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U‐Th‐Ra disequilibria and the extent of off‐axis volcanism across the East Pacific Rise at 9°30′N, 10°30′N, and 11°20′N

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[1] There is widespread interest in the distance that mid‐ocean ridge magmatism extends beyond the neovolcanic zone. Off‐axis magmas also provide a means to map out variations across the melting zone. We present ²³⁸U-²³⁰Th-²²⁶Ra data for 35 well-characterized samples that extend up to 50 km away from the ridge axis across the East Pacific Rise at 9°30′N, 10°30′N, and 11°20′N. The $(^{230}Th/^{238}U)$ ratios range from 1.00 to 1.19, and the $(^{226}Ra^{230}Th)$ ratios range from 1 to 2.78. The samples have a bimodal $(^{230}Th^{238}U)$ distribution with approximately half overlying published axial data on the U‐Th diagram and the remainder lying close to the equiline. The U series disequilibria in the majority of the samples can be explained by aging subsequent to eruption in a zone ∼8 km wide about the neovolcanic zone, consistent with visual evidence for sample age. Nevertheless, seven of the samples lie above calculated $(^{230}Th/^{238}U)$ axial decay curves and/or have ²²⁶Ra excesses implying eruption tens of kilometers off axis. These are consistent with evidence from seamounts and seismic interpretations that magmatism can extend up to 20 km off axis. The implication is that magma is not as efficiently focused beneath the ridge axis as has generally been believed. There is a decrease in initial $(^{230}Th/^{238}U)$ in both these and published samples inferred to have formed off axis, but there is no compelling evidence that this reflects source heterogeneity. Simple modeling suggests that this could be explained by a decrease in fertility and melt column length as the overlying lithosphere thickens with age and the solidus shallows.

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Theme: Geochemical Heterogeneities in Oceanic Island Basalt and Mid-ocean Ridge Basalt Sources: Implications for Melting Processes and Mantle Dynamics

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1. Introduction

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[2] More than 70% of global volcanism consists of basalts erupted at mid‐ocean ridges (MORB) and this is generally thought to be confined to a narrow axial zone [Fornari et al., 2004] and/or efficient magma focusing from a much broader, triangular melting region [e.g., Spiegelman and McKenzie, 1987; Plank and Langmuir, 1992; Korenaga and Kelemen, 1997]. However, the dynamics of melting and melt transport remain a subject of debate. In this regard, U series isotope studies have proved especially useful because of their ability to constrain melting rates and melt extraction porosity (see Elliott and Spiegelman [2003] for a review). One of the most intensively studied of the global ocean ridges is the fast spreading (11.1 cm/yr) East Pacific Rise (EPR) which appears to be underlain by axial magma chambers (see Sinton and Detrick [1992] or Fornari et al. [2004] and references therein for a summary). Although U series isotope data have been extensively analyzed from ∼ zero‐ aged, axial EPR samples, analogous to other settings (e.g., arcs) their interpretation is nonunique and open to question. For example, the effects of source compositional the style of melt extraction (such as dynamic melting versus equilibrium porous flow, versus two-porosity models) and the possibility that disequilibria could be created in the crust instead of the mantle remain controversial [e.g., Lundstrom et al., 1999; Jull et al., 2002; Saal and Van Orman, 2004].

[3] How the composition, eruption temperature and supply rate of MORB may vary as a function of time have been studied using samples (lavas) collected along ridge‐perpendicular transects across the EPR axis [Perfit et al., 1994; Batiza et al., 1996; Niu et al., 1999; Regelous et al., 1999; Reynolds and Langmuir, 2000]. It is important to note, however, that this approach implicitly assumes that all these samples studied were originally formed in the very narrow $\left($ <1–2 km) axial zone. This assumption may be reasonable to a first order but geophysical surveys indicate that the oceanic crust continues to thicken (by a factor of ∼2) over several km away from the ridge axis [e.g., Christeson et al., 1992; Harding et al., 1993; Sohn et al., 2004]. This requires the occurrence of magmatism beyond the axis including lavas flows breaching the axial trough and being deposited on the ridge flanks [e.g., Soule et al., 2005]. This inference has been supported by both submersible observation [e.g., *Perfit et al.*, 1994] and U series dating that has yielded ages younger than the spreading age (i.e., that predicted by combining the distance of the sample from the axis with the local spreading rate) [e.g., Goldstein et al., 1994; Rubin et al., 1994]. For example, in the region 9–10°N, Sims et al. [2003] used U series data to demonstrate the persistence of volcanism up to 4 km from the ridge axis. More strikingly, Zou et al. [2002] presented U-Th isotope evidence for magmatism as far off axis as 20 km at 9°30′N. The latter results were highly controversial, especially the occurrence of several samples with 238U excesses that are not anticipated in MORB [e.g., Elliott, 2002]. Nevertheless, Durant and Toomey [2009] present seismic evidence for the existence of a melt lens 20 km to the east of the ridge axis and numerical modeling that suggests that bending stresses may open tensile fractures which may accommodate off‐axis magmatism out to distances of ∼20 km from the axis [Sohn and Sims, 2005].

[4] In order to explore these substantial matters further, we have undertaken U‐Th‐Ra analyses of samples from traverses across the EPR at $10^{\circ}30'N$ and 11°20N as well as reanalyzing some of the more unusual samples from Zou et al. [2002] at 9°30′N. We have also obtained the first $2\overline{2}$ ⁶Ra data for samples from the three traverses. Our objectives are to (1) see if some of the more unusual results from the Zou et al. [2002] data from 9°30′N are correct, (2) conduct a similar experiment along the two other traverses across the EPR, (3) ascertain the frequency and location of any off‐axis flows, (4) investigate any contrasts in composition and/or melting dynamics between axial and off-axis MORB, and (5) determine whether the inferred temporal changes in eruption temperature and/or composition for the EPR MORB need reappraisal.

2. Sample Locations and Background Data

[5] The samples we analyzed were collected during Phoenix, leg 2 (R/V Melville) by dredging and rock coring along flow lines on both the Cocos and Pacific plates out to 40–50 km (or ∼800 kyr equivalent spreading age and the Brunhes magnetic reversal) from the ridge axis (Figure 1). They come from three traverses conducted across the EPR axis at 9°30′N, 10°30′N and 11°20N that were originally conducted in order to investigate temporal changes in axial magma chemistry over the last 800 kyr [Batiza et al., 1996]. All samples collected from the axis appeared extremely fresh in hand specimen (i.e., black and glassy). Samples collected away from the axis appear progressively older with

Figure 1. Map of the East Pacific Rise using UT [Wessel and Smith, 1991, 1998] showing major transform faults and the three across‐ridge traverses with sample localities indicated.

more alteration and thicker Mn coatings as a function of distance away from the EPR axis. The majority of the samples were dredged from small fault scarps [Batiza et al., 1996].

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[6] *Batiza et al.* [1996] concluded that both 9°30'N and 11°20N are magmatically robust ridge segments whereas 10°30′N appears to be a magmatically starved axis. Some major and trace element data for samples from 9°30′N were presented by Batiza et al. [1996] and U-Th and Pb isotope data by Zou et al. [2002]. Regelous et al. [1999] and Niu et al. [1999] present major and trace element and Sr‐Nd‐Pb isotope data for the samples from 10°30′N and 11°20N, respectively. The samples from 10°30′N are dominated by normal‐ or N‐MORB and show temporal changes in major elements that are not accompanied by systematic changes in trace element or radiogenic isotopes and have been interpreted to reflect changes in eruption temperature and, by inference, magma supply rate through time [Batiza et al., 1996; Regelous et al., 1999]. In contrast, samples from the traverse at 11°20N contain a number of enriched‐ or E‐MORB and exhibit significant variation in radiogenic isotopes interpreted to result from mixing with plume‐like heterogeneities [Niu et al., 1999].

[7] We chose the four $Zou et al.$ [2002] samples from 9°30′N that had the most extreme disequilibria for reanalysis (PH33-3 and PH8-1 with ²³⁰Th excess and PH34-1 and PH10-2 that had 238 U excess), plus a second sample from dredge 8 (sample PH8‐3). From 10°30′N we selected 15 samples extending out 31 km from the axis corresponding to spreading rate ages ranging from 8 to 556 kyr. These included 2 samples from dredge 79 (PH79‐1 and PH79‐2). Along the 11°20N traverse we selected 15 samples extending out to 48 km with an associated range in spreading rate ages ranging from 7 to 873 kyr.

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3. Preparation and Analytical Methods

[8] We used methods employed previously in the Macquarie University U series laboratory for submarine glass samples [*Beier et al.*, 2010]. Suitable glasses were crushed with an agate mortar and pestle to 0.5–1 mm sized chips, washed in deionized water and hand‐picked under a binocular microscope to avoid pieces with visible evidence for crystals, vesicles, hydrothermal alteration and/ or Mn coatings. Selected glass chips were then ultrasonicated in cold 2.5N HCl and 30% H_2O_2 (1:1 mix) for 10–20 min before being washed and ultrasonicated (20 min) in deionized water [e.g., Kokfelt et al., 2005; J. Gill, personal communication, 2005]. The glasses were then reinspected for alteration and hand picked again to obtain a 0.2–1 g quantity of clean glass chips. We analyzed U‐Th in leachates from 3 samples in order to appraise the effects of the cleaning procedure $(^{226}Ra$ concentrations below detection limit).

[9] U, Th and Ra concentrations and isotope ratios were determined on samples and leachates that were spiked with ²³⁶U⁻²²⁹Th and ²²⁸Ra tracers. Glass chips were dissolved using an HF-HNO₃-HCl mix in heated Teflon pressure bombs. All acids were Teflon distilled. The product was converted to chloride using 6N HCl and then 6N HCl saturated with H_3BO_3 to remove residual fluorides. The final product was then converted to nitrate using 14N HNO₃ and finally taken up in 7N HNO₃. U and Th purification was achieved via a single pass through a 4 ml anionic resin column using 7N $HNO₃$, 6N HCl and 0.2N HNO₃ as eluents. We purposefully avoided the use of EiChrom® resins for the U‐Th chemistry as these can bleed organics that lead to memory effects and interferences during MC‐ICP‐MS analysis. Concentrations and isotope ratios were measured in dynamic mode on a Nu Instruments® MC-ICP-MS at Macquarie University.
²³⁸U and ²³⁵U were analyzed on Faraday cups, using the ²³⁸U/²³⁵U ratio to determine the U mass bias, assuming $^{238}U/^{235}U = 137.88$, while ^{236}U and ^{234}U were alternately collected in the IC0 ion counter that is preceded by an energy filter. The IC0 gain was determined during interspersed dynamic analyses of CRM145 assuming a 234 U/ 238 U ratio of 5.286 \times 10⁻⁵ [Cheng et al., 2000]. Methods for Th isotope measurements employed a dynamic routine with 232 Th in Faraday cups and 230 Th and 229 Th alternating on IC0 and using bracketing measurements of

the Th"U" (Open University Th solution) standard to obtain the Th mass bias which can be up to 4% different to that for U [e.g., *Hoffmann et al.*, 2007]. Measurements at masses 230.5 and 229.5 were used to derive a linear correction for residual 232 Th tail interference as described in detail for the GEMOC laboratory in Appendix A of Sims et al. [2008]. For our setup, a polynomial tail correction can lead to a $\leq 0.6\%$ lower ²³⁰Th/²³²Th ratio (but see also discussion below). However, there is no a priori reason to assume that the tail has a polynomial form. Typical within run errors on $234 \text{U}/^{238}\text{U}$ and ²³⁰Th^{$/232$}Th ratio measurements were 2.61 × 10⁻⁸ and 1.10×10^{-8} , respectively.

[10] The Ra analysis procedure follows that used by Turner et al. [2000]. Ra was taken from the first elution from the anionic column and converted to chloride using 6N HCl. This was then loaded in 3 N HCl onto an 8 ml cationic column and Ra eluted using $3.75M HNO₃$ and the process repeated on a scaled‐down 0.6 ml column. The REE were then removed using a 150 μ l column of EiChrom® Ln-spec resin[™] and $0.1N$ HNO₃. Ra and Ba were finally chromatographically separated using EiChrom® Sr-spec resin[™] and 3N HNO₃ as eluent in a 150 μ l procedure. Samples were loaded onto degassed Re filaments using a Ta‐HF‐H3PO4 activator solution [$Birck$, 1986] and ²²⁸Ra/²²⁶Ra ratios were measured to a precision $\leq 1\%$ in dynamic ion counting mode on a ThermoFinnigan Triton® TIMS at Macquarie University. Organic interferences are often noted at low temperatures during TIMS analysis for Ra but were eliminated here by using the instrument fitted with a dry scroll pump instead of the standard rotary pump. This prevents leakage of organic molecules into the source during venting.

[11] In order to be able to demonstrate genuine departures from secular equilibrium in the EPR samples it is important to adopt a conservative estimate for analytical precision. Table 1 contains results for a number of equilibrium rock standards that were analyzed over the same 2 year period (2007–2008) as the traverse samples. The analyses of TML‐3, BCR‐1 and BCR‐2 are generally within error of published values and within 2 standard errors of secular equilibrium for $(^{230}Th/^{238}U)$ and $(^{226}Ra^{230}Th)$, excepting $(^{230}Th^{238}U)$ for BCR-1 that deviates by 4‰ outside of error and $(^{226}Ra^{230}Th)$ for TML-3 that deviates by 2% from equilibrium (brackets indicate activity ratios). On the basis of these results, we adopt 2σ precisions of 8‰ for $(2^{34}U/2^{38}U)$, 2% for $(2^{38}U/2^{32}Th)$, 1% for $(^{230}Th/^{232}Th)$, 2% for $(^{230}Th/^{238}U)$ and 6% for $\tilde{C}^{226}Ra^{230}Th$) for the data.

Standard/Rock	Th (ng/g)	U (ng/g)	$(1^{234}U)^{238}U$	(238) U (232) Th)	$(^{230}Th/^{232}Th)$	$(^{230}Th/^{238}U)$	226 Ra (fg/g)	$(226Ra)^{230}$ Th)
$BCR-1$	5.799	1.683	1.000	0.880	0.883	1.002	570.5	1.00
$BCR-1$	5.789	1.698	1.001	0.890	0.879	0.988		
$BCR-1$	5.552	1.644	1.000	0.898	0.874	0.972	555.5	1.02
$BCR-1a$	5.552	1.623	1.002	0.887	0.874	0.985	555.5	1.02
$BCR-1a$	5.552	1.641	0.999	0.897	0.874	0.974	556.1	1.02
Average	5.649	1.658	1.000	0.891	0.876	0.984	559.4	1.02
2 std. dev.	0.266	0.063	0.002	0.015	0.008	0.024	14.8	0.03
$BCR-2$	5.649	1.665	1.000	0.894	0.879	0.983		
$BCR-2^a$	5.739	1.665	1.000	0.880	0.877	0.996	549.1	0.97
$BCR-2^a$	5.754	1.665	1.000	0.878	0.875	0.996	543.1	0.97
Average	5.714	1.665	1.000	0.884	0.877	0.992	546.1	0.97
2 std. dev.	0.114	n/a	0.000	0.018	0.005	0.015	8.5	0.00
TML-3	29.182	10.186	1.010	1.059	1.089	1.028	3530.1	1.00
$TML-3$	27.504	9.826	1.003	1.084	1.083	0.999	3518.2	1.06
TML-3	29.627	10.582	1.006	1.084	1.077	0.993		
TML-3	29.793	10.476	0.995	1.067	1.075	1.008	3550.0	1.00
TML-3	29.062	10.505	0.993	1.097	1.086	0.990		
Average	29.034	10.315	1.002	1.078	1.082	1.004	3532.8	1.02
2 std. dev.	1.814	0.624	0.015	0.030	0.012	0.030	32.1	0.07
				2392-9 BBQ Flow				
This study	0.117	0.047	0.996	1.217	1.346	1.106	48.9	2.78
Rubin et al. [2005]	0.117	0.048		1.246	1.392	1.117		2.64
Sims et al. [2002]	0.118	0.048	1.002	1.231	1.393	1.131	48.5	2.66

Table 1. U Series Results for Secular Equilibrium Rock Standards and Interlaboratory Samples Over the Period 2007–2009

^aRepeat measurements of U or Th only.

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[12] Table 1 includes results from a MORB glass from the BBQ flow (sample 2392‐9) from the East Pacific Rise [Haymon et al., 1993; Rubin et al., 1994]. The U, Th and 226 Ra concentrations from the BBQ flow are in good agreement with Rubin et al. $[2005]$ and *Sims et al.* $[2002]$ although we obtained a $(^{230}Th/^{238}U)$ ratio that is closer to (and within our analytical reproducibility of) that reported by Rubin et al. [2005] than by Sims et al. [2002]. (Note that there is a typo in the work by Rubin et al. [2005], and sample 2392‐1 is in fact 2392‐9 (K. Rubin, personal communication, 2010).)

[13] The lower 230 Th is the opposite to the effect of seawater alteration or manganese coatings. Because this sample has a very low Th concentration compared with most of the other samples and standards analyzed, it may reflect overcorrection for the 232 Th tail on 230 Th because of the linear extrapolation employed and we also obtained a higher $(^{226}Ra^{230}Th)$ ratio for this sample. However, our λ^{230} Th/²³⁸U) and λ^{226} Ra/²³⁰Th) ratios are not reproducibly lower and higher, respectively, in the other standards analyzed and there may be real (minor) heterogeneity within the BBQ flow. Either way, the consequence of overcorrecting the ²³⁰Th tail would be that the U‐Th disequilibria reported here would represent minimum values; higher $(^{230}Th/^{232}Th)$ ratios would only lead to an increase

in the number of samples inferred to have formed off axis (see below).

4. Results

[14] The reader is referred to Regelous et al. [1999] and Niu et al. [1999] for a detailed description and discussion of the major and trace element and Sr‐Nd‐Pb isotope characteristics and variations in the 10°30′N and 11°20′N traverse samples, respectively. Previously unpublished major and trace element data for the 9°30′N samples are provided in Table 2; U-Th and Pb isotope data for these samples can be found in work by Zou et al. [2002]. As illustrated in Figure 2, the vast majority of the samples are N-MORB basalts; seven are E-MORB. Although we largely restrict the discussion below to the basalts, a few of the samples are more evolved; four are basaltic andesites (PH103‐2, PH46‐3, PH79‐2 and PH78‐2) and one plots in the andesite field (PH79‐1). However, it is unlikely that any of the samples have undergone fractional crystallization of accessory phases, such as apatite, that have the potential to affect U series disequilibria.

[15] The new U-Th-Ra results are presented in Table 3 and because the age of the samples is unknown, no age correction has been applied to any of the data. Seawater has high concentrations

Table 2. Major and Trace Element Data for 9°30'N Samples^a

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Table 2.

(continued)

Figure 2. Plot of $K_2O/TiO_2 \times 100$ versus SiO₂ showing that the majority of the samples are N‐MORB basalts.

of 234U, 230Th and 226Ra that could be inherited in Mn coatings on MORB glasses [Chen et al., 1986]. The leachates have $(23^2U/238U)$ ratios that range from 1.028 to 1.103 and unrealistically high $(^{230}Th/^{232}Th)$ and $(^{230}Th/^{238}U)$ ratios (see Table 3). However, with the exception of 3 of the samples, the $(^{234}U)^{238}U$) ratios of the leached glasses all lie within 2σ analytical error of secular equilibrium (Figure 3) and omitting the 3 outliers (PH10‐2, PH63-1 and PH54-3) does not alter any of the conclusions reached below. Thus, the sample preparation procedures appear to have effectively removed the effects of Mn coatings and/or seawater alteration. Note that Figure 3 also shows that there is no correlation between $(^{234}U)^{238}U$ and $(^{230}Th/^{238}U)$. Duplicate samples from the same dredges (i.e., PH8‐1 and PH8‐3, PH79‐1 and PH79‐2) and a duplicate analysis of the same sample (PH33-3) yielded $(^{230}Th/^{238}U)$ ratios that were reproducible within the precision estimates adopted above (Table 3). For clarity we begin by discussing the U‐Th‐Ra results from each traverse separately.

4.1. Samples at 9°30′N

[16] We obtained contrasting results to those of Zou et al. [2002] for the four most anomalous samples in their study. Our U and Th concentrations are in good agreement with the exception of sample PH34‐1 that has a higher concentration of Th than those reported by Zou et al. [2002]. Compared to the Zou et al. [2002] data, sample PH33‐3 has lower U/Th and $(2^{30} \text{Th})^{232}$ Th) ratios and a 1.5% measured ²³⁰Th excess. This is within error of secular equilibrium but this sample also has 25% ²²⁶Ra excess suggesting that sample represents an

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Figure 3. Plot of $(^{230}Th/^{238}U)$ versus $(^{234}U/^{238}U)$ showing that the vast majority of the samples have $(234U/238U)$ within 2σ analytical error (8‰) of secular equilibrium (gray field) and that there is no correlation with the extent of U‐Th disequilibria. Dashed line indicates secular equilibrium.

off‐axis eruption (see below). Sample PH8‐1 is in U‐Th isotopic equilibrium whereas sample PH8‐3, from the same dredge, is also essentially indistinguishable from secular equilibrium. We suggest that both of these latter two samples are effectively in secular equilibrium (and thus could have formed at the ridge axis). This is in contrast to the 13% 230 Th excess reported by Zou et al. [2002] for PH8‐1. The biggest difference between the two studies is that samples PH34‐1 and PH10‐2, that had significant 238 U excesses in the previous work, are either within error of secular equilibrium (sample PH34-1) or have a small but significant (4%) ²³⁰Th excess (sample PH10‐2). On the U‐Th equiline diagram, three of the samples lie close to the equiline at high U/Th whereas the remaining two lie at low U/Th, one with $230T$ Th excess (Figure 4a).

4.2. Samples at 10°30′N

[17] The samples from this traverse have Th and U concentrations that range from 0.16 to 1.55 and 0.06 to 0.57 ppm, respectively, which are in good agreement with ICP‐MS concentration data presented by Regelous et al. [1999]. ²²⁶Ra concentrations range from 33 to 98 fg/g overlapping the range observed by Lundstrom et al. [1999] and Sims et al. [2002, 2003]. $(^{230} \text{Th}/^{232} \text{Th})$ ranges from 1.12 to 1.34 and the majority of the samples have ²³⁰Th excess with $(^{230}\text{Th}/^{238}\text{U})$ ranging from within error of secular equilibrium up to 17%. Of the samples successfully analyzed for $226Ra$ (some TIMS analysis runs failed), four are within analytical error of secular equilibrium while the remaining four have significant ²²⁶Ra excesses (PH52-2, PH70-2, PH62-1 and PH63-1) with $(^{226}Ra)^{230}Th$ ratios that varies from 1.5 to 2.2.

[18] On the U-Th equiline diagram, three of the 10°30′N samples lie parallel to, and overlap, the

Figure 4. (a) U-Th equiline diagram, (b) $(^{226}Ra^{/230}Th)$ versus $(^{230}Th)^{238}U$) diagram, and (c) an expansion of the lower part of the $(^{226}Ra^{230}Th)$ versus $(^{230}Th^{238}U)$ diagram (error bars are 2σ). Also shown is the field for axial MORB samples from the region 9°–10°N on the EPR (data from Lundstrom et al. [1999] and Sims et al. [2002]). Dashed lines indicate secular equilibrium.

field for 9°30′–50′N axial samples [Lundstrom et al., 1999; Sims et al., 2002], nine lie close to the equiline and three plot in between these two groups (Figure 4a). Note that this field encompasses N‐, E‐ and D‐MORBs each with their own source chemistry and conditions of petrogenesis. On a plot of $(^{226}Ra^{/230}Th)$ versus $(^{230}Th^{/238}U)$, two samples lie within the negatively sloped EPR array of Sims et al. [2002]. The remainder fall below this array and/or within error of secular equilibrium. This is consistent with decay since the true age of the samples is unknown (see below).

4.3. Samples at 11°20′N

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[19] Samples from the 11°20′N traverse show slightly more restricted Th and U concentrations that range from 0.13 to 0.98 and 0.05 to 0.37 ppm, respectively, but again agree well with ICP‐MS concentration data from Niu et al. [1999]. 226 Ra concentrations range from 22 to 98 fg/g, again overlapping that observed in data from Sims et al. $[2002, 2003]$. $(^{230} \text{Th}/^{232} \text{Th})$ ranges from 1.12 to 1.38 and two samples have $(^{230} \text{Th}/^{238} \text{U})$ less than, but within error of, 1. The majority of the remaining samples have ²³⁰Th excess with $(230 \text{Th}/238 \text{U})$ ratios ranging up to 1.19. Samples PH82‐1, PH89‐1, PH100-1 and PH83-1 all have $(^{226}Ra^{230}Th)$ ratios that are effectively in secular equilibrium. In the remaining six samples $(^{226}Ra^{230}Th)$ ratios range from 1.09 to 2.13 (PH88‐1, PH92‐1, PH91‐1, PH101‐1, PH104‐1 and PH102‐1).

[20] On the U-Th equiline diagram, the $11^{\circ}20'N$ samples form a bimodal distribution and overlap the two main groups from 10°30′N samples. Within these two arrays $(^{230} \text{Th}/^{232} \text{Th})$ is strongly correlated with $(^{238}U/^{232}Th)$, commonly ascribed to source composition [e.g., Goldstein et al., 1992]. Seven samples fall within or close to the EPR 9°30′–50′N axial array whereas the remainder scatter around the equiline (Figure 4a). On the $(^{226}Ra^{230}Th)$ versus $(^{230}Th/^{238}U)$ plot, the 11°20′N samples again overlap the 10°30′N samples and two samples lie within the Sims et al. [2002] EPR array, indicating that their $(^{226}$ Ra $/^{230}$ Th) ratios may be near primary. Again, the remainder fall below this and/or within error of secular equilibrium (Figures 4b and 4c) suggestive of decay since eruption.

5. Off‐Axis Eruptions

[21] In Figures 5 and 6, we plot the U series disequilibria against distance from the ridge axis for all three traverses (error bars are 2σ). The half-life

of ²³⁰Th is 75 kyr and so the resolution of the ²³⁰Th-²³⁸U system is, in principle, ~600 kyr, corresponding to a distance of ∼33 km for the EPR half-spreading rate. The U-Th data are compared with two decay curves that assume a half-spreading rate of 5.5 cm/yr and initial, ridge-formed $(^{230}Th/^{238}U)$ ratios of 1.19 and 1.11 that correspond to the highest and lowest values, respectively, reported by Sims et al. [2002] for 9°30′–50′N axial EPR samples. While these data of are clearly pertinent to the 9°30′N traverse, they might be considered less so for the other two traverses further north. However, U‐Th data extending as far north as 10°N are consistent with the maximum axial $(^{230}Th/^{238}U)$ ratios being ~1.17–1.19 [Lundstrom et al., 1999]. Moreover, the initial $(^{230}\text{Th}/^{238}\text{U})$ would have to be >2 in order for all of the samples to lie on or below the axial decay curves and this value significantly exceeds any MORB measurement worldwide [Elliott and Spiegelman, 2003].

[22] *Zou et al.* [2002] did not undertake 226 Ra measurements in their study, but Figures 5d–5f show plots of Ra‐Th disequilibria against distance from the ridge axis for all three traverses along with a decay curve based on the highest $(^{226}Ra^{230}Th)$ ratio (4.2) measured by Lundstrom et al. [1999]. Note that the short (1.6 kyr) half-life of ^{226}Ra means that any sample formed at the ridge axis will have returned to 2^{26} Ra -2^{30} Th equilibrium within 10 kyr, or a distance of only 0.5 km for the EPR spreading rate. Alteration might lead to 226 Ra excesses but the samples have $({}^{234}U/{}^{238}U)$ within error of secular equilibrium and excessive leaching would be unlikely to produce $({}^{226}Ra/{}^{230}Th)$ > 1. It has also been suggested that diffusive interaction between melt and cumulate plagioclase could lead to 226Ra excesses in MORB [Saal and Van Orman, 2004]. However, that would not significantly alter the time scale implications made below. Note also that, because the age of the samples is unknown, the measured $(^{226}Ra^{/230}Th)$ ratios should be regarded as minima. Samples with ²²⁶Ra excess beyond analytical uncertainty are identified in Figures 5a–5c by gray filled symbols.

[23] Finally, Sims et al. [2003] have demonstrated that eruptions can occur up to 4 km off axis at $9^{\circ}50'N$ on the EPR consistent with the suggestion of Perfit and Chadwick [1998]. Indeed, the last eruption in this area flowed $2+$ km away from the ridge axis and erupted in places from vents that were almost a km off axis [Soule et al., 2007]. Thus, we also show a set of dashed gray curves in Figure 5 which permit that the samples were emplaced up to 4 km off axis with an initial $(^{230}Th/^{238}U)$ ratio of 1.19 and an initial $(^{226}Ra/^{230}Th)$ ratio of 4.2. This probably

Figure 5. Plots of $(^{230}Th)^{238}$ U) and $(^{226}Ra)^{230}$ Th) versus distance either side of the ridge axis for the three traverses (error bars are 2σ). Solid curves show decay of axially formed disequilibria, assuming a half-spreading rate of 5.5 cm/yr. For $(^{230}Th/^{238}U)$, these assumed maximum and minimum initial ratios of 1.19 and 1.11 as measured in axially formed samples from $9^{\circ}30'$ –50′N [Sims *et al.*, 2002]. The decay curve for $(^{226}Ra^{230}Th)$ assumed an initial ratio of 4.2 as measured on the A-B transform [Lundstrom et al., 1999]. The dashed curves reflect the same model but allow for the lavas to have formed at the edge of the neovolcanic zone that is thought to be up to 8–10 km wide [e.g., *Perfit et al.*, 1994; Sims et al., 2003]. Samples lying above the decay curves are inferred to have formed off axis, and those samples with significant ²²⁶Ra excess are indicated by gray symbols on the $(^{230}Th/^{238}U)$ diagrams. The shaded bar on the 9°30′N traverse indicates the location of the melt lens identified seismically by Durant and Toomey [2009].

represents an extreme case, since most of the offaxis samples analyzed by Sims et al. [2003] had somewhat lower $({}^{230} \text{Th}/{}^{238}\text{U})$ and $({}^{226} \text{Ra}/{}^{230}\text{Th})$ ratios than their axial samples. Nevertheless, these curves allow us to distinguish any samples that erupted >4 km from the axis.

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[24] In the case of samples for which the data overlap, within error, or lie beneath the decay curves there is no reason, from a U series isotope perspective, to suspect that they did not erupt at the ridge axis and become subsequently transported outward with the plate. On this basis, only one of the four 9°30′N samples (PH10‐2) lie above the U‐Th decay curves but that sample could, within error, have formed 4 km off axis (gray dashed curve in Figure 5a). Problematically, this is one of

Figure 6. Expanded view of the near-axis region for the $10^{\circ}30'N$ and $11^{\circ}20'N$ traverses.

the samples with a measured $(^{234}U/^{238}U)$ outside of analytical error from equilibrium. However, one sample ∼20 km to the west of the ridge (PH33-3) which lies above, but overlaps within error, the U-Th decay curves also has significant 2^{26} Ra excess indicating that this sample formed off axis close to its present location. Note that the combination of large ²²⁶Ra excess with low $(^{230}_{\sim} \text{Th}/^{238} \text{U})$ is consistent with the negative $(^{226}Ra)^{230}Th$ – $(^{230}Th)^{238}U$) array found in EPR samples [Sims et al., 2002].

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[25] At $10^{\circ}30^{\prime}$ N, two samples $10-20$ km to the east of the ridge clearly lie outside of error above all three of the U‐Th decay curves while the remaining thirteen lie within error of them. One of these samples (PH52-2) has significant 226 Ra excess, as does one sample ∼10 km to the west of the ridge (PH70‐2) that lies within error of the decay curves. Looking an expanded view close to the ridge axis (Figures 6a and 6b), all samples overlap the U‐Th decay curves yet two lie above the ²²⁶Ra decay curve suggesting eruption a few km outside of the neovolcanic zone.

[26] Finally at 11°20′N, one sample to the far east of the ridge (PH83‐1) lies just above the three U‐Th

decay curves while the remaining fourteen lie within error of them. However, a further three of the eastern samples have significant 226 Ra excess, again implying off‐axis eruption. Close to the ridge axis, one sample (PH108‐1) lies above the two axial U‐Th decay curves but could still have formed within 4 km of the axis while a further three samples have 226 Ra excesses suggesting that they must also have formed a few km outside of the neovolcanic zone (Figures 6c and 6d).

[27] In summary, our new data do not confirm the Zou et al. [2002] finding of *large* (i.e., axial magnitude) 230 Th excesses 10–20 km either side of the ridge at $9^{\circ}30'N$. Furthermore, no significant ^{238}U excesses have been found on any of the traverses [cf. Zou et al., 2002]. Close to the axis, our data support suggestions [e.g., Perfit and Chadwick, 1998; Sims et al., 2003; Soule et al., 2005] that volcanism extends several km beyond the neovolcanic zone (Figure 6). If this is taken into account in our models (i.e., the dashed gray curves in Figures 5 and 6), then 80% of the samples formed in the neovolcanic zone and aged with the plate as it migrated outward. These data are consistent with the observed alteration state of the samples.

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[28] Nevertheless, our data also permit that seven of the samples formed more than 4 km from the axis. Three of these have small to moderate 230 Th excesses and could have erupted within ∼10 km of their present location. However, 5 samples have 226 Ra excesses suggesting that they erupted within ∼0.5 km of their present location. This reinforces the conclusion of Zou et al. [2002] that MORB eruptions are not restricted to within a few km either side of the ridge axis [e.g., Goldstein et al., 1994; Rubin et al., 1994; Sims et al., 2003] but can occur as far out as 20 km through crust that is up to 550 kyr old (see Figure 5). The majority of the samples were dredged from small fault scarps [*Batiza et al.*, 1996] that could reflect the tensile fractures predicted by Sohn and Sims [2005] and some of these may have provided pathways for magma extrusion. It is also possible that some of the samples not analyzed for 226 Ra, but lying on the axis U‐Th decay curves, were also erupted off axis. These observations challenge the hypothesis that accretion of the crust is complete within a few km of the ridge axis and models that invoke highly efficient melt focusing [e.g., Spiegelman and McKenzie, 1987; Korenaga and Kelemen, 1997; Fornari et al., 2004]. They also imply that not all of the samples from the traverses record the temporal evolution of MORB magmatism at the axis [cf. Batiza et al., 1996; Regelous et al., 1999; Niu et al., 1999] (possibly explaining some of the scatter in the published trends) but may instead contain important information about melt generation beyond the axial zone [Niu et al., 1996].

[29] Before we explore some of the further implications of the potential off‐axis eruptions we recall that there is both seismic data [Garmany, 1989] and active seamounts on >1 Myr crust [e.g., Fornari et al., 1988; Batiza et al., 1990; Niu and Batiza, 1997; Niu et al., 1996] that provide compelling evidence for active magmatism up to 20–40 km from the ridge axis. In particular, Durant and Toomey [2009] found seismic evidence for a 5 km wide melt body 2 km deep beneath the Cocos plate 20 km to the east of the EPR axis at 9°30′N (see Figure 5a). Figure 5 suggests that off‐axis volcanism occurs in roughly this location on the 10°30′N traverse as well. Importantly, our data also suggest off‐axis eruption to the west, on the Pacific plate (see Figure 5), and this is consistent with the observations of significant U‐Th disequilibria in samples from the Lamont seamounts [*Lundstrom et al.*, 1999]. However, with the present coverage there appears to be no simple pattern to the distribution of the few inferred off‐axis eruptions.

6. Origin of Low $(^{230}Th/^{238}U)$ Ratios in Off‐Axis Samples

[30] One aspect of our data is that the subset of samples inferred to have erupted off axis may afford the possibility to interrogate melting dynamics beneath the EPR over a wider across‐axis area that had previously been thought possible. In this respect, one of the most intriguing issues with the observations is that even those off‐axis samples with significant ²²⁶Ra excess still have lower $(^{230}Th/^{238}U)$ than MORB erupted at the ridge axis (for these samples, the measured $(^{230}Th/^{238}U)$ can be considered primary). A similar observation is apparent in other data sets although its significance has not been discussed. Off-axis samples from the Southwest Indian Ridge have lower $(230 \text{Th}/238 \text{U})$ and $(^{226}Ra/^{230}Th)$ than axial samples [Standish and Sims, 2010], and Lundstrom et al. [1999] found that $(^{230}Th)^{238}$ U) ratios decreased with distance from the EPR ridge axis across the Lamont seamounts. For example, *Lundstrom et al.* [1999] reported an analysis from the Siqueiros Transform with $(^{230}Th/^{238}U)$ of 1.01 in a sample that had $(^{226}Ra)^{230}$ Th) > 3. Even EPR samples taken within 4 km of the ridge axis exhibit less U‐Th‐Ra disequilibria than those formed at the ridge axis [Sims et al., 2003]. This suggests that the off-axis samples were derived from magmas that had lower initial $(^{230}Th/^{238}U)$ and we note in passing that this may compromise age determinations that assume a generic axial $(^{230}\text{Th}/^{238}\text{U})$ value. Possible explanations for lower initial $({}^{230} \text{Th}/^{238} \text{U})$ ratios off axis include the role of source heterogeneity and changes in melting conditions.

6.1. The Role of Source Heterogeneity

[31] Numerous studies have argued that there is evidence for heterogeneities in the MORB source and some have argued that this could provide an important control on U series disequilibria [e.g., Lundstrom et al., 1998; Russo et al., 2009]. Indeed, Davies [2009] has recently argued for intimate mixing between melts from both depleted and enriched sources beneath all ridges and inferred that the enriched components are sourced from ancient recycled mafic material (e.g., pyroxenite [cf. Sobolev et al., 2007]). Perfit et al. [1994] noted the intimate juxtaposition of N‐MORB and E‐MORB across the EPR, and some of our samples do have E‐MORB compositions (cf. Figure 2). Furthermore, the samples from 11°20′N show a significant range in radiogenic isotope composition demanding source heterogeneity [Niu et al., 1999].

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[32] Experiments have shown that pyroxenite has a higher melt productivity (%/GPa decompression) than peridotite [Petermann and Hirschmann, 2003]. U‐Th disequilibria are highly sensitive to melting rate (see discussion below) and the faster melting rate of pyroxenite should result in lowered ²³⁰Th excesses [Prytulak and Elliott, 2009; Russo et al., 2009], all else being equal. A caveat to this is that Elkins et al. [2008] have suggested that the higher partition coefficient for U relative to Th in pyroxenite would offset the effects of the faster melting rate. Nevertheless, in principle, the presence of pyroxenite in the melting region could explain the lower $(^{230}Th/^{238}U)$ ratios in off-axis samples, although this would require that pyroxenite is fortuitously present off axis but not beneath the neovolcanic zone [cf. Zou et al., 2002]. How-

ever, we did not find any correlation between the U series disequilibria and radiogenic isotopes. Furthermore, the P/Nd ratios of the EPR traverse samples average 62 ± 9 , apparently precluding a significant role for pyroxenite [O'Neill and Mallmann, 2007]. We conclude that there is no convincing evidence that the lowered 230Th excesses in the inferred off-axis samples reflect source heterogeneity [cf. Zou et al., 2002]. Therefore, it is likely that any putative recycled components (as indicated by trace element or radiogenic isotope data) are present in the form of peridotite that has been metasomatized by pyroxenite melts rather than pyroxenite itself [e.g., Niu and O'Hara, 2008; Elkins et al., 2008; Prytulak and Elliott, 2009].

6.2. Variations in Melting Conditions

[33] An alternative explanation for the decrease in primary $(^{230}Th/^{238}U)$ ratios away from the ridge axis is a change in the melt regime between axial and off‐axis melt production. Models for the generation of U series disequilibria in MORB have

Figure 7. (a and b) Illustration of possible melting models to explain decreased, off-axis $(^{230}Th/^{238}U)$ ratios using the dynamic melting model of Williams and Gill [1989] and input parameters given in Table 4. In the first model, shown as the short‐dashed curve in Figure 7b, the length of the melting column decreases away from the axis solely as a result of thickening of the overlying plate. The oceanic crust is assumed to thicken by a factor of 1.5 away from the ridge axis, and the thickness of the lithosphere (age is indicated in italics) was calculated following Parsons and McKenzie [1978] and Turcotte and Schubert [2002]. Note that the melting region is required to dip into the garnet peridotite zone on the basis of the large 230 Th excesses observed [see also *Sims* et al., 2002], and the location of the spinel to garnet transition is based on Robinson and Wood [1998]. In the second model, shown as the long‐dashed curve, the length of the melting column additionally shortens away from the axis because it is also assumed that the effects of melt extraction in the central upwelling region results in cooler, more depleted peridotite moving away from the axis such that the depth of intersection with the solidus shallows (see text for details). This is inferred to result in a decreased clinopyroxene mode (assumed mineralogy is shown at the base of the illustration). Solid mantle flow lines are shown schematically in Figure 7a, and melt flow is assumed to be vertical. The location of the melt lens identified by Durant and Toomey [2009] and the Lamont seamounts on the Pacific plate are also shown. Samples plotted in Figure 7b are those inferred to have erupted off axis (cf. Figures 5 and 6) supplemented by the Lamont seamounts data from Lundstrom et al. [1999]. Ship not to scale.

Parameter	Units		
Age $(kyr)^a$	$0 - 900$		
Total melt extent $X(\%)$	0.1		
Porosity ϕ (%)	0.001		
ρs (kg/m ³)	3300		
ρf (kg/m ³)	2800		
Melting column top $(km)^{6}$	$6 - 18$		
Melting column bottom (km)	$100 - 60$		
Upwelling rate W_s (cm/yr)	5.5		
Average pressure (Gpa)	$1.5 - 2.2$		
Partition Coefficients ^c and Mode	U	Th	Ra
Olivine $(57%)$	0.00006	9.52E-06	5.75E-08
Orthopyroxene $(28-34\%)$	0.0030	0.0017	$6.00E-07$
Clinopyroxene $(7-15%)$	$0.0119 - 0.0121$	$0.0101 - 0.0108$	$4.13E-06$
Garnet $(8-0\%)$	0.015	0.002	7.00E-09

Table 4. Input Parameters for Dynamic Melting Model

^aCalculated assuming a half-spreading rate of 5.5 cm/yr.

^bCalculated assuming $z = 2.32 \cdot \sqrt{(k \cdot t)}$ as discussed in the text.

From Blundy and Wood [2003], Salters et al. [2002], and Adam and Green [2006]; clinopyroxene U-Th partitioning was calculated following Landwehr et al. [2001] for the average pressure of the melt columns.

recently been reviewed by Elliott and Spiegelman [2003] to which the reader is referred for a detailed discussion. Key parameters are the mineral melt partition coefficients [cf. Wood et al., 1999; Landwehr et al., 2001], the porosity of the melting region (ϕ) , the length of the melt column and the melting rate (Γ) . The latter is linked to the upwelling rate (W_s) that will, in turn, equate to the halfspreading rate directly beneath the ridge axis. Of these, it is decreasing melt column length (which reduces the time available for daughter ingrowth) that seems the most likely candidate to explain the off-axis decrease in initial $(^{230}Th/^{238}U)$ that we now explore further.

[34] A model illustrating the effects of decreasing melt column length, due to the thickening of the overlying plate, on U‐Th disequilibria is illustrated in Figure 7a. The thickness of the lithosphere $(z₁)$ and increase in bathymetry away from the ridge axis is proportional to its age [e.g., Parsons and McKenzie, 1978; Turcotte and Schubert, 2002]:

$$
z_l = 2.32 \cdot \sqrt{(k \cdot t)},
$$

where k is the thermal conductivity, $\sim 2^{-6}$ m²/s [*Chai et al.*, 1996], and t is time. Accordingly, the lithospheric thickness increases from near zero at the ridge axis to ∼18 km at age 800 kyr reducing the length of putative melt columns. In Figure 7b we show the results of a dynamic melting model, based on Williams and Gill [1989], to simulate this scenario. At the axis, the magnitude of measured ²³⁰Th excesses indicate that the axial melt region reaches into the garnet stability zone [cf. Sims et al., 2002] which is located at ∼83 km, based on the experiments of Robinson and Wood [1998]. The porosity was set at 10^{-3} to achieve or exceed the minimum $(^{226}Ra^{/230}Th)$ ratios observed in the data (resultant model $(^{226}Ra)^{230}Th$) values were 2.2–2.9). Other parameters are given in Table 4. Our calculations show that $(^{230} \text{Th}/^{238} \text{U})$ is indeed predicted to decrease with increasing distance away from the ridge axis but that the effect is insufficient to simulate the data (Figure 7b).

[35] Depending on the geometry of the matrix flow lines, off‐axis regions experience a decrease in the depth of the solidus as melt extraction leads to the source region becoming cooler and possibly more depleted or refractory [e.g., Reynolds and Langmuir, 2000]. The latter would likely result in an absence of garnet and decreasing clinopyroxene mode in the residue such that Th becomes less incompatible relative to U (see Figure 7a and Table 4). Thus, Figure 7b shows the results of a second model in which the depth of onset of melting shallows, and fertility decreases away from the axis (see Figure 7a). We assumed that the solidus for depleted peridotite occurs at ∼2 GPa at a temperature of ∼1380°C [Jaques and Green, 1980]. This model is highly simplistic but, nevertheless, provides a better simulation of the off-axis decrease in $(^{230}Th/^{238}U)$. Note that we kept the total melt fraction and the upwelling rate constant and it is important to recognize that the probable decrease in both of these away from the axis would lead to higher 230 Th excesses that would worsen the approximation of the data. Clearly, even relatively modest changes in porosity, extent of

melting and partition coefficients would have significant affects on the results and one could envisage a family of models with broadly similar outcomes. However, fundamental limitations of the data set do not warrant more sophisticated calculations (equilibrium transport, two‐porosity, etc.) that must await further constraints on the numerous input assumptions (e.g., more detailed sampling, 231 Pa measurements to better constrain primary disequilibria, porosity etc). Mixing may also be important [e.g., Sims et al., 2002].

7. Summary

[36] U-Th-Ra analyses of samples from traverses extending up to 50 km either side of the East Pacific Rise at 9°30N, 10°30N and 11°20N are consistent with the bulk of samples having formed in an axial region that is ∼8 km wide. Although the large ²³⁰Th and ²³⁸U excesses reported by Zou et al. [2002] have not been substantiated we do find evidence for some anomalously young samples tens of kilometers from the axis. These are consistent with recent seismic interpretation of a magma lens at this distance, at least to the east of the ridge [Durant and Toomey, 2009].

[37] The samples inferred to have erupted off axis have low initial $(^{230}Th/^{238}U)$ compared to axial samples, both here and in other studies [e.g., *Sims*] et al., 2003; Standish and Sims, 2010]. There is little evidence that this reflects source heterogeneity. Decreasing melting column length and shallowing of the solidus could, in principle, explain the decreases in primary $(^{230}Th/^{238}U)$ ratios away from the ridge axis. The alternative, that the lower $(230 \text{Th}/238 \text{U})$ ratios reflect aging, seems precluded by the 226Ra data whereas the formed interpretation is supported by the observed decrease in $(^{230}Th/^{238}U)$ across the Lamont seamounts to the west of the EPR axis.

[38] The most profound conclusion is that magma focusing beneath the ridge axis may not be as efficient as widely believed, despite the earlier recognition of active off‐axis seamount volcanism [*Batiza et al.*, 1990]. This may have significant implications for recent chromatographic melt migration models [e.g., Spiegelman et al., 2001]. Future work should be aimed at increasing the spatial resolution of across‐axis data, include additional seamount analyses and 231 Pa measurements to further test the age implications and to constrain the porosity in the melt region.

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