 till thickness or bed topography. The instability theory predicts the formation of drumlins as
356 waves which grow from a pre-existing (level) interface. Instabilities grow in nature because
357 there are always perturbations present. Thus, a bedrock protuberance is just an obvious
358 perturbation, and since the theory in one version of its current form (Fowler, 2009) predicts
359 growth of a finite amplitude stationary state, it is a consequence that such perturbations will
360 give rise to drumlins. In short, the theory predicts that bedrock bumps are sufficient but not
361 necessary to seed drumlins. Dynamical analogies abound: the formation of standing waves in
362 rivers at bedrock steps, atmospheric lee waves behind mountains, sand dunes formed behind
363 or in front of obstacles.

364 More specifically, the theory assumes that till flux (q) increases with basal shear stress (τ) but
365 decreases as effective pressure (N) increases (section 2). Thus, depending on the shear stress
366 (which is, in turn, related to ice velocity and effective pressure), it is obvious that till flux
367 varies both spatially and temporally under different conditions. In some circumstances,
368 therefore, till fluxes will be relatively high through the system and in others, they might be
369 relatively low. Given a system of high till fluxes over an underlying surface of bedrock
370 undulations and where till flux from up-ice is insufficient to maintain continuity (or local
371 erosion is too low), it is natural to expect that the till thickness will decrease and bedrock will
372 become exposed at the surface of the till. Depending on the nature of the bedrock surface and
373 the pre-existing sediment thickness, drumlin forms will thus emerge with varying degrees of
374 bedrock 'control' through time, see Figure 4. The situation of bedrock bumps perhaps
375 seeding some drumlins still holds (Fig. 4b), but as till is preferentially removed from the
376 system (erosion dominates over deposition), bedrock bumps are likely to emerge (Fig. 4c and
377 d) and the ultimate progression sees the system evolve to an entirely bedrock surface (Fig. 4e).
378 Similar erosional processes within a deforming bed were envisaged by Boyce and Eyles
379 (1991) in the Peterborough drumlin field, Ontario. They noted drumlins with bedrock cores in
380 areas where the length of time available for erosion was greatest and sediments were thinnest.
381 The observation that bedrock cores might be found in various positions within drumlins (cf.
382 Stokes *et al.*, 2011) is also fairly readily understood. Rapid ice flow over a bedrock bump will
383 cause a large cavity to form in its lee, and in the presence of an adequate till supply, the
384 cavity will be infilled by sediments (Fig. 4d). On the other hand, if the ice flow is relatively
385 slow, then we would expect little cavitation, but till dragged towards the obstacle will pile up,
386 causing a stoss-side cavity infill (which some workers have reported and termed 'pre-crags',
387 e.g. Haavisto-Hyvärinen, 1997). In the absence of plentiful till cover, we may expect bedrock
388 bumps to emerge above the till veneer (and the extreme case of this is the whaleback).

389 Additionally, because effective pressure increases with bed elevation, and thus till mobility
390 decreases, till may simply not be able to reach the summits.

391

392 4.2. *Drumlins composed of mainly till*

393 4.2.1. *Observations*

394 It is no surprise that there are numerous reports of drumlins composed mainly of till (type 3,
395 Fig. 3) (e.g. Wright, 1962; Nenonen, 1994; Menzies *et al.*, 1997; Rattas and Piotrowski,
396 2003). Indeed, Stokes *et al.* (2011) noted that this is by far the most common constituent of
397 drumlins *reported* in the literature and previous studies draw the same conclusion (e.g.
398 Menzies, 1979; Patterson and Hooke, 1995). The emphasis is on '*reported*' because we do
399 not have a large enough sample size to judge whether observations to date are a
400 representative sample of drumlin composition (see discussion in Stokes *et al.*, 2011). In some
401 cases, the entire drumlin appears to consist of an essentially structureless/homogeneous unit
402 of till (Habbe, 1992), whereas others exhibit several units; sometimes horizontally bedded
403 (e.g. Stea and Brown, 1989) and sometimes conformable with the drumlin relief (e.g.
404 Nenonen, 1994). The degree to which till units (or any sedimentary units for that matter) are
405 conformable with the drumlin surface is often viewed as a key issue in drumlinology and is
406 discussed in section 5.1. Drumlins composed mainly of till also show a variety of
407 features/structures related to both ductile and brittle deformation (e.g. Menzies *et al.*, 1997),
408 although others do not and, again, this issue is true for other types of drumlin.

409

410 4.2.2. *Model Explanation/Prediction*

411 In many ways, this type of drumlin is the least problematic for the instability theory. In effect,
412 the instability theory models the surface of the till as a sinusoidal wave of varying thickness

413 (Fig. 2c), but geomorphologists have tended to only map the landform above the mean till
414 surface (e.g. Spagnolo *et al.*, 2012) at some (possibly arbitrary: see Smith *et al.*, 2006), level;
415 and it is these landforms that sedimentologists have tended to focus on in terms of sampling
416 the sub-surface. Because of the dependence of the theory on a mobile till unit that grows in
417 thickness to form the body of the drumlin that scientists map and sample, an obvious
418 prediction of the instability is, therefore, that the high-points of the sinusoidal wave (which
419 we map as drumlins) should be composed of this mobile till unit.

420 That the most commonly reported drumlins are those composed of till would seem to serve
421 the theory well, especially where ice flows over a metres thick sequence of tills. However, it
422 is less obvious how such drumlins can be formed in the absence of a pre-existing deep
423 (several metres) till layer, though whether this is a significant problem needs exploration
424 through mathematical modelling. Two mechanisms emerge from our previous discussion of
425 the instability model. When hydraulic connectivity is poor, we may expect bedforms to grow
426 as travelling waves, and these waves will sweep the underlying sediments together as they
427 move. Alternatively, or as well, in net depositional environments, till is gradually deposited
428 as a thickening layer on top of any pre-existing sediments.

429

430 *4.3. Drumlins composed of part till/part sorted sediments*

431 *4.3.1. Key observations*

432 The second most commonly reported type of drumlin (cf. Stokes *et al.*, 2011), after those
433 composed mainly of till, are those composed of large amounts of both till and sorted (often
434 glaciofluvial) sediments (type 4, Fig. 3). The location of the sorted sediments can vary. In
435 some cases they form a centrally-positioned core or ‘pod’ (e.g. Rattas and Piotrowski, 2003),
436 whilst in others they form a horizontal unit that separates two till units (e.g. Kerr and Eyles,

437 2007) or interbedded with till or *vice versa* (e.g. Whittecar and Mickleson, 1979). Arguably,
438 the most commonly reported architecture, however, is where the sorted sediments simply
439 underlie the till unit (e.g. Clapperton, 1989; Boyce and Eyles, 1991; Habbe, 1992; Jorgenson
440 and Piotrowski, 2003). In this situation, there are reports of the sorted sediments showing
441 evidence of glaciotectonic deformation and being incorporated into the overlying till unit, e.g.
442 drag folds and rafts/lenses of underlying sediments (Boyce and Eyles, 1991) from either
443 ductile or brittle deformation. In other cases, the erosional contact with the sorted sediments
444 may show minimal evidence of deformation (e.g. Habbe, 1992; Hart, 1995a).

445 A commonly reported sub-type of part till/part sorted drumlins are those where the sorted
446 sediments are preferentially found on the lee side of the drumlin. These observations are
447 dominated from locations in Ireland (e.g. Dardis, 1985; Dardis and McCabe, 1983; Dardis *et*
448 *al.*, 1984; Hanvey, 1987, 1989) but not exclusively (see Fisher and Spooner, 1984).

449

450 4.3.2. Model explanation/prediction

451 Whilst drumlins with components of bedrock and till are relatively easy to explain, the
452 presence of stratified sediments is seen by some (e.g. Schoof, 2007a) as introducing
453 additional complexity that is, perhaps, incompatible with the instability theory. As discussed
454 in section 2.2, in the simplest scenario ($E = 0$) the mean interface level remains constant, but
455 this is not an essential ingredient of the model. If there is limited sediment flux from
456 upstream, perhaps because there is exposed bedrock there, or the effective pressure is too
457 high (or shear stress too low) to promote till deformation, then the erosion rate $E > 0$ and the
458 mean level of the interface will lower, even as the wavy interface (i.e. drumlinised surface)
459 evolves.

460 The hydraulic potential of the water at the bed is lowest at the lowest parts of the ice-till
461 interface (i.e. around the base of drumlins and in inter-drumlin areas), and so we expect
462 meltwater to be concentrated there (e.g. Fowler, 2010b), see Figure 5. Moreover, because the
463 low levels of effective pressure associated with sediment deformation are consistent with a
464 description of stream flow as a distributed system (Walder and Fowler, 1994), the most likely
465 form of the basal water flow is a slow trickle through a swamp-like basal platform.
466 Conversely, the effective pressure at the tops of drumlins should be higher, so that the till
467 there is stiffer. The higher/thicker drumlin material will be eroded around the base where the
468 till is softer, and the excavated material can then be removed through the meltwater system.
469 In this way, the transport (erosion) of till is enhanced by both the ice motion-induced
470 sediment flux and through meltwater erosion, with the overall effect of a wavy interface
471 cutting vertically into underlying units. It follows therefore, that a way in which to build part-
472 stratified drumlins as shown in Figure 5d is to first build till-filled drumlins in a depositional
473 environment ($E < 0$), and then later have these drumlins subjected to a net erosional
474 environment ($E > 0$), assuming they overlie pre-existing sorted sediments.

475 Thus, in certain circumstances, the instability theory would predict a wavy interface cutting
476 down into any pre-existing sediments (Fig. 5). The opposite case is where upstream sediment
477 flux is larger than can be excavated out of the drumlin field and, in this way, pre-existing
478 units of sorted sediments may be buried by till units and, in some cases, show evidence of
479 being deformed upwards into the till (e.g. Boyce and Eyles, 1991). In this manner, pods or
480 cores of glaciofluvial material may be incorporated into the till layer and, because they are
481 generally coarser-grained (e.g. sands and gravels) compared to till, such sediments are likely
482 to be better drained and more likely to act as competent material within a deforming layer of
483 till. In this way, they act as boudins around which the deforming till will flow. This idea is
484 not new in the drumlin literature (e.g. Smalley and Unwin, 1968) and was encapsulated most

485 notably in Boulton's (1987) 'theory of drumlin formation by subglacial sediment
486 deformation', see Figure 6. The appeal of the instability theory therefore, is that it may be
487 able to explain drumlins with or without such cores, with the major accomplishment being
488 that such cores (which are not always present: Stokes *et al.*, 2011) are not a necessary pre-
489 requisite.

490 The presence of deformation features at the contact with underlying or en-drumlin units is
491 also to be expected in that any bump created at the ice-bed interface is likely to induce large
492 stress gradients (cf. Morland and Boulton, 1975). More specifically, units with differing
493 rheologies (e.g. till overlying glaciofluvial sediments) are especially conducive to the
494 production of both ductile folds, and faults caused by fracturing of non-yielding material.
495 Thus, whilst deformation fields are likely to be complex and vary from drumlin to drumlin
496 (depending on their constituents), the observed manifestations of deformation are entirely
497 consistent with the instability theory.

498 The mechanical properties of the till itself also creates horizontal variations in its properties.
499 The effective pressure in the till at the ice-till interface increases with elevation of this
500 interface, having a vertical gradient $\Delta\rho_{wi}g$, where $\Delta\rho_{wi}$ is the density difference between
501 water and ice (g is the acceleration due to gravity); while the effective pressure N increases
502 with depth below the ice-till interface at a rate $\Delta\rho_{sw}(1 - \phi)g$, where ϕ is the sediment
503 porosity and $\Delta\rho_{sw}$ is the density difference between sediment and water. Consequently, N
504 will increase along a horizontal level as we move from the stoss face to a position
505 immediately under the crest, and then will decline thereafter to the lee face. The higher value
506 of N must cause lower values of μ ; tensional stresses will be generated towards the lee of
507 drumlins, while compressive stresses are likely to be generated towards the stoss face of the
508 drumlin (cf. Morland and Boulton, 1975), and failure of the till may also lead to thrust faults,

509 which again have also been observed in numerous drumlins (e.g. Hart, 1995a; Hart 1997;
510 McCabe and Dardis, 1994).

511 In terms of stratified sediments found only on the lee-side of drumlins, their presence is
512 usually ascribed to deposition by meltwater that is preferentially routed towards cavities
513 behind the developing drumlin (e.g. Dardis *et al.*, 1984). As noted above, most proponents
514 suggest that such deposition requires a pre-existing drumlin in place and is, therefore,
515 unlikely to explain the drumlin-forming mechanism. In this sense, it is not a drumlin-forming
516 mechanism (although see related arguments in Shaw, 1983; Fisher and Spooner, 1994, etc.).
517 Thus, we simply note that the presence of stratified glaciofluvial sediments on the lee-side of
518 drumlins is to be expected as a result of cavitation, which (although initially regarded as an
519 undesirable feature), the instability model always predicts to occur (see Schoof, 2007a; b;
520 Fowler, 2009), and meltwater routing towards these low pressure cavities.

521

522

523 *4.4. Drumlins composed of mainly sorted sediments*

524 *4.4.1. Key observations*

525 Although they are the least commonly reported drumlins in the literature (cf. Stokes *et al.*,
526 2011), it has been known for a long time that some drumlins are simply composed of sorted
527 (typically glaciofluvial) sediments or have only a thin veneer of till, and they often lack any
528 evidence of widespread deformation (Alden, 1905, Gravenor, 1953; Shaw, 1983; Shaw and
529 Kvill, 1984; Sharpe, 1987; Menzies and Brand, 2007). In many cases, the sorted sediments
530 are horizontally bedded but sorted sediments show a range of architectures and their presence
531 has been attributed to a range of factors. Menzies and Brand (2007), for example, observed
532 pre-existing proglacial and deltaic sediments which acted as an obstacle around which a thin

533 veneer of till was emplaced and which showed evidence of thin-skinned deformation within
534 the till but minimal disturbance of the underlying sediments, see Figure 7. In other cases,
535 evidence of undisturbed sorted sediments has been interpreted as being intimately linked to
536 drumlin formation by subglacial mega-floods, e.g. drumlins represent the glaciofluvial
537 infillings of subglacial cavities (e.g. Shaw, 1983). Such an interpretation assumes that the
538 sediments inside a drumlin are unquestionably linked to the drumlin forming mechanism,
539 which is not always obvious, e.g. Fig. 7 (Menzies and Brand, 2007; Knight and McCabe,
540 1997a; Stokes *et al.*, 2011).

541

542 *4.4.2. Model explanation/prediction*

543 As explained in section 4.3.2, the most obvious way in which the instability theory can
544 explain the presence of sorted sediments (irrespective of whether they form a part of or a
545 whole drumlin) is through the vertical erosion of a deforming till layer into pre-existing
546 sedimentary units. In such cases, the instability theory would have it that the sorted sediments
547 are often unrelated to drumlin formation, other than their potential to act as a stiffer core (see
548 section 4.3.2 and Fig. 6). Whilst it may be easier to conceptualise this down-cutting as
549 producing part till/part sorted drumlins (Fig. 5), it is perhaps more difficult to envisage how
550 the instability theory might explain drumlins composed of mainly sorted sediments and with
551 only a minimal veneer of till and with minimal disturbance of underlying units (e.g. Fig. 7).
552 We conceive of these drumlins forming in the following way.

553 Given an ice sheet building up over a layer of stratified sediments with a largely flat surface,
554 the ice will thicken and it may reach the melting point and begin to produce basal water and
555 slide. The water saturates the underlying sediments, which then begin to deform in a thin (e.g.
556 cm to metre) layer. The thickness of the dilating active layer is not simply a property of the
557 sediments, but is also a consequence of the effective pressure and applied shear stress

558 (Boulton and Hindmarsh, 1987, eqn. (25); Hart *et al.*, 1990, eqn (1)). Thin layers are
559 associated with high effective pressures, which themselves are suggestive of very well
560 drained sediments. As noted in section 2.1, subglacial instabilities are largely independent of
561 whether a plastic or viscous till rheology is used (Schoof, 2007a). Thin dilatant layers are
562 consistent with the instability theory and exhibit instabilities (Dunlop *et al.*, 2008).

563 The instability theory predicts that bedforms grow, and we can expect that as they do so, they
564 begin to obliterate the structure of the underlying sediments. If the underlying sediments are
565 sufficiently porous (e.g. sands or gravels) and the overlying active till layer is sufficiently thin,
566 or non-cohesive, we may expect not only that the effective pressures are relatively high, but
567 also that the water in the till layer to be hydrostatic, and it is in this situation that the evolving
568 drumlins may be expected to be stationary (Fowler 2009). As discussed earlier, this leads to a
569 situation in which the internal sediment architecture is maintained. We thus envisage a suite
570 of bedforms consisting of hard resistant material residing in a basal platform of soft swampy
571 sediments, where the basal water flow is situated. The soft material should be erodible, and as
572 it erodes, we may imagine the drumlins collapsing as their foundations are removed (Fig. 5b).

573 Figure 8 shows the result of a numerical simulation in which this evolution is demonstrated,
574 the details of which are given in the appendix. The initially stratified sediments are indicated
575 by the horizontal coloured bands and, as the bedform descends (its initial range is from -5 to
576 0 m on the vertical axis), the near surface sediments are distorted and move in a thin veneer
577 along the ice/till interface. The figure shows the resulting stratification after a period of ten
578 years, when the wave form has eroded five metres of sediments (horizontal axis also in
579 metres). Such high erosion rates are compatible with recent observations from under an active
580 ice stream in W. Antarctica (Smith *et al.*, 2012). A brief movie of the evolution is included in
581 the supplementary online material.

582 We have included these calculations because, although internal stratification is a simple and
583 inexorable consequence of stationary drumlin formation in an erosive environment, it has
584 caused perceived difficulties with regard to acceptance of the theory (Schoof 2007a; Pelletier
585 2008).

586

587 **5. Discussion**

588 *5.1. Erosional versus depositional drumlins: a false dichotomy?*

589 It is clear from the preceding sections that both the composition and structure of drumlins are
590 important aspects for the instability theory to explain. In particular, the extent to which
591 drumlin composition conforms to the drumlin surface is an issue which has attracted much
592 attention and one which has often led to them being classified as ‘erosional’ or ‘depositional’
593 (cf. Patterson and Hooke, 1995); or ‘destructural’ versus ‘constructional’ (Hart, 1995b;
594 1997).

595 Structures (e.g. layered units of till) that are conformable with the surface of the drumlin form
596 have been noted in a number of studies (e.g. Hill, 1972; Nenonen, 1994; Hanvey, 1992) and
597 are usually interpreted as reflecting a mechanism of formation that involves accretion of
598 material around a core that builds up incrementally, layer by layer (e.g. Fairchild, 1929).
599 Such interpretations are usually supported by reports of clast macro-fabrics that show
600 expected patterns of divergence and convergence around the drumlin as till was deposited
601 and emplaced around a growing obstacle (e.g. Savage, 1968; Goldstein, 1989). As such,
602 drumlins with these surface conformable structures are often referred to as ‘depositional’
603 drumlins. In contrast, those with structures that are unconformable with the surface (e.g. beds
604 of sorted glaciofluvial or till units, e.g. Fig. 7) are, in most cases, assumed to reflect pre-
605 existing material that has been left behind by some form of erosional process – hence the

606 term 'erosional' drumlins. This dichotomy between erosional and depositional drumlins
607 pervades the early drumlin literature and, as Patterson and Hooke (1995) note, "any general
608 theory of drumlin formation must accommodate both possibilities" (p. 33). The alternative
609 requires two theories: one to explain depositional drumlins and one to explain erosional
610 drumlins.

611 We note that although serious questions have been raised over the plausibility of the
612 meltwater flood theory for drumlin formation (see Shaw, 1983; Shaw and Kvill, 1984; and
613 Benn and Evans, 2006 versus Shaw and Munro-Stasiuk, 2006), a notable strength of this
614 theory is its ability to explain both erosional and depositional drumlins. Shaw and co-workers
615 were some of the first to recognise the inherent patterning in glacial landscapes and develop a
616 unifying mechanism to create a surface of bedforms. In this sense, it is similar to the
617 instability theory, with the major difference being in terms of the 'fluid' media through which
618 bedforms are created.

619 Crucially, and like the meltwater flood theory, a strength of the instability theory is that it
620 predicts both depositional and erosional drumlins, depending on whether deposition or
621 erosion dominates in particular settings (e.g. section 4.3: Figure 4 and 5). Where till build-up
622 (deposition) is greater than till transport out of the system then drumlins will build-up, accrete,
623 migrate and deform; and this is likely to result in both homogenous and surface conformable
624 (accretionary) structures depending on the duration of ice and sediment flow and the effective
625 pressures on both developing drumlins and inter-drumlin areas. Generally speaking, these
626 depositional environments might be expected in slower flowing areas and/or towards the
627 margins of an ice sheet, where large sediment depocentres are inherited from previous
628 glaciations. In contrast, where sediment supply is limited from upstream, then the mean level
629 of the ice-bed interface will lower as the higher/thicker drumlin material is eroded around the
630 base. These environments might be more common down-ice from the core areas of an ice

631 sheet where previous glaciations may have stripped away the sediments and/or beneath fast-
632 flowing ice streams, where subglacial erosion is often focussed (cf. Smith et al., 2012).

633 Similar processes of net sediment removal or deposition were described by Schoof and
634 Clarke (2008) in a numerical model of flute formation, which also allowed for observations
635 of both erosional and depositional bedforms. As described above (section 4.3, 4.4 and
636 appendix), net sediment removal will produce drumlins that largely reflect the pre-existing
637 sedimentary units that pre-date drumlin formation (cf. Knight and McCabe, 1997a; Menzies
638 and Brand, 2007) and such units are unlikely to mimic the surface form of the down-cutting
639 drumlinised surface (e.g. Fig. 7). Using these sediments and structures to deduce the
640 fundamental mechanism of drumlin formation is, therefore, largely flawed (Stokes *et al.*,
641 2011).

642 The issue of a depositional versus erosional origin is also linked to the observation that
643 drumlins show varying degrees of deformation structures within them. It is very clear from
644 the preceding discussion that sediments inside drumlins show a range of features that attest to
645 both brittle and ductile deformation and which might occur extensively and throughout the
646 entire drumlin thickness (e.g. Menzies *et al.*, 1997) or in discrete locations or very thin layers
647 (e.g. Menzies and Brand, 2007). Such features might occur in till (e.g. Hart, 1995a; Hart 1997)
648 or initially sorted sediments (e.g. Ellwanger, 1992), or at the interface between the two (e.g.
649 Boyce and Eyles, 1992); or even between till and bedrock (e.g. McCabe and Dardis, 1994).

650 A further strength of the instability theory is that it predicts this range of deformation
651 histories. Where till flux into the system is greater than till flux out of the system, it is likely
652 that bumps (drumlins) will grow through accretionary mechanisms and compressive stresses
653 will develop on their stoss faces, whilst tensional forces develop towards their lee side, and
654 these simple concepts (see also Morland and Boulton, 1975, for a fuller treatment) can
655 explain a range of observed deformation structures that might result during drumlin formation

656 (e.g. boudins, drag folds, thrust faults, etc.). As noted above, however, if the till flux out of
657 the system is greater than that moving into the system then the wavy interface is likely to cut-
658 down through pre-existing sediments to create drumlins as erosional remnants that might
659 show minimal evidence of deformation and/or deformation structures that pre-date drumlin
660 formation. As noted earlier, these ideas are not new in the drumlin literature, with Hart
661 (1995b; 1997) invoking similar scenarios to account for both ‘constructional’ and
662 ‘destructional’ deformation to describe drumlins formed by the build-up or removal of
663 material within a deforming layer.

664 In summary, whilst the concept of erosional and depositional drumlins is useful, it does not
665 justify a view that suggests they have a different mechanism of formation. In this sense, it is a
666 false dichotomy (cf. Schoof and Clarke, 2008) and it is, perhaps, more helpful to view them
667 as end members of a continuum. The appeal of the instability theory is that it accounts for
668 both end members and various intermediate forms along this continuum (Fig. 4 and 5).

669

670 *5.2. Timescales of drumlin formation: inheritance, preservation and prediction*

671 It is clear from numerous studies described above that the material found within drumlins can
672 be related to processes representing several ice flow phases, as well as being inherited from
673 previous sedimentary environments that were not associated with ice flow in any sense. An
674 excellent example of this can be found in Stea and Pe-Piper (1999) who used whole rock
675 geochemistry to locate the source of igneous erratic material inside drumlins in Nova Scotia.
676 This provenance analysis revealed that the material inside the drumlins was delivered by at
677 least two ice flow phases with different source areas. Furthermore, Stea and Brown (1989)
678 noted that till units in some drumlins represented erosional remnants from older drumlins,
679 around which material was emplaced, see Figure 9. These processes are also expressed in the

680 surface morphology of drumlin fields, many of which are known to reflect a palimpsest
681 landscape of ‘cross-cutting’ ice flow landforms (cf. Clark, 1993). Furthermore, there are
682 several reports of different types of bedforms being superimposed on each other, e.g. ribbed
683 moraine superimposed on top of drumlins (Dunlop and Clark, 2006) or mega-scale glacial
684 lineations (Stokes *et al.*, 2008); or drumlins superimposed on top of ribbed moraine (e.g.
685 Dunlop and Clark, 2006; Knight and McCabe, 1997b; Hättestrand and Kleman, 1999) or
686 mega-scale glacial lineations (Clark, 1993). Thus, whilst it is clear that some material inside
687 drumlins might be unrelated to the drumlin forming mechanism (the ‘erosional’ drumlins
688 described above in section 5.1) it is also important to appreciate that material may also reflect
689 a time-integrated record of several ice flow phases and bedforming events, which
690 observations clearly support (e.g. Stea and Brown, 1989).

691 A key prediction of the instability theory (see Fowler, 2009) that is entirely consistent with
692 these observations is that the time-scale for growth of drumlins is of the order of years (see
693 also Smith *et al.*, 2007), whereas the time-scale of ice sheet occupation and associated
694 changes in ice dynamical behaviour is of the order of hundreds to thousands of years. An
695 obvious consequence is that drumlins are likely to be remoulded by episodic changes in ice
696 flow direction and, in some cases, completely erased. The instability theory would thus
697 predict that, under most circumstances, pre-existing drumlin sediments will form the cores of
698 drumlins from a younger ice flow, if the vertical erosion of the wavy interface or the till flux
699 into the system is not great enough to remove them altogether. In other cases, all evidence of
700 pre-existing drumlins might be removed and yet in other cases, the time window for drumlin
701 formation might be so small as to conduct minimal landform/bedform creation, leaving a pre-
702 existing drumlin field barely modified at all.

703 These simple concepts suggest that the most important factors controlling drumlin
704 composition and internal structure are: (i) the pre-existing sediments; (ii) the balance between

705 till deposition and erosion; and (iii) the timing and duration of ice flow, which includes
706 episodic changes in flow direction, basal thermal regime, and subglacial water conditions.
707 These controls are encapsulated in Figure 10. If we take a common situation often found
708 towards the margins of an ice sheet with thick layers of till overlying pre-existing
709 glaciofluvial sediments (bottom left in Fig. 10a), the system could evolve to produce
710 drumlins with quite different contents based on the duration and variability of ice flow (Fig.
711 10a). A similar situation might occur with till overlying bedrock (Fig. 10b). Indeed, these
712 ideas are applicable to a range of settings and, furthermore, provide a useful framework for
713 interpreting drumlin composition (e.g. Boyce and Eyles, 1991).

714 With the above in mind, it is possible to make some general predictions about where, under a
715 continental ice sheet, drumlins with different compositions are more likely to occur based on
716 the pre-existing substrate (largely influenced by previous cycles of ice sheet
717 erosion/deposition). Figure 11 shows a simplified terrain from a previous glaciation with a
718 core of pre-existing crystalline bedrock (zone 1) surrounded by a transitional zone of bedrock
719 and thin till (< few metres: zone 2) that progressively thickens towards the ice sheet margins
720 (zones 3 and 4). Such a terrain is not dissimilar to an idealised Laurentide Ice Sheet bed
721 during the Late Pleistocene, which is thought to have changed from an all soft-bedded to a
722 mixed hard-soft bedded ice sheet during the Middle Pleistocene through glacial erosion of a
723 thick regolith and resulting exposure of unweathered crystalline bedrock (Clark and Pollard,
724 1998). Given the waxing and waning of ice sheets during both glacial-interglacial and stadial-
725 interstadial time-scales, it would also be expected that the periphery of the previous ice sheet
726 extent would be characterised by thicker sequences of proglacial/deglacial sediments for
727 subsequent overriding (zone 5).

728 Given the pre-existing terrain in Figure 11, it is not hard to predict (at a general level) which
729 of the main types of drumlins might be expected to occur in each zone as a result of the last

730 glaciation. Zone 1 would be characterised by purely bedrock forms (type 1; Fig. 3); zone 2 by
731 part bedrock/part till (type 2); zone 3 by mainly till (type 3); zone 4 by part till/part sorted
732 sediments (type 4); and zone 5 by mainly sorted sediments (type 5). This is, of course, a
733 generalisation and localised variations are bound to exist but it emphasises the importance of
734 the pre-existing substrate conditions in influencing drumlin composition.

735 Unfortunately, it is not yet possible to ascertain whether the spatial patterns in drumlin types
736 shown on Fig. 11 actually exist (e.g. on the Laurentide Ice Sheet bed). This is because
737 observations of drumlin composition are typically restricted to just a small sample of
738 drumlins within a particular drumlin field, and because most observations are tightly
739 clustered towards the southern marginal areas of the last mid-latitude ice sheets (see Fig. 17
740 in Stokes *et al.*, 2011). However, observations from these regions suggest that ‘mainly till’
741 (type 3) and ‘part till/part sorted’ (type 2) are by far the most commonly reported (Stokes *et*
742 *al.*, 2011), which may lend support to the ideas encapsulated in Figure 11. Furthermore,
743 whilst there are very few observations of drumlin composition from interior regions of former
744 ice sheets, reports from Fennoscandia suggest that drumlins dominated by bedrock
745 components (zone 2, type 2 on our Fig. 11) are commonly found towards interior zones (e.g.
746 Glückert, 1973; Minell, 1979; Möller 1987; Haavisto-Hyvärinen, 1997).

747

748

749 **6. Summary and Conclusions**

750 The instability theory proposes that a range of subglacial bedforms (including drumlins) arise
751 from an instability that occurs at the ice-bed interface as the result of the coupled flow of ice
752 and till and is one of the few explanations to be described in physically-based
753 numerical/analytical models (cf. Fowler, 2000; 2010a; Hindmarsh, 1998a; b; Schoof, 2002b;

754 Chapwanya *et al.*, 2011). Predictions from these models have been shown to match
755 observations of bedform dimensions (e.g. Dunlop *et al.*, 2008; Fowler, 2000; 2009; 2010b;
756 Chapwanya *et al.*, 2011). A key strength of the instability theory, therefore, is that it offers a
757 promising unifying explanation for a range (continuum?) of subglacial bedforms (cf. Aario,
758 1977; Rose, 1987) of which drumlins are the most ubiquitous. Observations of drumlin
759 composition and internal structure, however, are incredibly diverse and this is often seen as a
760 major obstacle to a unifying theory. In this paper, we have compared the key observations of
761 drumlin composition and internal structures in the literature and considered, theoretically,
762 how they might arise based on firm physical principles that form the basic ingredients of the
763 instability theory. Contrary to a view which suggests that certain observations (e.g. the
764 presence of undeformed stratified sediments) are inconsistent with the instability theory, we
765 suggest that one would actually *expect* a range of drumlin constituents, including at least
766 some occurrences of drumlins with stratified cores of glaciofluvial material.

767 In terms of the five main types of drumlin composition (Figure 3) identified in the literature
768 (Stokes *et al.*, 2011) and excluding mainly bedrock forms (type 1) the instability theory
769 suggests:

- 770 • Drumlins composed of part-bedrock/part till (type 2) occur because bedrock bumps
771 act as perturbations that give rise to drumlins
- 772 • Drumlins composed of mainly till (type 3) occur because of the dependence on a
773 mobile till unit that grows in thickness in a depositional environment to form the body
774 of the drumlin
- 775 • Drumlins composed of part till/part sorted sediments (type 4) occur through the
776 advection of till across and erosion into pre-existing sorted sediments and around
777 cores of sorted sediments

778 • Drumlins composed of mainly sorted sediments (type 4) through the vertical erosion
779 of both till and meltwater into pre-existing sorted sediments at the wavy ice-bed
780 interface

781 Related to the above, and with specific reference to the structure of the sedimentary units, the
782 instability theory predicts:

783 • Drumlins which are built by successive episodes of till influx and deposition will
784 naturally build a structure in which the separate till units are conformable with the
785 drumlin surface, i.e. accretionary or depositional drumlins where more sediment flows
786 into bumps than out of them

787 • Drumlins with internal structures that are unconformable with the drumlin surface in
788 conditions where till flux out of the system is greater than till supply into the system,
789 i.e. erosional cores may be preserved as the wavy interface cuts vertically downwards

790 • Drumlin formation and shaping can occur rapidly (few decades), such that changes in
791 ice flow direction will lead to inherited cores from previous flow directions.

792 Within the framework of the instability theory, the varied content of drumlins can be
793 explained by three key factors: (i) the pre-existing sediments; (ii) the balance between till
794 erosion and deposition; and (iii) the variability and duration of ice flow. These simple
795 concepts offer an interpretative and predictive framework for where specific types of drumlin
796 composition might be found on an ice sheet bed and how they might be interpreted in terms
797 of ice dynamics and sediment flux. We conclude that the instability theory represents the
798 most promising solution to the ‘drumlin problem’ thus far and offers a unifying explanation
799 for the creation of a range of subglacial bedforms.

800

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808

809

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1072

1073

1074 **Figure Captions:**

1075 **Figure 1:** Instabilities create familiar patterned surfaces of bedforms in a range of
1076 environments not only restricted to Earth: (A) HiRISE image of the Herschel Crater, Mars,
1077 showing a field of barchan dunes (from Kocurek *et al.*, 2010); (B) hillshade image from
1078 LiDAR-derived DEM of laterally linked barchans dunes to form crescentric dunes in the
1079 White Sands Dune Field, New Mexico (from Kocurek *et al.*, 2010); (C) multi-beam
1080 bathymetric images of Mississippi River bed at Audubon Park, Louisiana, during high
1081 discharge of $34,300 \text{ m}^3 \text{ s}^{-1}$ (left) and low discharge of $8,900 \text{ m}^3 \text{ s}^{-1}$ (right) (from Kocurek *et*
1082 *al.*, 2010); (d) DEM of down-ice transition from barchan-like ribbed moraine (left) to
1083 drumlins (right) in north central Ireland (from Clark, 2010); (E) aerial photograph of
1084 classical-type ribbed moraine located at Lake Rogen, Härjedalen, central Sweden; (F)
1085 modelled subglacial bedforms using the Hindmarsh-Fowler instability theory as formulated in
1086 Chapanwanya *et al.* (2011), taken from Clark (2010). Note that these modelled features are
1087 almost identical to the ‘real’ ribbed moraine in (E), in terms of their dimensions and
1088 wavelengths, with ‘drumlin-like’ culminations appearing in red shading ($\sim 1 \text{ m}$ high) after 50
1089 years).

1090

1091 **Figure 2:** Schematic diagram showing the basic ingredients and underlying principles of the
1092 instability theory. When ice and sediment are allowed to deform and sliding can occur at the
1093 ice-till interface (a), the system is prone to the development of an along-flow instability
1094 which creates waveforms (bedforms) at the ice-till surface (b) that emerge as drumlins of
1095 dominant wavelength (c). See section 2 for detail.

1096

1097 **Figure 3:** Schematic illustration of the five main types of drumlin (and sub-types) identified
1098 in a systematic review of the literature reported in Stokes *et al.* (2011), who further suggest

1099 that purely bedrock forms (type 1) should be referred to as symmetrical or asymmetrical
1100 whalebacks, rather than drumlins (cf. Dionne, 1987; Evans, 1996).

1101
1102 **Figure 4:** Schematic illustration of how drumlins can emerge with varying degrees of
1103 bedrock ‘control’ through time in a system where till continuity is inhibited, e.g. where till
1104 transport > till supply. The instability theory predicts that the ice-till interface, under certain
1105 circumstances, is unstable and becomes wavy (Fig. 2). Depending on till continuity (i.e. the
1106 balance between till transport within the deforming layer and till supply from erosion and or
1107 advection from up-ice) the wavy interface can cut downwards. Given a setting with a metres-
1108 thick layer of pre-existing till overlying bedrock, drumlins will emerge from the instability
1109 with minimal bedrock control (A). As till is removed from the system (because till transport >
1110 till supply), the till thickness will be reduced and some drumlins will be anchored by pre-
1111 existing perturbations that might act as cores in a variety of locations (B). Further till
1112 exhaustion might lead to more obvious drumlin cores, as in (C), crag-and-tail features (D),
1113 and an entirely bedrock surface (E).

1114
1115 **Figure 5:** Schematic illustration of how drumlins can emerge with stratified cores of
1116 glaciofluvial material (or similar units) through time in a system where till continuity is
1117 inhibited, e.g. where till transport > till supply. Given a wavy interface, the effective pressure
1118 (N) is predicted be highest at the tops of the drumlins and lowest around the base of the
1119 drumlin and in inter-drumlin areas, where the hydraulic potential of the water is also likely to
1120 be lowest and where meltwater is likely to be concentrated (A). We envisage a slow trickle
1121 through a swamp-like basal platform that erodes around the stiffer drumlins (B). As in Figure
1122 4, if the wavy interface cuts vertically into pre-existing sediments, it is likely that pre-existing
1123 sediments will be incorporated into the drumlins and one might envisage a situation evolving

1124 through time from drumlins composed of mainly till (C); part till/part sorted sediments (D),
1125 to mainly sorted sediments (E). This schematic illustrates pre-existing glaciofluvial sediments
1126 but applies to any pre-existing sedimentary units (e.g. till units, deltaic deposits, etc.).

1127

1128 **Figure 6:** Schematic illustration of how pre-existing sediments may influence drumlin
1129 formation (from Boulton, 1987). In this case, areas of ‘stiffer’ well-drained glaciofluvial
1130 material act as cores around which till deforms. Clark (2010) refers to these as drumlin clones
1131 (or obstacle drumlins where an obvious bedrock protuberance occurs) to distinguish them
1132 from drumlins formed purely from the instability (emergent drumlins). The formation of
1133 drumlin clones is consistent with the instability theory but a further appeal of the instability
1134 mechanism is that it can also account for drumlins without obvious cores based on purely
1135 fluid dynamical principles.

1136

1137 **Figure 7:** Cross section of the Port Byron drumlin, New York State, USA (redrawn from
1138 Menzies and Brand, 2007) that clearly illustrates the presence of mainly stratified sediment
1139 overlain with only a thin veneer of till. These observations clearly show minimal disturbance
1140 of pre-existing sediments that are unrelated to drumlin sediments but probably acted as a
1141 stiffer core (see Fig’s 5 and 6).

1142

1143 **Figure 8:** Evolution of an initially stratified layer of sediments in an erosive environment, as
1144 described in the appendix and shown schematically in Figure 5. Horizontal and vertical axes
1145 are in metres and parameters used are for a 5 m high drumlin eroding down at $E = 0.5 \text{ m y}^{-1}$
1146 for 10 years. Deformable till depth (d_T) = 0.5 m and ice velocity (u_0) = 18 m y^{-1} . The
1147 drumlin profile is $S = \frac{1}{2}a_0 \cos kx$ where a_0 is the drumlin height, x is the distance along
1148 flow in metres, and k is the wave-number ($= 2\pi/l$, where l is the length (period) of the

1149 drumlin in metres). A brief movie of the evolution is available online as supplementary
1150 information.

1151

1152 **Figure 9:** Changes in ice flow direction over time are likely to result in complex
1153 stratigraphies within drumlins, where pre-existing material from older bedforms may be
1154 wholly or partially removed. This is detailed in Stea and Brown (1989) who interpreted
1155 material from some drumlins in central Nova Scotia as relicts from older bedforms, shown in
1156 (A). Shaded areas under stratigraphy and form are thought to represent till units formed at the
1157 same time as the drumlin shaping process whereas unshaded areas under stratigraphy
1158 represent erosional remnants of earlier units (redrawn from Stea and Brown, 1989). A
1159 satellite image of cross-cutting bedforms from Wollaston Peninsula, Victoria Island,
1160 Canadian Arctic Archipelago, is shown in (B), depicting three populations of drumlins
1161 (selected bedforms highlighted with red (inferred oldest), green and yellow (youngest) lines:
1162 from Stokes *et al.* (2006).

1163

1164 **Figure 10:** The instability theory suggests that the most important determinants of drumlin
1165 composition are: (i) the pre-existing sediments (shown bottom left in each panel); (ii) the
1166 balance between till erosion and deposition (y-axis on each panel); and (iii), the duration of
1167 flow (x-axis on each panel). Variations in the above are predicted to produce a variety of
1168 drumlin types from initial substrate conditions, depicted here as till overlying sorted
1169 sediments (A) and till overlying bedrock (B), although the concepts apply to any pre-existing
1170 terrain. The drumlin types refer to those described in section 4 and Figure 3.

1171

1172 **Figure 11:** Schematic illustration of the predicted occurrence of drumlins with different
1173 composition under an idealised ice sheet, which bears some similarity to the Laurentide Ice

1174 Sheet bed, but is used to make the general point that drumlin composition is likely to largely
1175 reflect pre-existing sediments and their position in relation to the ice sheet margin. The
1176 drumlin types refer to those described in section 4 and Figure 3.

Appendix

In order to simulate the evolution of subsurface stratified sediments under the evolution of the instability which causes drumlins to grow, it is necessary to specify a subsurface transport field. In the development of the theory (Fowler 2009), no reference was made to any specific rheology, other than that the sediment flux q was assumed to increase with increasing basal stress τ and decrease with increasing effective pressure N . In a two-dimensional region (coordinates x and z) of sediment bounded above by the ice-till interface at $z = s$, we may write

$$q = \int_{-\infty}^s u dz, \quad (\text{A.1})$$

where the till velocity has horizontal and vertical components u and w .

In a complete theory such as that of Fowler (2009), we derive an evolution equation for s based only on a prescription for q . In order to facilitate our present objective, we will *prescribe* $s(x, t)$, and use its form to infer subglacial sediment transport patterns, based on a realistic assumption about the till velocity. Specifically, we make the assumption that

$$u = u_0 \exp[-b(s - z)], \quad (\text{A.2})$$

where u_0 is the sliding velocity at the ice-till interface, and may be taken to be constant.¹ The exponent b measures the depth ($\sim b^{-1}$) of the deforming till layer, and we expect values $b \sim 1 \text{ m}^{-1}$, although necessarily, b cannot be constant. In fact, the assumption (A.2) implies that

$$b = \frac{u_0}{q}. \quad (\text{A.3})$$

Notice in particular that q must remain positive (as we expect).

In the present situation, we are interested in the case where net erosion of the sediment causes downcutting of the ice into the sediment, and in this case we pose a modified form of the Exner equation as

$$s_t + q_x = -E, \quad (\text{A.4})$$

where E represents a net erosion rate with units of metres per year. We do not conceive of this erosion as being the plucking and grinding of bedrock (which would not present such a term in the Exner equation), but rather a superfluous removal term by subglacial stream flow.

Given the horizontal velocity u in (A.2), we can solve for w to find (bearing in mind that the kinematic condition at $z = s$ is $w = s_t + us_x + E$)

$$w = us_x + (s_t + E)[1 + b(s - z)] \exp[-b(s - z)], \quad (\text{A.5})$$

¹Equation (A.2) is of course inconsistent with a finite thickness of deforming till, but the distinction is only cosmetic, and the present assumption is made purely for algebraic convenience.

and individual particles can be tracked by solving the ordinary differential equations

$$\dot{x} = u, \quad \dot{z} = w. \quad (\text{A.6})$$

In particular, we can track the evolution of different layers of sediment by assigning a variable c which is an indicator for material content. For example, we might take $c = -1$ for clay, $c = 0$ for sand, and $c = 1$ for till. If (ξ, ζ) marks the initial location of a particle, then the sediment type is given by a function

$$c = c(\xi, \zeta), \quad (\text{A.7})$$

and ξ and ζ are the initial values at $t = 0$ for x and z , i. e.,

$$x = \xi, \quad z = \zeta \quad \text{at} \quad t = 0. \quad (\text{A.8})$$

In practice we sequentially plot the surface (x, z, c) parametrically at successive times in terms of the parameters ξ and ζ , using Matlab's `scatter` command.

Choice of parameters

To be specific, we choose the interface position s to be given by the function

$$s = -Et + a \cos k(x - vt), \quad (\text{A.9})$$

where a is the interfacial amplitude, and v is the interfacial wave speed. Generally, a is a function of time, and a representative choice is the function

$$a = \frac{1}{2}a_0(1 - e^{-rt}), \quad (\text{A.10})$$

where a_0 is the final drumlin elevation, and r is a measure of the growth rate.² The wavenumber k is defined in terms of the wavelength l by

$$k = \frac{2\pi}{l}, \quad (\text{A.11})$$

and the inlet sediment flux is chosen as

$$q_0 = \frac{1}{2}u_0d_T, \quad (\text{A.12})$$

where d_T is an estimate of deformable till thickness. From (A.4) and (A.9), we have the expression for q ,

$$q = q_0 - \frac{\dot{a}}{k}[\sin(k(x - vt) + \sin kvt)] + av[\cos k(x - vt) - \cos kvt], \quad (\text{A.13})$$

where we apply the condition $q = q_0$ at $x = 0$.

²More realistic choices might be made to reflect initial exponential growth, but there is little purpose to this.

Symbol	Meaning	Typical value
a_0	elevation	10 m
b^{-1}	till shear exponent	$\sim d_T \sim 1$ m
d_T	till deformation thickness	1 m
E	erosion rate	0.1 m y ⁻¹
l	wavelength	300 m
q_0	upstream sediment flux	50 m ² y ⁻¹
r^{-1}	growth time scale	10 y
u_0	sliding velocity	100 m y ⁻¹
v	wave speed	0, 50 m y ⁻¹

Table 1: Typical values of the parameters in the model.

Table 1 gives our estimate of typical values of the parameters. An awkwardness occurs in making the simulations. Because q must remain positive, we see from (A.13) that we must have $q_0 > \frac{2\dot{a}}{k}$, and thus the growth time

$$r^{-1} > \frac{a_0}{\pi d_T} \frac{l}{u_0}, \quad (\text{A.14})$$

which is roughly the time $t_D = \frac{l}{u_0}$ for a sediment particle at the ice-till interface to move one drumlin length. Equally, it is necessary that $q_0 > a_0 v$, and thus

$$a_0 < \frac{u_0 d_T}{2v}. \quad (\text{A.15})$$

A more thorough analysis of (A.13) shows that the precise condition is that both (A.14) and (A.15) must be satisfied, or simply

$$a_0 < d_T \min\left(\frac{u_0}{2v}, \frac{\pi u_0}{rl}\right). \quad (\text{A.16})$$

The awkwardness lies in the fact that since we typically expect the growth time $r^{-1} \sim \frac{l}{u_0}$, and for wave instabilities, $v \sim u_0$, the constraint on amplitude is that $a_0 \lesssim d_T$, yielding unnecessarily small amplitude drumlins. This point was one of Schoof's (2007a) objections to the instability theory of drumlin formation, and can be seen to be a purely kinematic consequence of the Exner equation. In reality, cavities form for $a_0 > d_T$, and the Exner equation cannot be applied in the same way (Fowler 2009).

We have sidestepped this issue here by considering only the case where the waves are stationary ($v = 0$) as found by Fowler (2009). It is essentially obvious that a downcutting drumlin with only a thin veneer of mobile till will maintain subsurface

stratification; but an illustration nevertheless illuminates the point. It is also obvious that a travelling drumlin will churn up the subsurface sediments, and in fact we consider this to be a mechanism to provide till-formed drumlins, despite only having near-surface mobility. As explained in the text, we may associate travelling drumlins with non-hydrostatic water pressure, i. e., less well-drained material (Chapwanya *et al.* 2011). Figure 8 in the main text shows the results of a computation as described above.

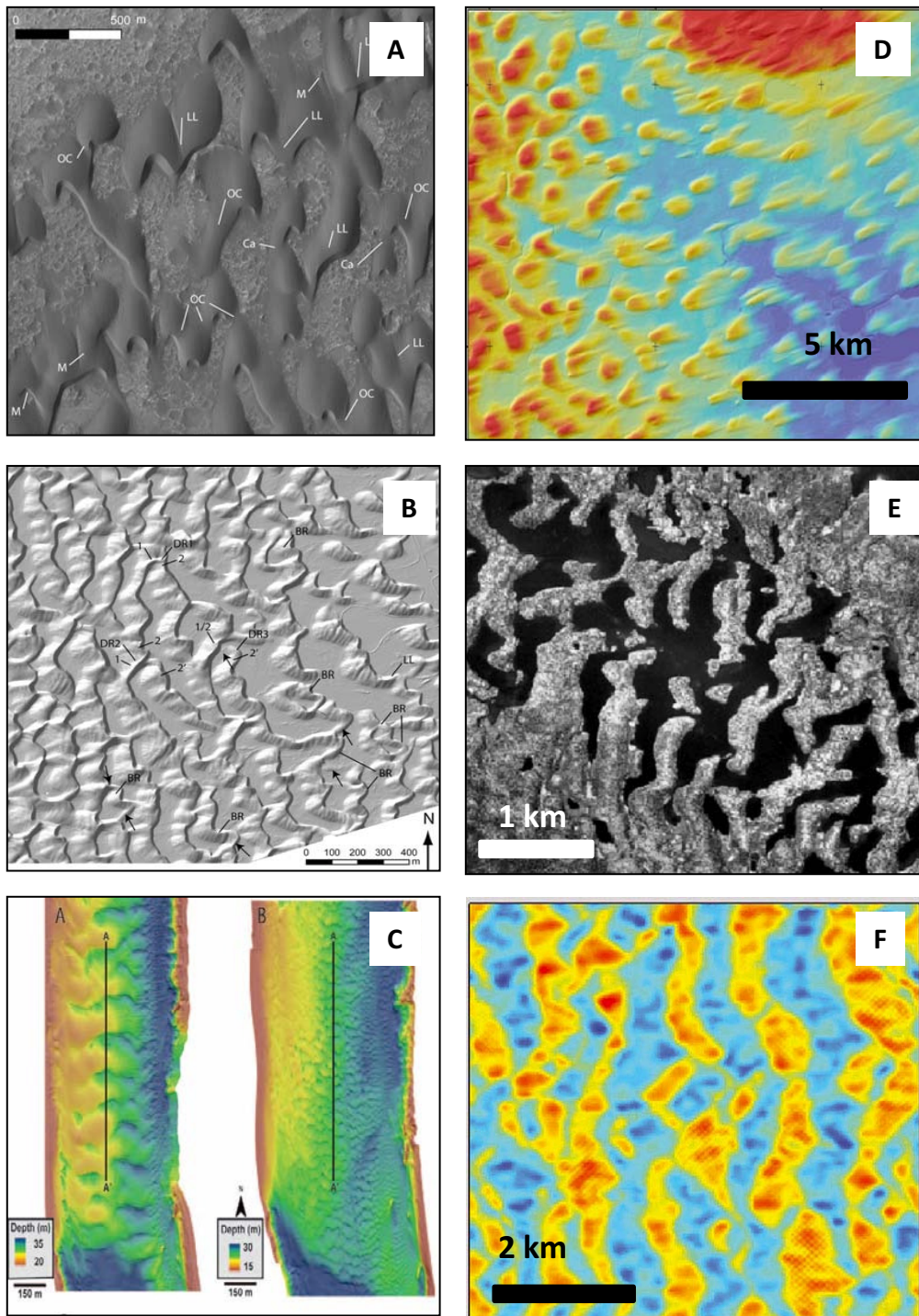
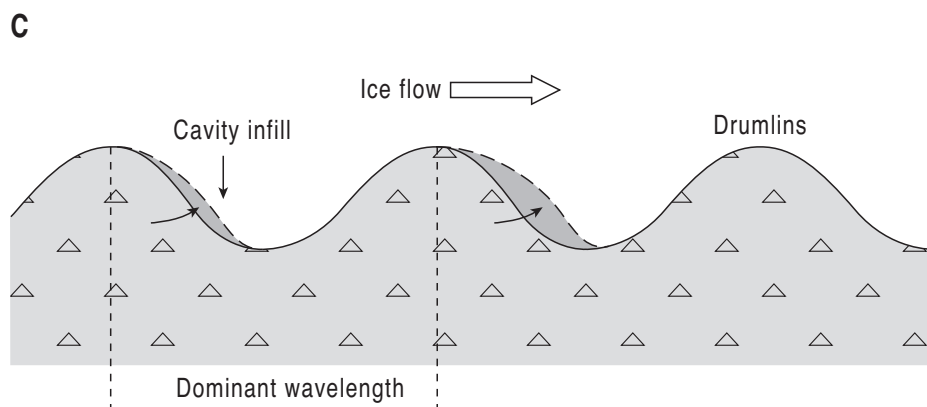
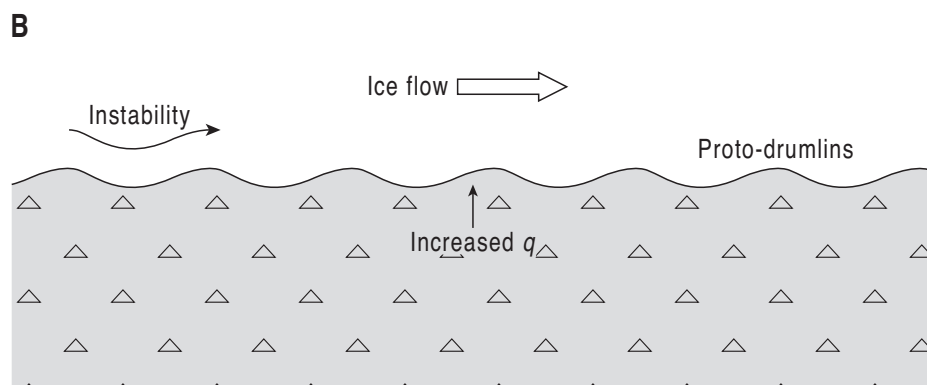
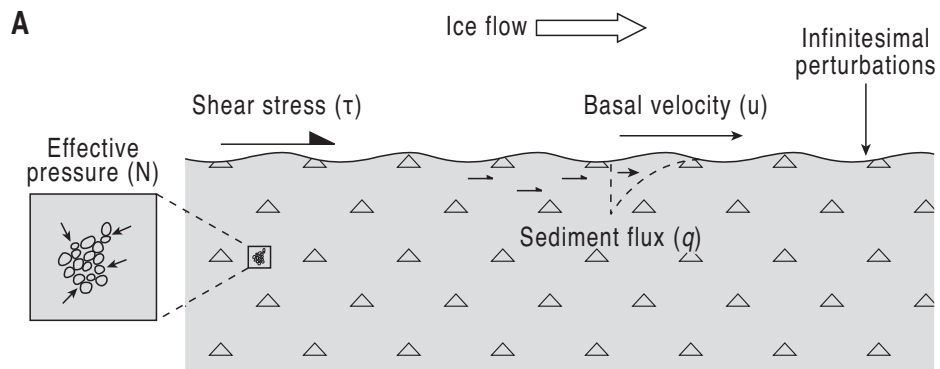
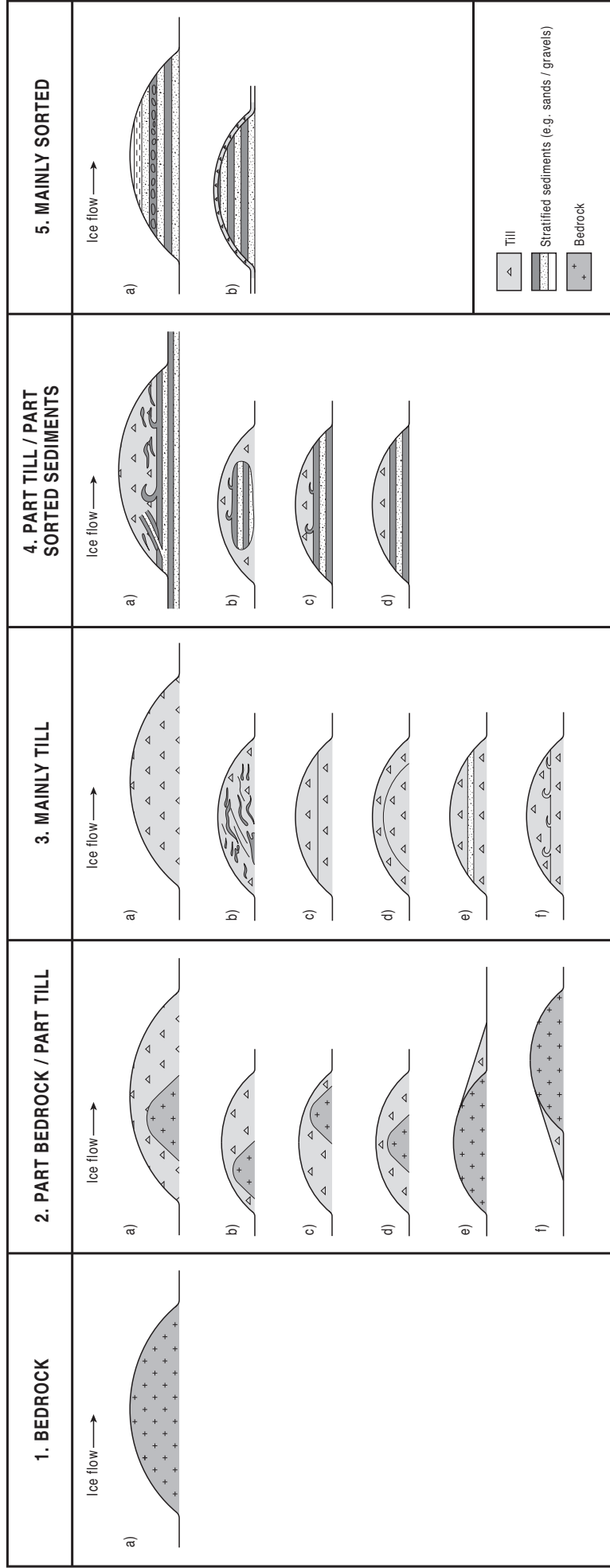
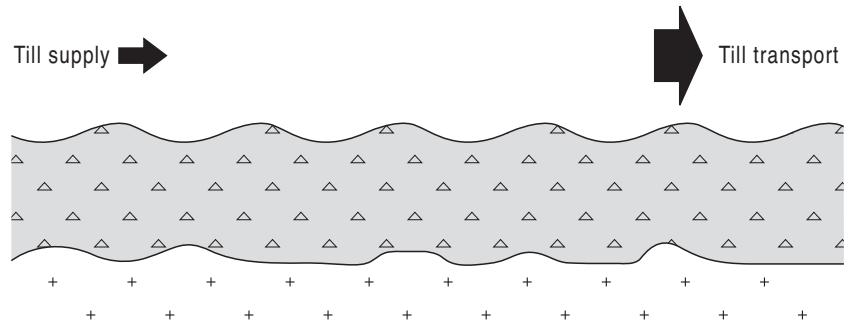


Figure 1

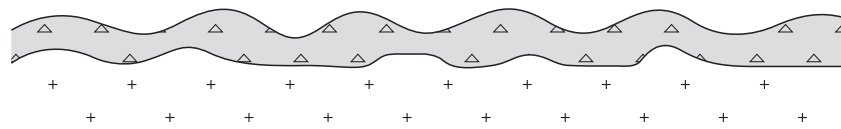




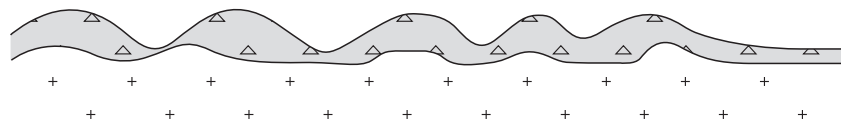
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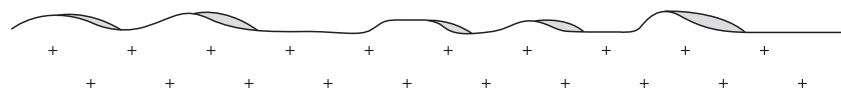
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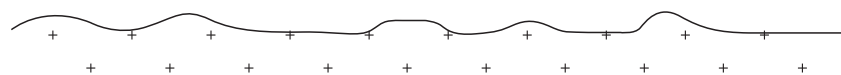
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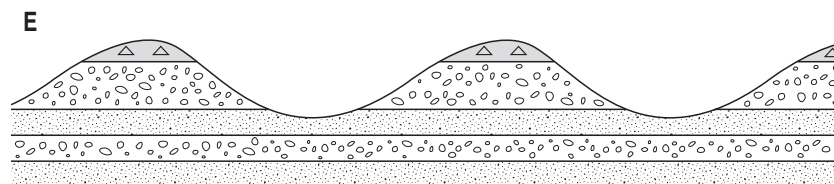
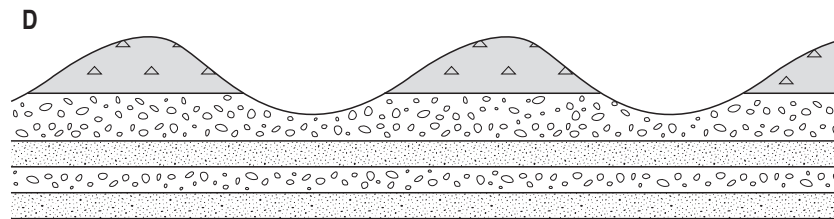
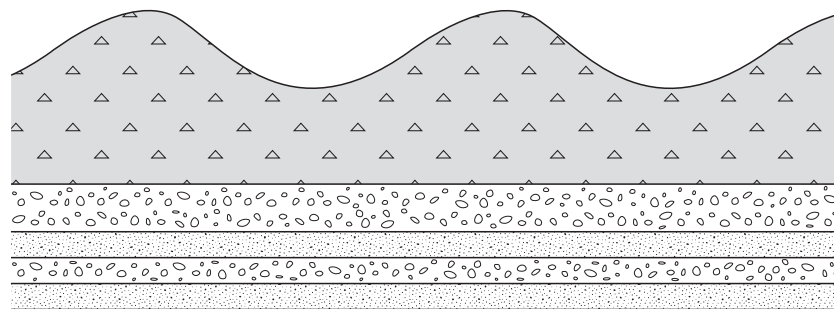
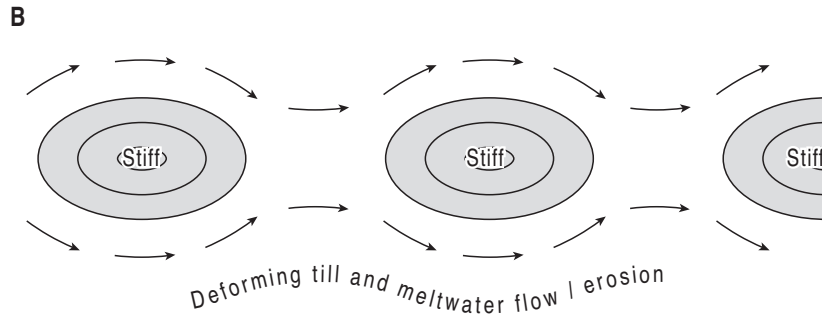
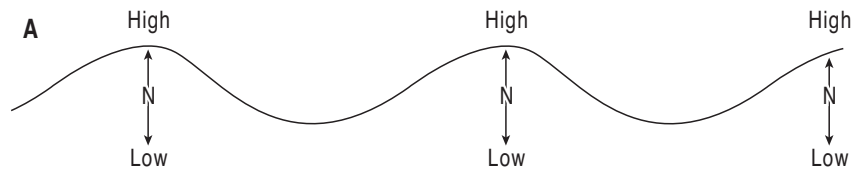
D)



E)



△ Till + Bedrock



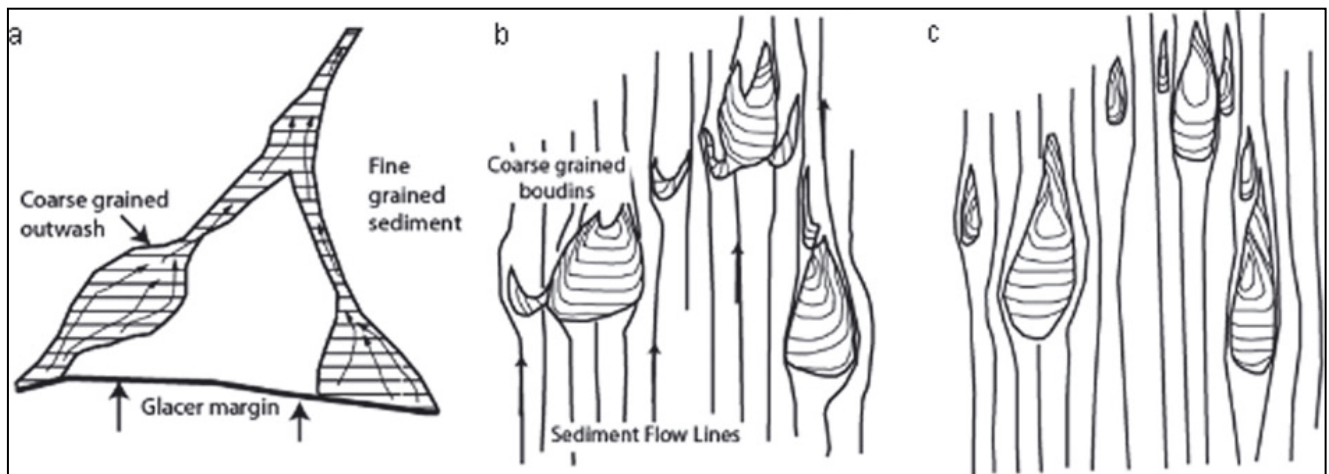
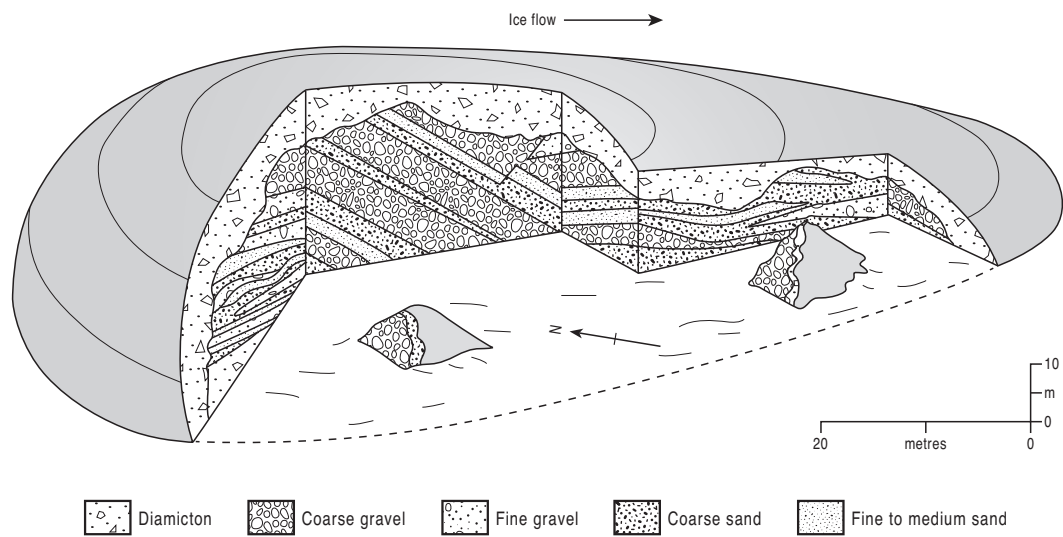
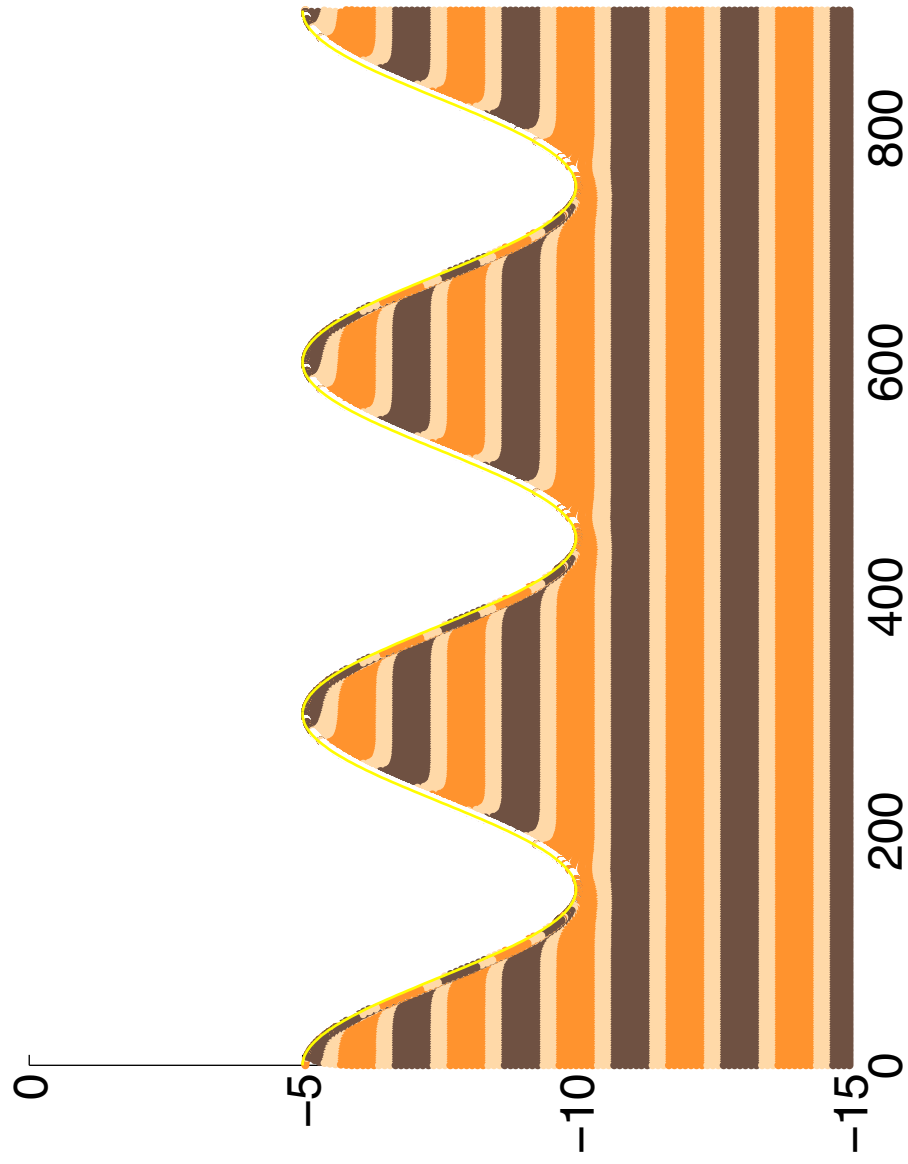


Figure 6

PORT BYRON DRUMLIN



$t = 10.0000$



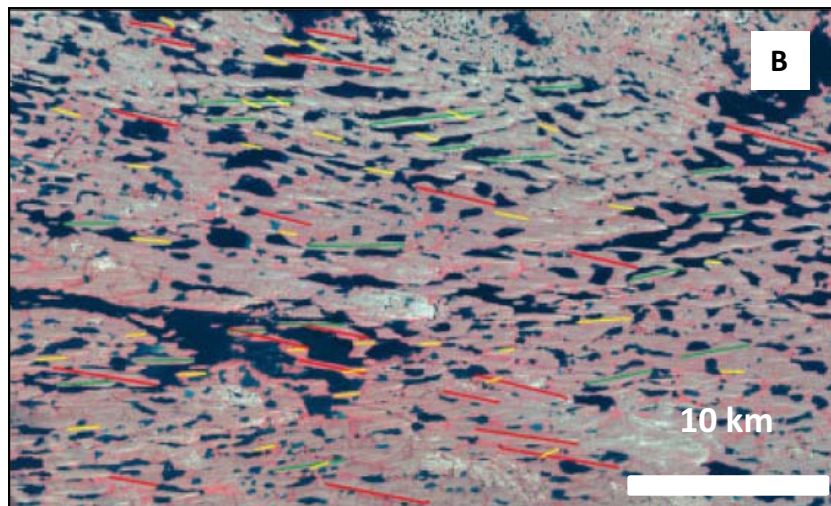
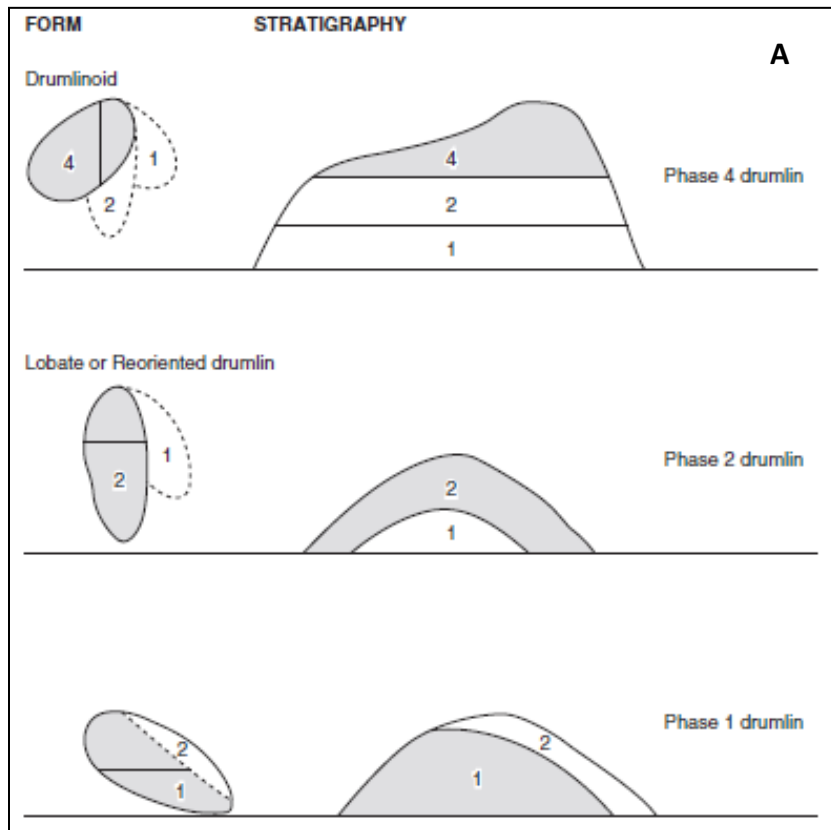
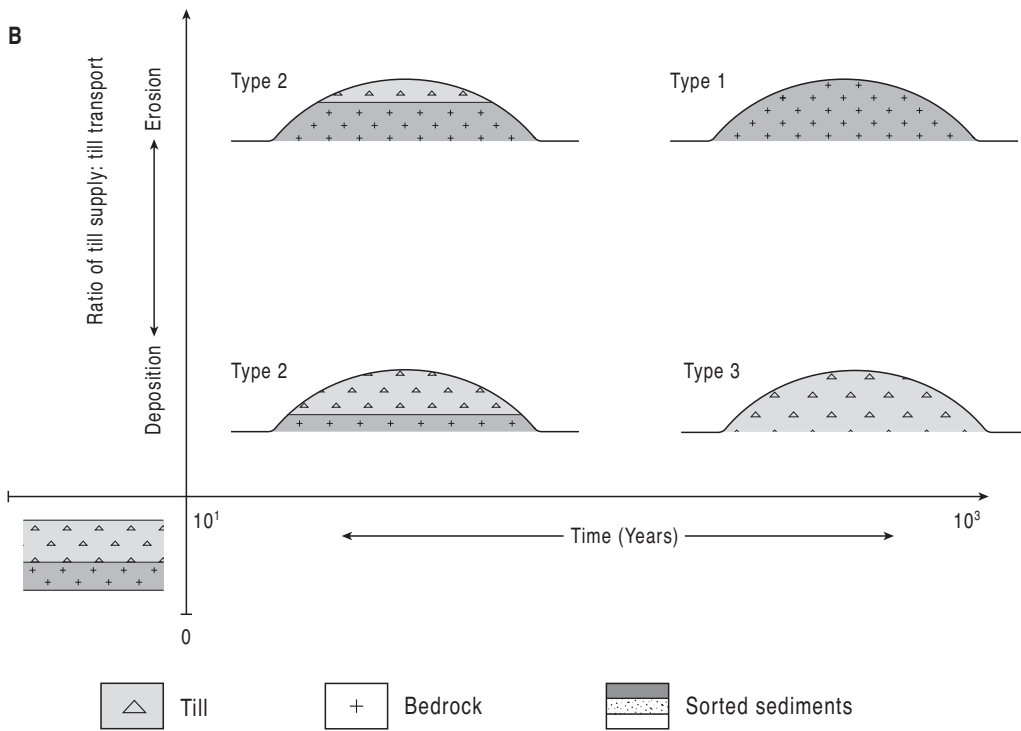
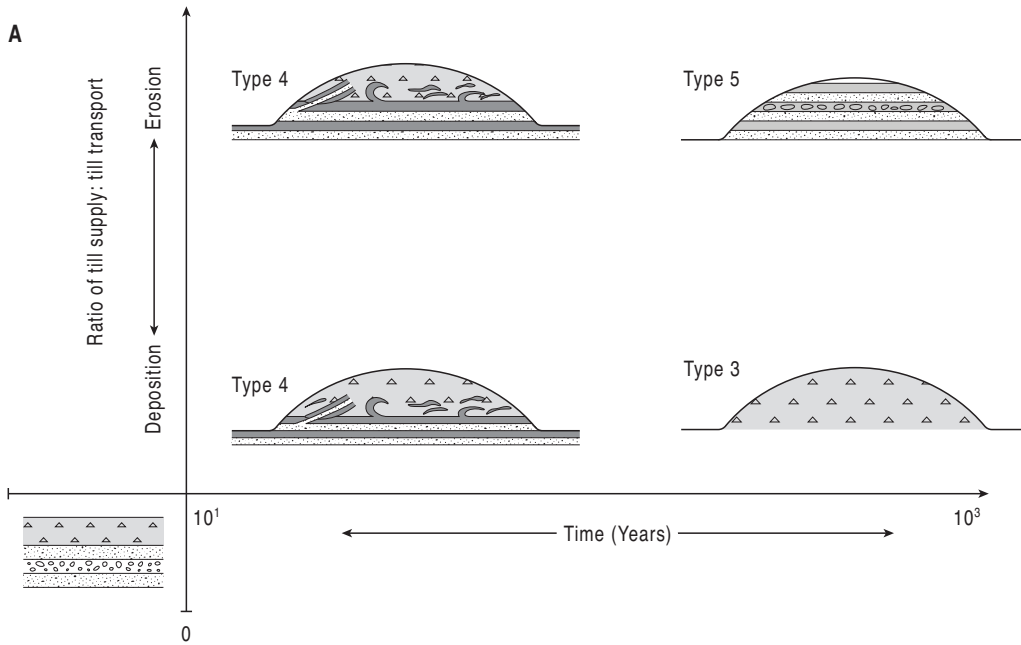
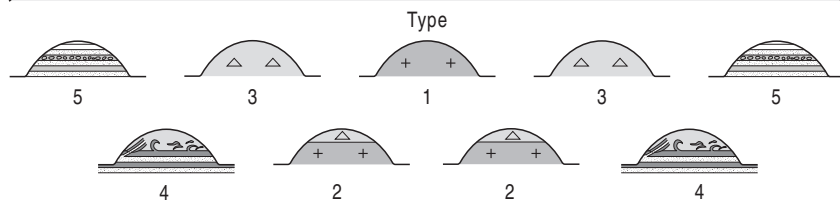
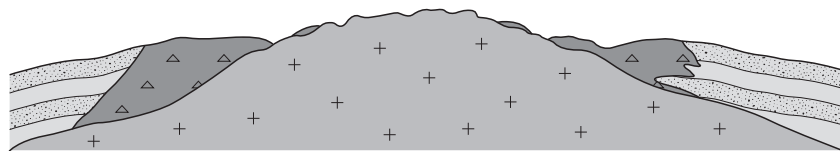
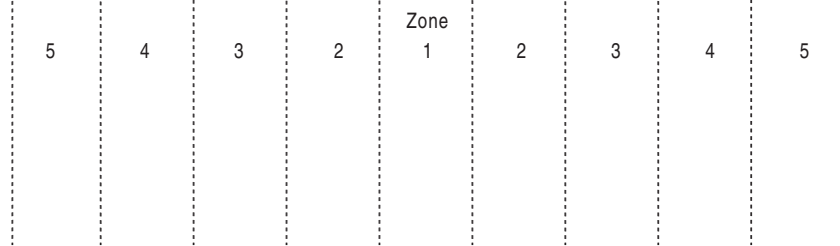
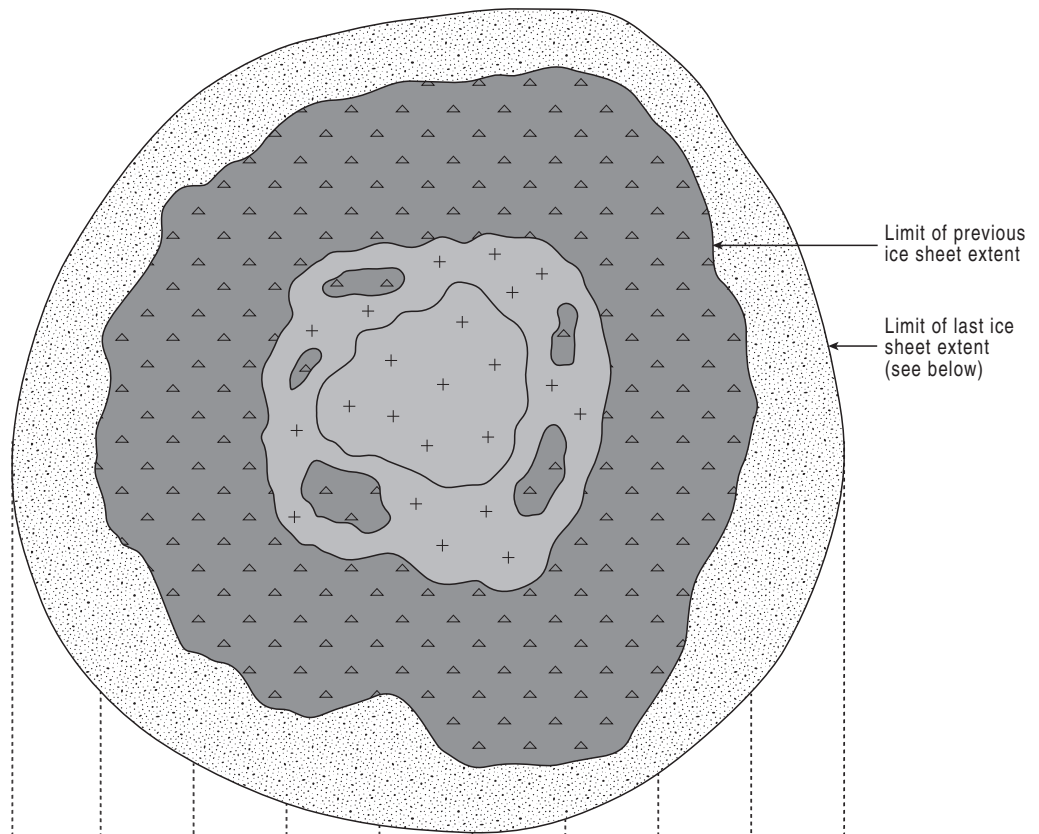


Figure 9





(Stokes *et al.*, 2011)

