



The SCUBA-2 Cosmology Legacy Survey: the submillimetre properties of Lyman-break galaxies at $z = 3\text{--}5$

K. E. K. Coppin,^{1★} J. E. Geach,¹ O. Almaini,² V. Arumugam,^{3,4} J. S. Dunlop,³ W. G. Hartley,⁵ R. J. Ivison,^{3,4} C. J. Simpson,⁶ D. J. B. Smith,¹ A. M. Swinbank,⁷ A. W. Blain,⁸ N. Bourne,³ M. Bremer,⁹ C. Conselice,² C. M. Harrison,⁷ A. Mortlock,³ S. C. Chapman,¹⁰ L. J. M. Davies,¹¹ D. Farrah,¹² A. Gibb,¹³ T. Jenness,^{14,15} A. Karim,¹⁶ K. K. Knudsen,¹⁷ E. Ibar,¹⁸ M. J. Michałowski,³ J. A. Peacock,³ D. Rigopoulou,^{19,20} E. I. Robson,^{3,21} D. Scott,¹³ J. Stevens¹ and P. P. van der Werf²²

Affiliations are listed at the end of the paper

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ABSTRACT

We present detections at 850 μm of the Lyman-break galaxy (LBG) population at $z \approx 3, 4$, and 5 using data from the Submillimetre Common User Bolometer Array 2 Cosmology Legacy Survey in the United Kingdom Infrared Deep Sky Survey ‘Ultra Deep Survey’ field. We employ stacking to probe beneath the survey limit, measuring the average 850 μm flux density of LBGs at $z \approx 3, 4$, and 5 with typical ultraviolet luminosities of $L_{1700} \approx 10^{29} \text{ erg s}^{-1} \text{ Hz}^{-1}$. We measure 850 μm flux densities of (0.25 ± 0.03) , (0.41 ± 0.06) , and $(0.88 \pm 0.23) \text{ mJy}$, respectively, finding that they contribute at most 20 per cent to the cosmic far-infrared (IR) background at 850 μm . Fitting an appropriate range of spectral energy distributions to the $z \approx 3, 4$, and 5 LBG stacked 24–850 μm fluxes, we derive IR luminosities of $L_{8-1000 \mu\text{m}} \approx 3.2, 5.5$, and $11.0 \times 10^{11} L_{\odot}$ [and star formation rates (SFRs) of $\approx 50\text{--}200 M_{\odot} \text{ yr}^{-1}$], respectively. We find that the evolution in the IR luminosity density of LBGs is broadly consistent with model predictions for the expected contribution of luminous-to-ultraluminous IR galaxies at these epochs. We observe a positive correlation between stellar mass and IR luminosity and confirm that, for a fixed mass, the reddest LBGs (UV slope $\beta \rightarrow 0$) are redder due to dust extinction, with $\text{SFR(IR)}/\text{SFR(UV)}$ increasing by about an order of magnitude over $-2 < \beta < 0$ with $\text{SFR(IR)}/\text{SFR(UV)} \sim 20$ for the reddest LBGs. Furthermore, the most massive LBGs tend to have higher obscured-to-unobscured ratios, hinting at a variation in the obscuration properties across the mass range.

Key words: galaxies: evolution – galaxies: formation – galaxies: high-redshift – galaxies: star formation – submillimetre: galaxies – ultraviolet: galaxies.

1 INTRODUCTION

Lyman-break galaxies (LBGs) are currently the largest population of star-forming galaxies known to be at high redshift, $z > 3$, and as such have provided valuable insights into the mass assembly of galaxies during the first few Gyr of the Universe. From a practical standpoint, this has been due, in part, to the efficiency of the simple selection of LBGs in broad-band colours that span the eponymous ‘Lyman break’ as it is redshifted through the optical bands at $z > 3$ (Steidel & Hamilton 1993). This has made it possible to identify

large samples of LBGs, and from this we have learned that they are massive ($M_{\star} \sim 10^9\text{--}10^{11} M_{\odot}$; Sawicki & Yee 1998; Reddy et al. 2006; Rigopoulou et al. 2006; Verma et al. 2007; Magdis et al. 2008, 2010a; Stark et al. 2009), rapidly star-forming (tens to hundreds of $M_{\odot} \text{ yr}^{-1}$; e.g. Sawicki & Yee 1998; Shapley et al. 2001, 2005; Chapman & Casey 2009; Magdis et al. 2010b; Rigopoulou et al. 2010) and numerous ($\phi^* \sim 0.005 \text{ Mpc}^{-3}$; e.g. Reddy & Steidel 2009). As such, they have been interpreted as a phase in the formation of ‘typical’ galaxies (e.g. Somerville, Primack & Faber 2001; Baugh et al. 2005) and are the progenitors of a reasonable fraction of massive ($L > L^*$) galaxies today.

Despite this progress, there are still a number of open questions surrounding the nature and properties of typical LBGs. Some

★E-mail: k.coppin@herts.ac.uk

LBGs have dust-corrected star formation rates (SFRs) of up to $100 M_{\odot} \text{ yr}^{-1}$, which is enough to form present-day elliptical galaxies, and indeed LBGs at $z > 3$ are attractive progenitors for the rather substantial population of passive red galaxies already in place by $z \sim 2\text{--}3$ (Kriek et al. 2006; van Dokkum et al. 2006; Stark et al. 2009). Thus, LBGs should contribute a significant portion of the submm background at $850 \mu\text{m}$ at high redshifts (Adelberger & Steidel 2000). However, determining their contribution to the submm background has been hampered by the large uncertainties that go hand in hand with deriving dust-corrected UV luminosities. The most reliable way of measuring their dust content and contribution to the submm background is *directly* through submillimetre (submm) observations, which will reveal any dust-obscured star formation activity.

Several attempts have been made to detect the rest-frame far-infrared (IR) emission in LBGs, with mixed success. There are only a handful of individual submm detections of LBGs, including Westphal-MMD11 (Chapman et al. 2000) and Westphal-MM8 (Chapman & Casey 2009), targeted as part of a sample of LBGs with high UV-derived SFRs, as well as a highly magnified gravitationally lensed LBG at $z \sim 3$ (Baker et al. 2001). Two early Submillimetre Common User Bolometer Array (SCUBA)-based statistical studies of the submm flux of LBGs with high UV-derived SFRs (Chapman et al. 2000) and of ‘typical’ canonically selected $z \sim 3$ LBGs (Webb et al. 2003) yielded no detection of the population, with an average below 1 mJy. On the other hand, Peacock et al. (2000) detected a mean flux of (0.2 ± 0.04) mJy for $0 < z < 6$ starburst galaxies with UV-SFRs of $> 2 M_{\odot} \text{ yr}^{-1}$, with the mean submm flux density increasing with UV-SFR. More recently with *Herschel*, several authors have revisited LBG stacking. Lee et al. (2012) find a marginal signal ($\sim 3\sigma\text{--}4\sigma$ level) in a single band ($500 \mu\text{m}$) for a UV-bright subset of $z \sim 4$ LBGs, and Rigopoulou et al. (2010) and Magdis et al. (2010b) have reported statistical detections of mid-IR-selected LBGs with *Herschel*, but this latter sample is clearly biased to the IR-brightest most massive and/or dustiest subset of LBGs (Huang et al. 2005; Rigopoulou et al. 2010; see also Oteo et al. 2013). As the available submm maps continue to increase in size and depth (e.g. Weiss et al. 2009), stacking is now able to sometimes yield successful statistical detections of specially chosen subsets of LBGs at $z \sim 3$ (split by stellar mass and UV-estimated SFRs; Davies et al. 2013), which has helped to reliably constrain the far-IR luminosities, obscured SFRs, dust masses, and dust temperatures in the most massive and UV luminous $z \sim 3$ LBGs. There have also been attempts to detect the dust emission in a handful of LBGs at $z \sim 5$ (e.g. Stanway et al. 2010; Davies et al. 2012), although no detections were made, making it difficult to say anything conclusive. Thus, the far-IR emission from *typical* LBGs (i.e. those selected by the canonical UV/optical broad-band colour criteria), hence their dust content and contribution to the submm background, as well as tracing their IR luminosity density as a function of redshift all remain open questions.

Here, we present a statistical (stacking) analysis of LBGs at $z \sim 3, 4$, and 5 in the United Kingdom Infrared Telescope (UKIRT) Infrared Deep Sky Survey ‘Ultra Deep Survey’ (UKIDSS-UDS; Lawrence et al. 2007) field. The UDS has been mapped with the *Herschel* Spectral and Photometric Imaging Receiver (SPIRE) at $250, 350$, and $500 \mu\text{m}$ and at $850 \mu\text{m}$ as part of the SCUBA-2 (Holland et al. 2013) Cosmology Legacy Survey (S2CLS; e.g. Geach et al. 2013). The $850 \mu\text{m}$ map is the largest submm map yet obtained at this depth ($1\sigma \approx 2$ mJy over 1 deg^2 , see Section 2.1). The availability of deep optical/near-infrared (OIR) imaging in this field, combined with the new wide-area deep $850 \mu\text{m}$ map makes this an

ideal resource to study the submm properties of LBGs, without biasing the analysis to the most massive and/or mid-IR or UV-brightest LBGs (e.g. Rigopoulou et al. 2010; Davies et al. 2013). The goal of this work is to measure the average submm flux density (which is a direct probe of obscured star formation) of LBGs selected at $z = 3, 4$, and 5 in order to estimate the bolometric luminosity of typical massive galaxies at these epochs and to assess the LBG contribution to the IR luminosity density to compare with models.

This paper is organized as follows. The data and LBG sample selections are described in Section 2. In Section 3, we present the main analysis and results of the LBG stacking in submm maps, including the LBG contribution to the submm background, LBG SEDs, and the evolution of the IR luminosity density of LBGs. In Section 4, we discuss the implications of the results and, finally, our conclusions are given in Section 5. Throughout the paper, we assume cosmological parameters of $\Omega_{\Lambda} = 0.73$, $\Omega_{\text{m}} = 0.27$, and $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (Spergel et al. 2003).

2 THE DATA

2.1 OIR imaging

Our samples are based on the deep K -band image from the UKIDSS (Lawrence et al. 2007) UDS¹ data release 8 (DR8; Almaini et al., in preparation) and co-incident multiwavelength data. The parent catalogue was extracted using SExtractor (Bertin & Arnouts 1996) on the deep K -band image ($K_{\text{AB}} \leq 24.6$) and after merging the source lists from two independent extractions. The extraction parameters were designed to recover both point-like and extended low-surface-brightness objects (see Hartley et al. 2013 for further details). In addition to the three *JHK* UKIDSS bands, the UDS has been mapped by the Canada–France–Hawaii Telescope (CFHT) Megacam u' band ($u'_{\text{AB}} \leq 26.75$; Foucaud et al., in preparation), by Subaru Suprime-cam in the optical ($B = 27.6$, $V = 27.2$, $R = 27.0$, $i' = 27.0$, and $z' = 26.0$, 5σ AB mags; Furusawa et al. 2008) and by the *Spitzer* Infrared Array Camera (IRAC; Fazio et al. 2004) in channels 1 ($3.6 \mu\text{m}$; 5σ depth of 24.2 AB mags), and 2 ($4.5 \mu\text{m}$; 5σ depth of 24.0 AB mags) via the UDS *Spitzer* Legacy Program (PI: Dunlop). The overlap of all these data sets (after masking bad regions) is 0.62 deg^2 . In addition, available X-ray (Ueda et al. 2008) and radio (Simpson et al. 2006) data were employed to remove bright active galactic nuclei (AGN). The matched multiwavelength photometry was extracted using 3 arcsec apertures centred on the positions of the K -band catalogue sources (see Simpson et al. 2012 for full details). Point spread function photometry corrections were required and applied to three of the bands (the CFHT u' band and the two IRAC channels) in order to obtain correct colours (see Hartley et al. 2013 for details). We use this final matched multiwavelength catalogue to perform our LBG selection.

2.1.1 LBG selection

The canonical LBG selection utilizes the *UGR*, or *BVR*, colour space, with ‘drop-outs’ in the bluest band efficiently isolating galaxies at $z \approx 3$ (Steidel et al. 1996). Since the selection requires detection in G and R , as the Lyman break is redshifted through the G or V band, the selection function falls off with redshift. The technique can be applied to higher redshift by simply moving the entire colour space to longer wavelengths. Ouchi et al. (2004) show how a similar

¹ <http://www.nottingham.ac.uk/astronomy/UDS/>

selection in BRi , isolates LBGs at $z \approx 4$, and Viz , or Riz , extends the selection to $z \approx 5$.

We adopt the following colour selections for LBGs at $z \approx 3$ (equation 1), $z \approx 4$ (equation 2), and $z \approx 5$ (equations 3 and 4):

$$\begin{aligned} R < 27, & \quad (U - V) > 1.2, \\ -1.0 < (V - R) < 0.6, & \quad (U - V) > 3.8(V - R) + 1.2; \end{aligned} \quad (1)$$

$$\begin{aligned} I < 27, & \quad (B - R) > 1.2, \\ (R - I) < 0.7, & \quad (B - R) > 1.6(R - I) + 1.9; \end{aligned} \quad (2)$$

$$\begin{aligned} Z < 26, & \quad (V - I) > 1.2, \\ (I - Z) < 0.7, & \quad (V - I) > 1.8(I - Z) + 2.3; \end{aligned} \quad (3)$$

$$\begin{aligned} Z < 26, & \quad (R - I) > 1.2, \\ (I - Z) < 0.7, & \quad (R - I) > (I - Z) + 1.0. \end{aligned} \quad (4)$$

For the $z = 5$ selection, we require selection using either equation (3) or equation (4), slightly improving our yield (see Ouchi et al. 2004), and for all selections we also require that the source has not been identified as a possible star.

11-band photometric ($UBVRIZJHK[3.6][4.5]$) redshifts have been estimated for all of the galaxies in the DR8 parent sample. The overview of these redshifts and how they are calculated is discussed in more detail in Hartley et al. (2013) and Mortlock et al. (2013). In summary, the photometric redshifts are computed using the EAZY template-fitting code, adopting the six default EAZY spectral energy distribution (SED) templates (Brammer, van Dokkum & Coppi 2008). In addition to this, we further include a template which is the bluest EAZY template, but with a small amount of Small Magellanic Cloud (SMC) like extinction added. Photometric redshifts are then computed from a maximum likelihood analysis. To test how accurate these photometric redshifts are, we compare the values we calculate to spectroscopic redshifts that are available in the UDS. A large fraction of these are from the UDSz, a European Southern Observatory large spectroscopic survey (ID:180.A-0776; Almaini et al., in preparation) and also from previous published values (see Simpson et al. 2012 and references therein). We have in total about 1500 secure spectroscopic redshifts in the UDSz, and around 4000 redshifts from other sources, although this reduces to 2146 after the removal of AGN. Excluding catastrophic outliers ($|z_{\text{phot}} - z_{\text{spec}}|/(1 + z_{\text{spec}}) > 0.15$), we find an average $|z_{\text{phot}} - z_{\text{spec}}|/(1 + z_{\text{spec}}) = 0.031$ (Hartley et al. 2013).

To further improve the LBG colour selections described above, we enforce a photometric redshift $z_{\text{phot}} \geq 2$ selection to eliminate any low- z interlopers (where z_{phot} is the solution with the minimum χ^2 from the EAZY fitting). The photo- z fitting procedure evaluates a probability density distribution, $P(z)$, for each galaxy, with the peak of that distribution representing the best estimate for the redshift. Apart from the $z_{\text{phot}} \geq 2$ criterion, we do not use the photometric redshifts further. However, we can make use of $P(z)$ to investigate the efficacy of the colour selections described above. In Fig. 1, we show the (normalized) sum of the individual $P(z)$ for the galaxies in the selections defined in equations (1)–(4). This clearly demonstrates the effectiveness of the colour selections, and also provides us with an estimate of the widths of the redshift distributions, reflecting both intrinsic spread (due to the ‘window’ where the redshifted Lyman break can be identified in each case), and uncertainty in the photometric redshift estimates. The resulting LBG sample properties are summarized in Table 1. Note that all three probability density distributions show a low-redshift tail, with a slight peak at $z \sim 2.5$. We

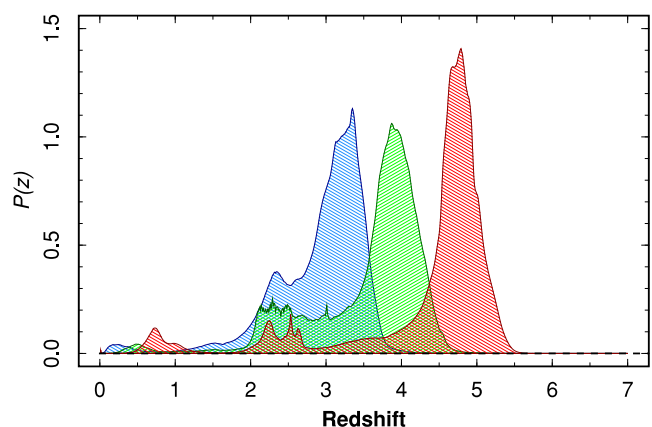


Figure 1. Stacked photometric redshift distributions from EAZY (see the text) of the sources selected in our three LBG samples (all normalized to a common integral), clearly showing the three distinct peaks in redshift space, centred at $z \sim 3, 4$, and 5 .

emphasize that our samples are constructed from colour selections described above, with a best-fitting z_{phot} lower limit strictly enforced to eliminate obvious low- z interlopers. We have verified that using a stricter $z_{\text{phot}} \geq 3$ lower limit for the $z \sim 4$ and 5 samples does not change the significance of our results in the following analysis of statistical submm detections, showing that the low- z tails in Fig. 1 do not represent significant contamination of the samples.

An important point to note is that our LBG selection is predicated on a $K < 24.6$ selected sample, that is the basis of the UKIDSS-UDS catalogue described above. As a result, we will not include, for example, $z \approx 3$ LBGs with $R < 27$ and $K > 24.6$, but these will be at the low stellar mass end of the distribution ($M_* < 10^9 M_\odot$, see Section 3.5). As we explore in Section 3.5, the average submm flux of LBGs is expected to be a function of mass, and so we treat the $K > 24.6$ limit as a proxy mass limit for the galaxies in the sample. Note also that the optical limits in equations (1)–(4) result in a slight bias in UV luminosity selection with redshift. At $z = 3$, we are sensitive to galaxies more luminous than $M_{1700} \sim -19$ mag, but only $M_{1700} \sim -20$ mag at $z \sim 5$. Therefore, when we examine trends with UV luminosity, we focus on the $z \sim 3$ sample.

2.2 SCUBA-2 observations and map making

As part of the S2CLS campaign, observations of UKIDSS-UDS were conducted in Band 2/3 weather ($0.05 < \tau_{225\text{GHz}} \leq 0.1$) between 2011 October and 2013 February. The mapping centre of the UDS field is $\alpha = 2^{\text{h}}18^{\text{m}}00^{\text{s}}$, $\delta = -05^{\circ}05'00''$, and a 3300 arcsec diameter PONG mapping pattern was used, resulting in a circular map that reaches a uniform depth over an area of ~ 700 arcmin². The total mapping time was 130 h, split into 195 scans of 40 min each. Individual scans are reduced using the dynamic iterative map-maker (MAKEMAP) of the SMURF package (Chapin et al. 2013). Geach et al. (2013) presented the reduction procedure for CLS data, but in summary, we apply the following steps.

- (i) Raw data are flat-fielded using ramps bracketing every science observation, and the data are scaled to units of pW.
- (ii) The signal each bolometer records is then assumed to be a linear combination of (a) a common mode signal dominated by atmospheric water and ambient thermal emission; (b) the astronomical signal (attenuated by atmospheric extinction, see Dempsey et al. 2013); and finally (c) a noise term (including $1/f$ -type correlations), taken to be the combination of any additional signal not accounted

Table 1. Stacked mid-IR-to-submm photometry of LBGs.

$\langle z_{\text{phot}} \rangle$	N	850 μm^a (mJy)	850 μm^b (mJy)	250 μm (mJy)	350 μm (mJy)	500 μm (mJy)	24 μm (μJy)
3.35	4201	$0.249 \pm 0.029 (8.5\sigma)$	$0.228 \pm 0.029 (7.8\sigma)$	$0.751 \pm 0.077 (9.8\sigma)$	$0.938 \pm 0.109 (8.6\sigma)$	$0.634 \pm 0.143 (4.4\sigma)$	$9.5 \pm 0.9 (10.5\sigma)$
3.87	869	$0.411 \pm 0.064 (6.4\sigma)$	$0.388 \pm 0.064 (6.0\sigma)$	$0.783 \pm 0.177 (4.4\sigma)$	$1.116 \pm 0.163 (6.8\sigma)$	$1.024 \pm 0.164 (6.2\sigma)$	$2.1 \pm 0.5 (4.2\sigma)$
4.79	68	$0.875 \pm 0.229 (3.8\sigma)$	$0.811 \pm 0.229 (3.5\sigma)$	$2.023 \pm 0.798 (2.5\sigma)$	$2.337 \pm 0.916 (2.6\sigma)$	$2.022 \pm 0.686 (2.9\sigma)$	$4.0 \pm 2.0 (2.0\sigma)$

Notes: ^aSubmm sources $\geq 3.5\sigma$ and further than 8 arcsec away from an LBG have been removed. We adopt this column for the SED fitting.

^bNo submm sources have been removed.

for by (a) and (b). The dynamic iterative map maker attempts to solve for these model components, refining the model until convergence is met, or an acceptable tolerance has been reached.

The reduction also includes the usual filtering steps of spike removal ($>10\sigma$ deviations in a moving boxcar) and DC step corrections. Throughout the iterative map-making process, bad bolometers (those significantly deviating from the model) are flagged and do not contribute to the final map.

(iii) Time streams are finally re-gridded into a map with 4 arcsec pixels, with each bolometer weighted according to its time-domain variance (which is also used to estimate the χ^2 tolerance in the fit).

Maps made from individual scans are co-added in an optimal, noise-weighted manner, using the `MOSAIC_JCMT_IMAGE` recipe in the `PICARD` environment (Jenness et al. 2008). Finally, to improve the detectability of faint point sources, we apply a beam-matched filter to improve their detectability using the `PICARD` recipe `SCUBA2_MATCHED_FILTER`. The average exposure time over the nominal 3300 arcsec mapping region in the co-added map is approximately 1.5 ks per 4 arcsec pixel, and the rms noise value is 1.9 mJy. Due to the scanning strategy, there are data beyond the 3300 arcsec perimeter, but because these receive less total exposure, the noise increases accordingly. The LBG catalogues cover a 0.6 deg^2 area within the central uniform noise region of the SCUBA-2 map.

An important component of our stacking analysis is to create a map with bright significantly detected point sources near to an LBG (but not near enough to be associated) removed. For this, we run a peak-finding algorithm on the signal-to-noise ratio (S/N) map down to a level of 3.5σ , and this was done for each LBG-redshift subsample separately. This S/N threshold was chosen to strike a balance between detecting the majority of real sources in the map while maintaining a low contamination rate due to false positives (5–10 per cent) expected from pure Gaussian noise (e.g. Geach et al. 2013). Each time a peak is found further than 8 arcsec away (about half a SCUBA-2 beam), the flux of the contributing source is removed by subtracting a model point response function (PRF) at that position, scaled to the flux of the source. The peak-normalized model PRF is generated from the data by stacking the map at the positions of all $>10\sigma$ peaks. This residual map still contains astronomical signal, but with deemed LBG-unassociated submm $>3.5\sigma$ sources removed, so that we do not inadvertently bias our stacking results high by including flux from nearby unrelated positive beams. This procedure is virtually identical to the approach taken in the Webb et al. (2003) SCUBA LBG stacking analysis paper. The remaining 36 submm sources² may not be genuine LBG-SCUBA-2 associations (and in fact the calculated Poisson probabilities that they are chance coincidences are not particularly low), but because they *could* be genuine, we must not remove these sources from the

map. These candidate LBG-SCUBA-2 associations will be evaluated using a rigorous multiwavelength approach in Coppin et al. (in preparation).

3 ANALYSIS AND RESULTS

3.1 Statistical results

With the LBG samples defined, we measure the submm flux at each LBG position in the 850 μm beam-convolved, LBG-unassociated $\geq 3.5\sigma$ submm source-subtracted flux map, weighted by the noise at the corresponding positions in the noise map. For visualization, we simultaneously stack 80 arcsec \times 80 arcsec cut-outs, again applying a simple inverse variance weighting scheme, obtaining the weights from the noise maps (note that measuring the stacked flux directly from these images necessarily gives the same result). The stacked thumbnail images for each sample are shown in Fig. 2, indicating the detection of significant 850 μm emission from LBGs at $z \approx 3$, 4, and 5 (see Table 1). We repeat the exercise above on the original flux map (i.e. including all $>3.5\sigma$ sources which had been removed) and list these results in Table 1. These stacked fluxes are statistically indistinguishable from the stacked fluxes calculated previously on the source-subtracted maps. As an additional check, we compare the radially averaged profiles of the stacks with the shape of the beam (empirically derived by stacking many high-significance peaks in the detection image). The stacks are indistinguishable from the beam, and this indicates that if there is any significant contribution to the observed signals from large-scale structure correlated with the LBGs, it is on scales below ~ 100 kpc.

To verify the significance of our results, we have stacked at the same number of uniform random positions within the same map area and repeated this 10 000 times for each LBG redshift slice subset. We do not, however, remove any sources from the map; this should provide a conservative estimate of the likelihood of measuring the same (or greater) level of flux in the stack as observed for the LBGs simply by chance. In these random realizations, we never find a result as statistically significant as ours in the $z \sim 3$ and 4 samples (the most significant stacked result we find at random in the $z \sim 3$ and 4 samples are at the 3.8σ and 3.4σ levels, respectively, occurring in each <0.02 per cent of the time), and less than 0.02 per cent of the time in the $z \sim 5$ sample. Thus, in light of these simulations, our statistical detections appear to be robust.

3.1.1 Assessing the level of bias in our statistical results

Our simple methodology implicitly assumes that the galaxy population we have stacked is uncorrelated (i.e. the LBGs are not clustered). However, like most populations, LBGs are known to be clustered ($r_0 \simeq 4 h^{-1}$ Mpc; Adelberger et al. 2005). We investigate the level of bias due to LBG clustering in our stacking results by following the formalism of Viero et al. (2013) for unbiased

² We find 26, 9, and 1 candidate LBG-SMG counterparts for our LBG redshift subsamples at $z \sim 3$, 4, and 5, respectively.

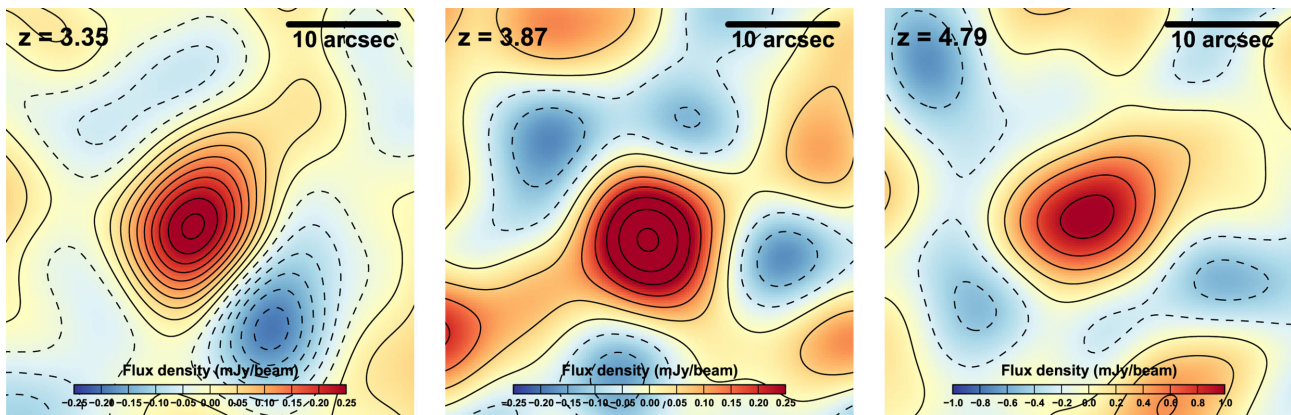


Figure 2. Thumbnail images of the stacked 850 μm flux in the SCUBA-2 map, centred on the LBG positions in the $z \sim 3, 4$, and 5 subsamples (the mean redshift of each LBG subsample is indicated on each image). The colour scale indicates the average flux density and the solid contours represent significance levels of $0\sigma, 1\sigma, 2\sigma, 3\sigma, \dots$, and the dashed lines are the corresponding to negative deviations.

stacking of galaxy catalogues on *Herschel* maps (see also Kurczynski & Gawiser 2010 and Roseboom et al. 2010), which takes into account the presence of multiple sources in the same beam. In brief, we create a ‘hits’ map (H) of the same size as the SCUBA-2 map for a given LBG catalogue, where each pixel in the map contains the integer number of sources that fall into it. The hits map is then convolved with the *actual* PRF of the instrument (note that a different result, yielding a non-negligible positive flux bias, occurs if the beam is wrongly assumed to be a simplistic Gaussian). The convolved hits map is then regressed against the 850 μm flux map (S_{850}) to find the minimal flux residual $|(S)H - S_{850}|$, where (S) is the average flux density of the LBG population. The method safely assumes that galaxies in different redshift slices are uncorrelated and we thus treat each of our LBG lists independently. The fitting routine yields minima at $S_{850} = (0.21 \pm 0.04), (0.39 \pm 0.07)$, and (0.86 ± 0.26) mJy for $z \sim 3, 4$, and 5 , respectively (with the error bars calculated from Monte Carlo simulations). These results are indistinguishable from the results in columns 3 and 4 of Table 1. It is reassuring that our statistical detections do not appear to be dominated by an upwards bias in flux introduced by potentially ‘double counting’ flux for multiple sources falling in the same beam.

There is also the legitimate worry that LBGs are also correlated with some fainter (as yet undetected, and potentially numerous) population that contributes to the observed submm flux density, thus boosting our statistical detections (e.g. Chary & Pope 2010; Kurczynski & Gawiser 2010; Serjeant et al. 2010). However, for the purposes of our analysis, we would say that this was flux associated with the LBG. Whether or not the emission is always physically within the LBG is beyond the scope of our current study.

3.2 The LBG contribution to the submm background at 850 μm

The contribution to the submm background from LBGs is still poorly constrained; however, our data can finally address this question, since we now have a robust detection of the average 850 μm flux density of LBGs at three epochs, at least those with ultraviolet luminosities of $L_{1700} \approx 10^{29} \text{ erg s}^{-1} \text{ Hz}^{-1}$, characteristic of galaxies in our sample (Table 2). Using our average stacked 850 μm flux densities, we estimate surface brightness densities of 1700, 600, and 100 mJy deg^{-2} of LBGs at $z \sim 3, 4$, and 5 , respectively. By comparison, the total background at 850 μm inferred from *COBE-FIRAS* is $3.1\text{--}4.4 \times 10^4 \text{ mJy deg}^{-2}$ (Puget et al. 1996; Fixsen et al.

Table 2. Derived average IR and 1700 \AA luminosities and corresponding obscured (IR) and unobscured (UV) SFRs of $z \sim 3\text{--}5$ LBGs.

$\langle z_{\text{phot}} \rangle$	L_{IR} ($10^{11} L_{\odot}$)	SFR_{IR} ($M_{\odot} \text{ yr}^{-1}$)	L_{1700} ($10^{29} \text{ erg s}^{-1} \text{ Hz}^{-1}$)	SFR_{UV} ($M_{\odot} \text{ yr}^{-1}$)
3.35	$3.2^{+0.8}_{-0.6}$	55^{+14}_{-10}	0.9	18
3.87	$5.5^{+0.3}_{-0.4}$	93^{+5}_{-7}	1.7	33
4.79	$11.0^{+4.2}_{-5.9}$	186^{+71}_{-101}	1.7	33

1998; Hauser & Dwek 2001; Lagache, Puget & Dole 2005). Summing these separate surface brightness densities together, we find that the LBGs with $L_{1700} > 10^{29} \text{ erg s}^{-1} \text{ Hz}^{-1}$ in our $z \sim 3, 4$, and 5 samples comprise around 6–8 per cent of the submm background at 850 μm (where the range of values simply reflects the uncertainty in the *COBE-FIRAS* result). However, the true contribution from LBG-like galaxies to the submm background will come from a wider range in redshift, not just from the rather narrow redshift slices we have sampled (see Fig. 1), and from sources falling out of our samples due to incompleteness. To determine the total (corrected) contribution from LBGs over $3 < z < 5$, we assume that LBGs have a constant comoving number density (which is a reasonable assumption since the bright end of the luminosity function for LBGs shows little evolution over $3 < z < 5$; Bouwens et al. 2007; McLure et al. 2009; Reddy & Steidel 2009). Starting with the $z \sim 3$ sample (which is our most complete subsample of LBGs), we can integrate over the comoving volume element for this redshift range, scaling the LBG submm background contribution accordingly. We find that the total contribution to the submm background from LBGs over the redshift range $3 < z < 5$ is likely to be closer to 14–20 per cent. This result is consistent with Webb et al. (2003), who estimated an upper limit to the contribution to the submm background from $1 < z < 5$ of less than 20 per cent.

3.3 The bolometric luminosities and total SFRs of LBGs

In order to put LBGs in context with other IR-luminous galaxy populations, we must determine some intrinsic properties, such as their IR luminosities and obscured and unobscured SFRs. Our approach is to make use of the full (rest-frame UV to submm) average SEDs of LBGs at each epoch, where the photometry are the average for the sample.

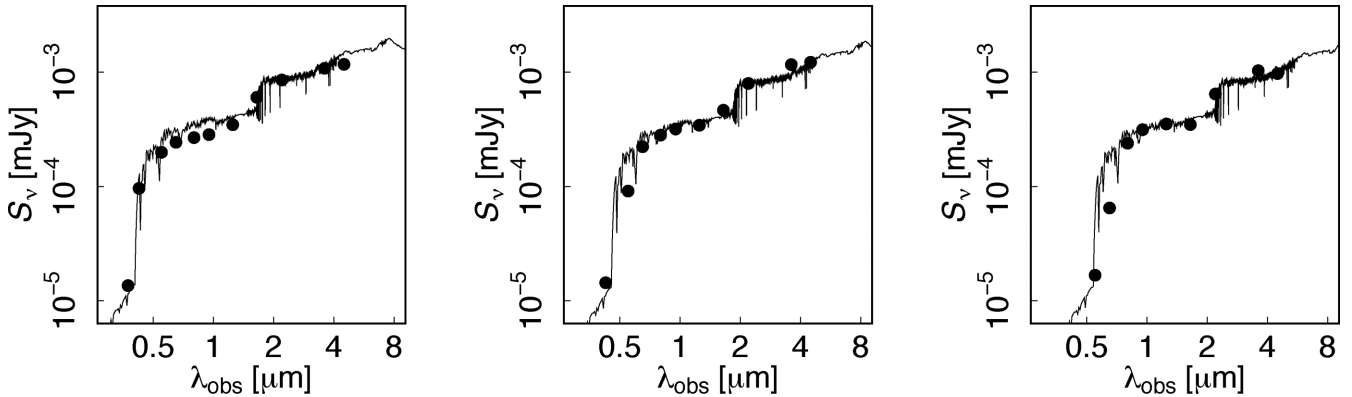


Figure 3. Best-fitting Bruzual (2007) model for $z = 3$ LBGs selected at $24\ \mu\text{m}$ from Magdis et al. (2010b), normalized to the observed K -band flux of each LBG sample at $z = 3, 4$, and 5 . This simple SED scaling provides an overall excellent description of the observed average photometry, with only a slight discrepancy in the UV emission at $z = 4$ and 5 , most likely due to increasing $\text{Ly}\alpha$ forest absorption at $z > 3$ (see the text).

For the average rest-frame UV–near-IR SED of our LBGs, we find that the best-fitting Bruzual (2007) model for $z = 3$ LBGs selected at $24\ \mu\text{m}$ selected from Magdis et al. (2010b) provides a reasonable model of the stellar emission in our *non*-IR selected sample. In Fig. 3, we show the average photometry of the three LBG samples, with the Magdis et al. (2010b) SED redshifted and scaled to the observed K -band flux. The template SED provides an excellent description of the observed average photometry, with only a slight discrepancy in the UV emission at $z = 4$ and 5 , most likely due to increasing $\text{Ly}\alpha$ forest absorption at $z > 3$. We use this template to calculate k -corrections in the UV–near-IR part of the SED.

For the mid-IR-to-submm part of the SED, we use the 185 SED template library constructed by Swinbank et al. (2014), which includes local galaxy templates from Chary & Elbaz (2001), Rieke et al. (2009), and Draine et al. (2007), as well as the SEDs from well-studied high-redshift starburst galaxies SMM J2135–0102 ($z = 2.32$) and GN20 ($z = 4.05$) from Ivison et al. (2010) and Carilli et al. (2011), respectively. We note that the range of temperatures in the templates is fairly broad: 19–60 K. We fit these templates to our observed far-IR stacked photometry of the LBGs using a χ^2 minimization approach, but allowing the redshift and normalization of the templates to vary. To improve the constraints and accuracy of the SED fitting, we have also extracted stacked photometry for our LBG subsamples based on *Herschel* SPIRE maps of the UDS at 250, 350, and 500 μm from the *Herschel* Multi-tiered Extragalactic Survey (Oliver et al. 2012). The Level 2 data products from the *Herschel* European Space Agency (ESA) archive were retrieved, aligned, and co-added to produce maps. For the stacking, we use the raw map from Swinbank et al. (2014), which has had the mean flux of 1000 random positions subtracted, effectively removing any systematic contribution from the background or confusion. In addition, this map includes re-injected flux from any LBG identified in the $24\ \mu\text{m}$ catalogue used to make a deblended version of the map³ in Swinbank et al. (2014, including 142, 15, and 2 in the $z \sim 3, 4$, and 5 samples, respectively). The stacked results are given in Table 1. We note that significant (i.e. $>3\sigma$) stacked detections in

all three SPIRE bands are made at $z \sim 3$ and 4 , but not at $z \sim 5$. To provide a constraint on the Wien side of the spectrum, we have also stacked the LBGs in $24\ \mu\text{m}$ Multiband Imaging Photometer for *Spitzer* (MIPS; Rieke et al. 2004) imaging using a 13 arcsec aperture radius and apply the standard aperture corrections. This yields stacked detections of the $z \sim 3$ and 4 LBGs but not for $z \sim 5$ LBG sample, and the results are tabulated in Table 1. Thus, we use the stacked fluxes and errors from Table 1 for the $z \sim 3$ and 4 samples in the SED fitting, but convert the SPIRE and MIPS fluxes of the $z \sim 5$ sample to 3σ upper limits. The best-fitting templates are shown in Fig. 4, along with the range of models which lie within 1σ of the best fit that are used to calculate uncertainties on our SED-derived parameters. Note that our derived SED parameters do not change appreciably (i.e. are well within the 1σ errors) when we exclude the MIPS data from the fitting.

We use the best-fitting SEDs to compute the basic properties L_{IR} and SFR. To calculate the L_{IR} values of our average LBG at $z \sim 3, 4$, and 5 , we integrate under the best-fitting SEDs between 8 and 1000 μm and tabulate the results in Table 2. We then calculate the corresponding SFR for our average LBG (see Table 2) following Kennicutt (1998): $\text{SFR} (M_{\odot} \text{ yr}^{-1}) = 1.7 \times 10^{-10} L_{\text{IR}} (L_{\odot})$. This relation assumes that the IR luminosity is predominantly powered by star formation (i.e. a negligible contribution from an AGN, which is assumed to be a good assumption in general for LBGs; e.g. Shapley et al. 2005; Huang et al. 2007; Rigopoulou et al. 2010) and comes from a starburst less than 100 Myr old, with a Salpeter (1955) initial mass function (IMF). We find that the best-fitting SED template for the $z \sim 3$ sample has a dust temperature of $T_{\text{dust}} = 37\ \text{K}$ and the best-fitting SED template for the $z \sim 4$ and 5 samples has $T_{\text{dust}} = 38\ \text{K}$ (although note that the derived dust temperature for the $z \sim 5$ LBG SED is not a very reliable constraint since the photometry are all upper limits except for 850 μm flux density data point). For reference, we also fit the SPIRE+SCUBA-2 photometry using a simple modified blackbody assuming a dust emissivity index of $\beta = 1.5$ (e.g. Blain et al. 2002) with the Wien side of the spectrum modified by a power law of the form $S_{\nu} \propto \nu^{-1.7}$. The derived dust temperatures from the SED template fitting are lower (by 5–25 per cent, with the percentage difference increasing with redshift) compared to the modified blackbody method, and the L_{IR} values from the modified blackbody fits are a factor of ≈ 2 –3 times higher than the SED templates, demonstrating the additional systematic

³ The raw SPIRE maps, deblended model, catalogues and residual maps used for this analysis are available at <http://astro.dur.ac.uk/~ams/HSODEblend/UDS/>

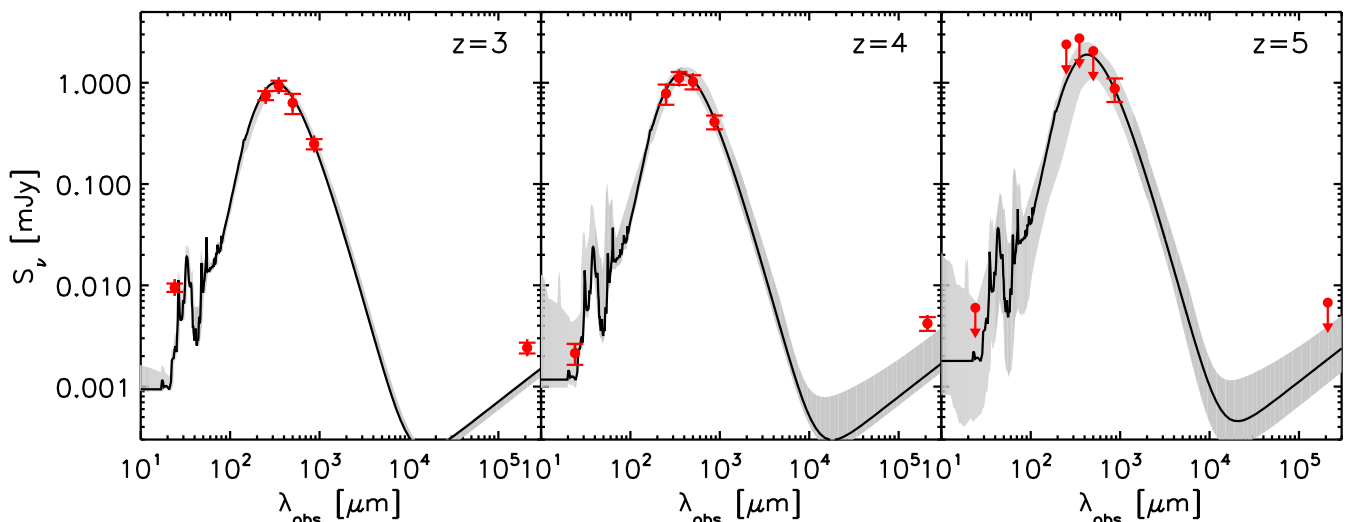


Figure 4. Best-fitting far-IR SEDs of the LBG samples at $z = 3, 4$, and 5 , with the shaded regions showing the range of acceptable solutions (see the text). We have used the stacked fluxes (or 3σ upper limits, where appropriate) at $24, 250, 350, 500$, and $850 \mu\text{m}$ in the SED fitting. Also overplotted (i.e. not fitted) are the median stacked 1.4 GHz radio fluxes, which appear to lie mildly (1σ – 3σ) higher than the best-fitting far-IR constrained SEDs would predict for the $z \sim 3$ and 4 samples, hinting that either the radio emission in the template is not representative of $z \sim 3$ – 4 star-forming galaxies and/or that there is some AGN contribution to the radio fluxes of typical LBGs (see the text).

uncertainty involved in far-IR SED fitting that should be taken into account.

For each LBG redshift subset, we also perform median stacking⁴ of the LBG samples in the UDS Jansky Very Large Array 1.4 GHz data. The median stacked radio fluxes are $S_{1400} = (2.42 \pm 0.30)$, (4.20 ± 0.65) , and $(5.06 \pm 2.25) \mu\text{Jy}$, for our $z \sim 3, 4$, and 5 LBGs, respectively. The fluxes were obtained using the JMFIT task within AIPS, where the width of the Gaussian fit to the stack was fixed to be the full width at half-maximum (FWHM) of the corresponding stack at the positions of sources in a K -band catalogue. The whole K -band catalogue was used to provide a high S/N image from which the Gaussian parameters could be determined. This is used over the synthesized beam FWHM to account for bandwidth smearing effects. The radio flux extraction will be described in more detail in Arumugam (in preparation). The radio stacks are simply overplotted on the SEDs in Fig. 4 as detections (at $z \sim 3$ and 4) or upper limits ($z \sim 5$) using a nominal $>3\sigma$ threshold. We note that the detected radio stacks at $z \sim 3$ and 4 lie 1σ – 3σ higher than the best-fitting far-IR constrained SEDs would predict, hinting that either the radio emission in the template is not representative of $z \sim 3$ – 4 star-forming galaxies and/or that there is some AGN contribution to the radio fluxes of typical LBGs (although, only ~ 3 per cent of LBGs show any AGN activity; Steidel et al. 2002; Laird et al. 2006).

Nevertheless, the LBGs appear to be of the Luminous IR Galaxy (LIRG)-to-Ultraluminous IR Galaxy (ULIRG) class ($L_{\text{IR}} \sim 10^{11}$ – $10^{12} L_{\odot}$), with corresponding SFRs of several tens to hundreds of $M_{\odot} \text{ yr}^{-1}$. These results are consistent with previous work within the photometric uncertainties and systematic errors involved in SED fitting, but our deep stacked fluxes have enabled us to provide physical constraints on the far-IR SEDs of canonically selected ‘typical’ LBGs for the first time, and are not limited to just the most massive or UV-brightest subsets (e.g. Lee et al. 2012; Davies et al. 2013).

⁴ A median stack minimizes any bias due to a few number of radio-loud outliers.

3.4 The evolution of the IR luminosity density of LBGs

We have shown that LBGs contribute a non-negligible fraction of the submm background. Our stacking results provide a new suite of constraints for models of luminosity density and galaxy formation as they directly probe the far-IR density of the Universe from $z \sim 3$ to 5 down to the LIRG level of energy output ($\approx 10^{11} L_{\odot}$).

We can estimate the IR luminosity density of LBGs at $z \sim 3, 4$, and 5 by multiplying the LBG volume densities by the corresponding average LBG L_{IR} . To calculate the volume densities, we integrate the UV luminosity functions from Reddy & Steidel (2009), Bouwens et al. (2007), and McLure et al. (2009) for $z \sim 3, 4$, and 5 , respectively, down to an absolute UV magnitude (M_{1700}) of $M_{1700}^* + 1$ at each redshift. Our choice of UV luminosity limit is motivated by the limitation that $M_{1700}^* + 1$ is approaching the limiting optical depth of the survey, with the R, I , and z band sampling the rest-frame 1700 \AA emission at $z = 3$ – 5 . Integrating the luminosity function to a lower luminosity introduces uncertainty since we enter a regime where the infrared emission of low-(UV)-luminosity LBGs is not well measured, and therefore their contribution to the IR luminosity budget is uncertain. Limiting to $M^* + 1$ ensures a reasonably conservative estimate of the LBGs’ contribution to the luminosity density, whilst sampling a representative proportion of the population, below the knee in the luminosity function. Obviously, integrating further down the UV luminosity function whilst applying the same canonical average L_{IR} would result in a higher luminosity density, but it is likely that generally L_{IR} is falling with decreasing L_{UV} .

For LBGs with UV luminosities above $M_{1700}^* + 1$, we find corresponding IR luminosity densities of $(4.0^{+2.3}_{-1.5}) \times 10^8$, $(5.0^{+2.9}_{-1.8}) \times 10^8$, and $(7.9^{+7.9}_{-4.0}) \times 10^8 L_{\odot} \text{ Mpc}^{-3}$ for $z \sim 3, 4$, and 5 , respectively. In Fig. 5, we have compared our results alongside the models of the IR luminosity density from Béthermin et al. (2011) and other available measurements from the literature (Pascale et al. 2009; Rodighiero et al. 2010). Pascale et al. (2009) stacked $24 \mu\text{m}$ -selected sources, limiting their study to $z < 1.5$ – 2 , in the submm using BLAST data (Devlin et al. 2009). Here, using the efficient

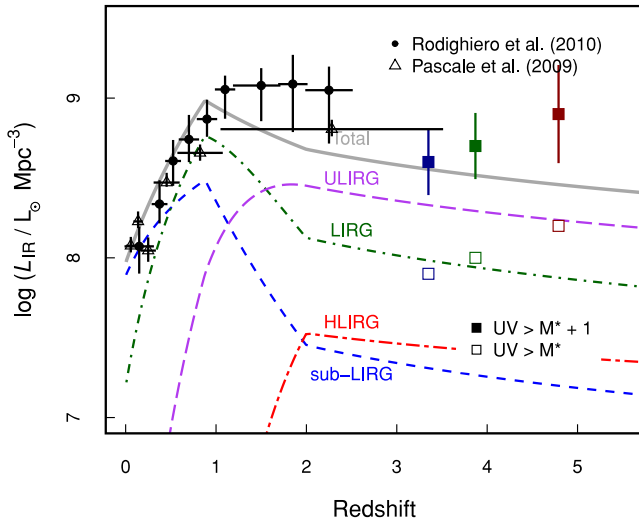


Figure 5. Evolution of the bolometric IR luminosity density (solid line) as a function of redshift from the parametric backward evolution model of Béthermin et al. (2011). Also shown are the individual contributions in the models from normal galaxies ($L_{\text{IR}} < 10^{11} L_{\odot}$), LIRGs ($10^{11} < L_{\text{IR}} < 10^{12} L_{\odot}$), ULIRGs ($10^{12} < L_{\text{IR}} < 10^{13} L_{\odot}$), and HyLIRGs ($L_{\text{IR}} > 10^{13} L_{\odot}$). Measurements from Pascale et al. (2009) and Rodighiero et al. (2010) are overplotted, along with our results. To derive the IR luminosity density, we have multiplied the space density of LBGs at each redshift by the average L_{IR} we have measured, which is obviously affected by the choice of luminosity integration limit; we show the effect in integrating down to M^* and $M^* + 1$. These results extend observational constraints of the models from $z = 3$ to 5, showing that the observations of LBGs are broadly consistent with the model predictions for the LIRG class of galaxies for our $z \sim 3$ and 4 LBGs and with the ULIRG class for the $z \sim 5$ LBGs.

canonical LBG selection, we are able to reliably extend the model constraints to $z \sim 3$ and beyond for the first time. It appears that our measured IR luminosity densities for canonically selected LBGs at $z \sim 3$ –5 lie mildly high on average, but are broadly consistent with the Béthermin et al. (2011) model predictions of LIRGs at $z \sim 3$ –4 and ULIRGs at $z \sim 5$. It should be noted that the relative contribution of different luminosity classes to the infrared luminosity density predicted by Béthermin et al. (2011) is poorly constrained at $z > 1$, especially for low-luminosity galaxies. The contribution of LIRGs and sub-LIRGs which is decreasing towards higher redshifts is a result of the evolving luminosity function in the Béthermin model, which sees a steady increase in L_* to $z \approx 2$ and a sharp decline in ϕ_* across $z \approx 1$ –2, whilst the faint end slope is held constant. The poor empirical constraints on the actual form of the bolometric luminosity function, in particular the faint end slope, at $z > 2$ should be taken into account when interpreting our results in the context of models such as this.

Obviously the derived luminosity density depends on our choice of luminosity limit when integrating the UV luminosity function, since in this calculation we are assuming a fixed average L_{IR} for all LBGs at each epoch. We have chosen $M_{1700}^* + 1$ to roughly match the depth of our catalogue; however, we show the effect of integrating down to M_{1700}^* in Fig. 5, which bring the luminosity densities more in line with the predictions of Béthermin et al. (2011) for galaxies of this luminosity class. Interpreting these observations as a lower limit to the total infrared luminosity density at $z \approx 3$ –5, and assuming that the luminosity density peaks at $z \approx 1$ –2, our observations are consistent with little, or slow, evolution in the infrared luminosity density over $z \approx 3$ –5.

3.5 Variation of submm luminosity with stellar mass

Although the LBG approach is efficient at selecting relatively normal, massive galaxies at high redshift, it is important to consider that the ‘LBG’ population is rather a broad demographic. By definition, the rest-frame ultraviolet/optical colours of LBGs are quite uniform, as enforced by the colour selection, but the lack of constraints on the observed near-infrared (and longer wavelength) properties results, not surprisingly, in a rather large range of (rest-frame) OIR colours in LBG samples (e.g. Rigopoulou et al. 2006). As pointed out in Davies et al. (2013), there could be a substantial variation in the average submm flux density within each of the LBG samples that we calculated in Section 3.1. This can be investigated directly by examining the individual submm flux densities of the LBGs. Unfortunately, the 1σ scatter in the flux densities of individual sources in each LBG subset is similar to the map noise (≈ 2.0 mJy), indicating that we do not have sufficient S/N to detect any real variation or intrinsic scatter in the submm properties of the bulk of the LBG population. Nevertheless, a reasonable expectation is that the far-IR luminosity could be a strong function of stellar mass (e.g. Davies et al. 2013), which we now investigate.

The UV/OIR SEDs described above provide a means of selecting LBGs within each redshift bin by stellar mass, thus allowing us to apply the appropriate k -correction to estimate the absolute rest-frame K -band luminosity from the observed K -band magnitude. Using UKIRT filters, the k -corrections are 2.6, 2.8, and 2.8 mag for $z = 3, 4$, and 5, respectively. We have assumed a fixed template for the LBGs to evaluate the k -corrections; however, in reality the galaxies will have a range of spectral shapes, due to (primarily) variations in extinction and age. In the following analysis, we apply the canonical factors above because we lack the information to k -correct on a case-by-case basis, but as a guide to the likely range in k -corrections due to varying spectral shapes, an additional $A_V = 1$ mag of extinction on the templates correspond to approximately 0.3 mag differences in the k -corrections. This ‘uncertainty’ should be kept in mind in the following analysis and interpretations.

Adopting M_K as an empirical proxy for stellar mass (thus allowing us freedom in the interpretation by assuming different M_*/L_K), we repeat the stacking procedure in bins of M_K . The results are shown in Fig. 6, which reveals a positive correlation between submm flux density and stellar mass at the highest redshifts, with the most massive LBGs at $z \sim 5$ tending to be the most submm luminous on average. There is also a mild hint of evolution in the data – with the $z \sim 5$ LBGs being more IR luminous than similarly massive LBGs at $z \sim 3$ and 4 – albeit not at a very significant level ($\approx 2\sigma$). On closer inspection, we find that the stacked submm flux result for the highest mass bin in the $z \sim 5$ sample is dominated by a single $\sim 4\sigma$ submm source in the map at one of the LBG locations. So to be conservative, one could just consider the bins with $M_K \geq -21$ mag and the overall trend persists. We discuss this further in Section 4.

3.6 Comparison of submm-derived SFRs with UV-derived SFRs

For star-forming galaxies, we expect that the rest-frame UV continuum and the far-IR emission should be correlated if $\tau \leq 1$, since for a given dust covering fraction, the UV continuum emission should on average scale with the energy re-radiated in the far-IR. It also follows that the dust absorption should be correlated with UV reddening, and this has been seen for local starburst galaxies (Meurer, Heckman & Calzetti 1999) and at higher redshifts from $z \sim 1$ –3 (e.g. Reddy et al. 2006, 2012; Magdis et al. 2010a,b; Heinis et al. 2013).

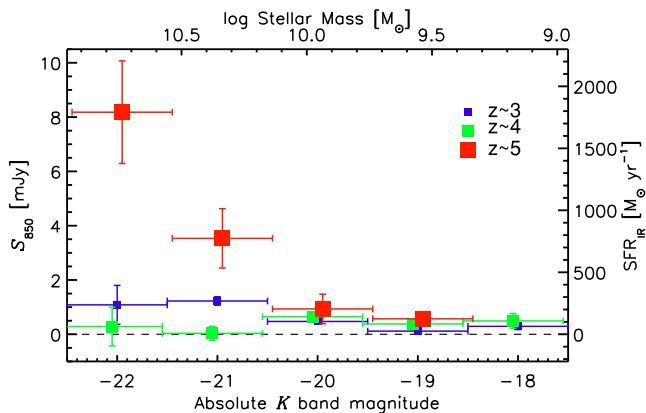


Figure 6. Measured 850 μm flux density (as a proxy for IR luminosity or SFR, shown on right-hand axis, using an average approximate conversion from our best-fitting SEDs of $1 \text{ mJy} = 220 M_{\odot} \text{ yr}^{-1}$) as a function of absolute rest-frame K -band magnitude (as a proxy for stellar mass, shown on the top axis, assuming an average mass-to-light ratio here for illustration purposes for the LBGs of 1.15, which assumes a Salpeter 1955 IMF). The plot reveals a positive correlation between submm flux density and stellar mass at $z \sim 5$, with the most massive LBGs being more submm luminous on average, although the first bin at $M_K \approx -22$ mag is dominated by a single $\sim 4\sigma$ source. However, to be conservative, one could just consider the data at $M_K \geq -21$ mag, where we see the trend persist.

It is therefore difficult to distinguish between a galaxy with a truly high SFR and one with an intrinsically low SFR and large photometric errors. Correcting for dust attenuation using only optical data is inherently difficult and uncertain, given the degeneracy between the average age of the stellar population and effect of reddening, and thus the need for excellent sampling of the UV/OIR spectral range. Submm observations provide a clean route to assessing the level of dust extinction, and hence the amount of star formation that is ‘obscured’ from view in the optical.

Shapley et al. (2001) discuss that the more intrinsically luminous LBGs appear to be dustier, with redder colours (and should therefore be brighter in the submm). Chapman et al. (2000) reported that submm-bright LBGs tended to have extremely red colours compared to the average of the population. This was confirmed by Reddy et al. (2012) who showed using *Herschel* 100 and 160 μm data that the reddest of their 146 LBGs at $z \sim 2$ indeed contain more dust, finding remarkable agreement between the local and high-redshift UV attenuation curves (e.g. Meurer et al. 1999). With our large LBG catalogues and deep SCUBA-2 map, we can now explore these issues in more detail and with greater statistical significance. Here, we focus on our largest LBG sample at $z \sim 3$ since this is where we have sufficient S/N to split up the sample further. To estimate the rest-frame UV (1700 \AA) luminosities, we simply calculate the luminosity from the observed flux in the observed R band, with no k -correction. The associated UV-derived SFRs are obtained by employing the relation in Kennicutt (1998) to the absolute M_{1700} magnitudes of the LBGs (see Table 2).

In Fig. 7, we plot measured 850 μm flux (as a proxy for obscured SFR) versus the UV-derived SFRs (uncorrected for dust) – this provides a measure of the amount of ‘dust-obscured’ versus ‘unobscured’ or ‘visible’ star formation. We find an inverse correlation of submm flux with UV SFR, that is, the LBGs that are faintest in the UV are also the most submm bright on average. Are the faint UV population intrinsically the dustiest or is this simply a mass selection effect? To explore this, we split up our sample into three bins of equal mass width across our range of stellar mass:

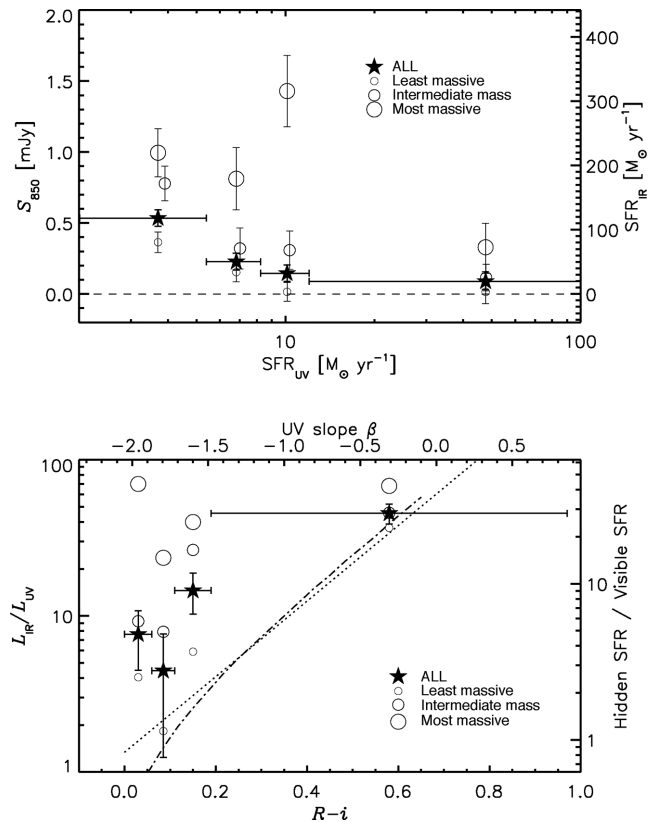


Figure 7. Trends of submm flux, optical colours, and UV-estimated SFRs. Top panel: the weighted mean 850 μm flux density as a function of UV-estimated SFR for LBGs at $z \sim 3$, showing an increase in submm flux density with lower UV SFR. The error bars represent the standard error of the mean. Approximate corresponding IR-SFRs are given on the right-hand axis for reference and are calculated using an average approximate conversion of $1 \text{ mJy} = 220 M_{\odot} \text{ yr}^{-1}$ from our SED fitting results for the $z \sim 3$ sample. Bottom panel: the mean $L_{\text{IR}}/L_{\text{UV}}$ as a function of $R - i$ colour for different mass bins (with the approximate corresponding UV slope, β , given on the top axis for reference), showing that $z \sim 3$ LBGs with the reddest colours are also the most dust obscured, and that this trend scales up with mass. The error bars represent the standard error of the mean. The dotted curve shows the expected $L_{\text{IR}}/L_{\text{UV}}$ as a function of optical reddening, assuming an SMC-like extinction curve (for reference, the dot-dashed curve of Meurer et al. 1999 is also shown). Taken together, these plots show directly that for a fixed mass, the reddest LBGs are the most submm luminous (i.e. are the most dust obscured), with the average submm luminosity-to-UV luminosity ratio increasing with mass.

$M_{\star} < 5 \times 10^9 M_{\odot}$ (least massive); $5 \times 10^9 \leq M_{\star} \leq 1 \times 10^{10} M_{\odot}$ (intermediate mass); and $M_{\star} > 1 \times 10^{10} M_{\odot}$ (most massive). At a fixed mass, we again see that submm flux density increases with decreasing UV SFR, but now also that the average submm flux increases with the stellar mass.

What is the distribution of obscured-to-unobscured star formation in the LBGs? We can explore this by plotting the average $L_{\text{IR}}/L_{\text{UV}}$ ratio versus (observed) $R - i$ colour for the LBGs, using the same stellar mass binning as above (Fig. 7). This plot demonstrates that on average, LBGs have $L_{\text{IR}}/L_{\text{UV}} \sim 5$, but for the reddest ($\beta \approx -0.3$) LBGs in all mass subsets the obscured-to-unobscured ratio is up to an order of magnitude larger than this, with $L_{\text{IR}}/L_{\text{UV}} \sim 50$. As Fig. 7 shows, for the population as a whole, the increase in $L_{\text{IR}}/L_{\text{UV}}$ scales with β in a manner broadly in agreement with that expected if we assume an SMC-like extinction curve and assume that all of

the optical extinction is caused by dust that is absorbing and re-emitting the UV radiation field. However, we also see evidence that there is a variation in the $L_{\text{IR}}/L_{\text{UV}}-\beta$ trend with mass, with the most massive LBGs showing the highest $L_{\text{IR}}/L_{\text{UV}}$ for all colour bins. The combination of these two plots demonstrates that (a) at a fixed mass, we can relate optical reddening, as parameterized by the UV continuum slope β , to dust-reprocessed emission in the observed submm, and (b) the most massive LBGs also tend to have higher obscured-to-unobscured ($L_{\text{IR}}/L_{\text{UV}}$) ratios (are more extinguished) on average than the less massive LBGs, hinting at different dust properties/geometries across the mass range. We discuss this further below.

4 DISCUSSION: EXPLORING THE LIRG POPULATION AT $z > 3-5$

We detect a statistical signal in the submm for the LBG populations at $z \sim 3, 4$, and 5. These stacked fluxes tell us about the ‘average’ submm properties of LBGs at each epoch, and the mean submm flux density tends to increase with redshift. A more careful analysis reveals that it appears to be a simple mass-selection effect – at higher redshifts lower mass (and lower overall luminosity) objects drop out of the sample, mimicking the trend that LBGs increase in IR luminosity with redshift. When this is taken into account, by splitting our redshift subsamples further into stellar mass bins, the data show no convincing evidence of strong redshift evolution in submm flux (as a proxy for bolometric luminosity) with absolute rest-frame K -band flux (as a proxy for stellar mass). For comparison to the well-established ‘main sequence’ of star-forming galaxies, which is seen to evolve with redshift, we plot our $z \sim 3, 4$, and 5 LBG total (unobscured+obscured) SFRs versus stellar mass in Fig. 8, alongside observed correlations from $z = 0-3$ (Elbaz et al. 2007; Noeske et al. 2007; Daddi et al. 2008; Magdis et al. 2010a). When plotted in this way, our data show that LBGs at $z \sim 3, 4$, and 5 might follow the so-called ‘main sequence’ of star formation activity, with the most massive LBGs tending to have the highest overall SFRs. The current data are not strongly constraining on the evolution of the scaling between stellar mass and SFR for LBGs over redshifts 3–5, but hint that the ‘main sequence’ (for LBGs at least) shows little evolution over this range.

Because we are not prone to dust obscuration effects in the submm (in contrast to the UV/OIR), we can place strong and definitive constraints on the submm background for the first time in this luminosity and redshift regime. We have estimated that LBGs at $3 < z < 5$ contribute 6–8 per cent of the 850 μm background, which rises to 14–20 per cent once we have corrected for our narrow redshift slices across this range. To put this in perspective, the bright ($S_{850} > 2$ mJy) submillimetre galaxy (SMG)⁵ population makes up 20–30 per cent (e.g. Coppin et al. 2006; Weiss et al. 2009; Swinbank et al. 2014) of 850 μm background, and the contribution from > 1 mJy SMGs peaks at ~ 50 per cent of the SFRD at $z \sim 2$ (Wardlow et al. 2011). It was thought that if SMGs and LBGs are essentially the same population, but with LBGs just being fainter and more numerous, then LBGs could make up the rest of the ‘missing’ submm background at 850 μm (75–80 per cent; Adelberger & Steidel 2000). Our measurement shows that at $3 < z < 5$ LBGs with (dust-uncorrected) UV luminosities greater than $M_{1700}^* + 1$ contribute at most 20 per cent

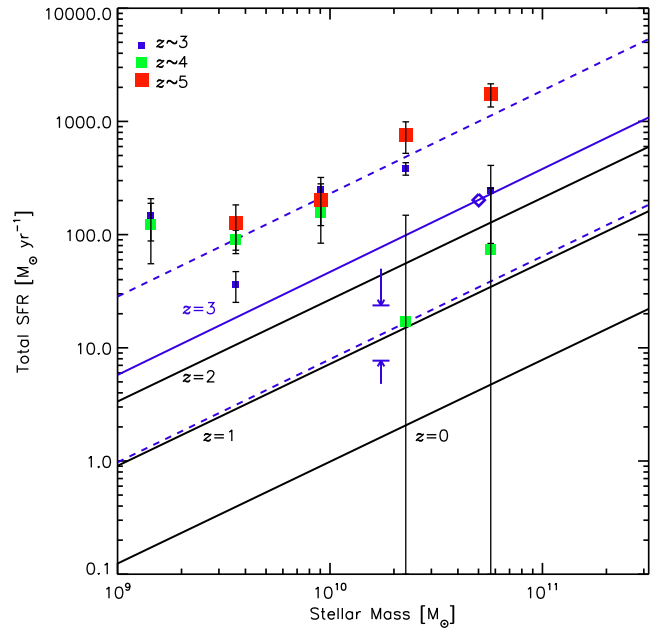


Figure 8. The ‘main sequence’ of star-forming galaxies predicts increasing star formation activity with stellar mass, and the relation is known to evolve with redshift. The observed correlations are plotted at $z = 0$ (Noeske et al. 2007), $z = 1$ (Elbaz et al. 2007), $z = 2$ (Daddi et al. 2008), and $z = 3$ (from Magdis et al. 2010a, based on IRAC-detected LBGs). The blue dashed lines simply represent the extent of the 1σ spread in the Magdis et al. (2010a) data, to guide the eye to the likely range of values at $z \sim 3$. This plot has been adapted from Davies et al. (2013) and their limits for typical $z \sim 3$ LBGs are shown as blue arrows, as well as their submm detection of IRAC-22.5 (high-mass) subsample. All of the data points and correlations shown take into account the total (obscured+unobscured) SFRs. Similar to Fig. 6 (which only shows the trend for the obscured SFRs with stellar mass), our new data hint that the SFRs of LBGs at $z \sim 3, 4$, and 5 scale with mass, with the most massive LBGs tending to have the highest overall SFRs.

of the 850 μm submm background, implying that the majority of the background is emitted at $z < 3$. Of course, the remaining ~ 50 –60 per cent of the 850 μm background, unaccounted for by SMGs and LBGs, could be feasibly contributed by star-forming galaxies at lower redshifts, e.g. BX/BMs at $z \approx 2$, or less UV/optically luminous ‘normal’ galaxies (Smith et al., in preparation).

Now that we know the average 850 μm flux of the LBG populations at different epochs, we can test this conjecture by looking at the amount of overlap of LBGs and SMGs in the 850 μm number counts. Turning to the deepest and most tightly constrained 850 μm number counts available from the cluster lensing fields of Zemcov et al. (2010), at our average detected stacked flux of ≈ 0.25 mJy and surface density of $z = 3$ LBGs suggests that LBGs comprise ≈ 35 per cent of the SMG population in this sub-mJy regime. We find that the LBGs comprise increasingly smaller fractions of the SMG population as we probe to higher 850 μm fluxes, namely ≈ 15 per cent at ≈ 0.4 mJy for the $z \sim 4$ LBGs and ≈ 3 per cent at ≈ 0.9 mJy for the $z \sim 5$ LBGs. It appears that we are finally starting to see a unifying link between the LBG and SMG populations, with LBGs becoming increasingly important contributors to the submm background as we probe down into the sub-mJy regime. Note that some classical SMGs are also LBGs, or at least have UV/optical properties consistent with LBGs (Chapman et al. 2005; Simpson et al. 2014), again indicating overlap between the populations. Further progress in unifying the SMG and LBG populations can be made

⁵ Here, we consider an SMG to be a galaxy with $S_{850} \geq 1$ mJy, since for the bulk of these sources, $S_{\text{IR}} > S_{\text{UV}}$, which makes a natural dividing line between whether the submm or UV is more important energetically.

by examining our candidate SMG-LBG counterparts, which will appear in a future paper.

We also find that the submm flux increases with LBG optical faintness and redness – that is to say, LBGs appear to be dustier and redder on average as they become more optically faint. This may seem at odds with Shapley et al. (2001), who claim that more intrinsically luminous LBGs appear to be dustier, with redder colours. But recall that our UV-SFRs have not been corrected for dust attenuation, so their intrinsic luminosities will be higher once the obscured SFRs (as measured in the submm) have been factored in. Submm flux can be taken as a proxy for the bolometric luminosity of an LBG.⁶ So in fact, we do see that the most bolometrically luminous LBGs (as traced in the submm) contain the most dust (by definition) and have redder optical colours on average, but these are not the most optically luminous LBGs. Surveys trying to detect LBGs in the submm individually should therefore focus on the optically faintest and/or reddest LBGs.

5 CONCLUSIONS

We have presented an 850 μm stacking analysis of LBG samples at $z = 3, 4$, and 5, to measure the average rest-frame far-infrared ($\sim 150\text{--}200\ \mu\text{m}$) luminosity of ‘normal’ galaxies in the first 1–2 Gyr after the big bang. We have several main findings.

(i) The average 850 μm flux density of LBGs is $\langle S_{850} \rangle \approx 0.2\text{--}0.9\ \text{mJy}$ at $z = 3\text{--}5$ (increasing with redshift), and they contribute up to 20 per cent of the submm background.

(ii) Assuming reasonable templates for the broad (8–1000 μm) infrared spectrum, the results imply that LBGs straddle the LIRG/ULIRG population, with $L_{\text{IR}} \sim 3\text{--}11 \times 10^{11}\ L_{\odot}$ (increasing with redshift). Our observed evolution of the IR luminosity density of LBGs is broadly consistent with model predictions.

(iii) We see a positive correlation between submm flux density and stellar mass: the most massive LBGs at all three epochs are the brightest at 850 μm . But we do not see evidence in our data for a continued strong evolution of the main sequence of star formation from $z = 3$ to 5.

(iv) We have determined that for a fixed mass, the optically reddest LBGs are the most submm luminous (i.e. are the most dust obscured, by definition), and that the average submm luminosity-to-UV luminosity ratio increases with mass.

These results provide an empirical baseline to motivate and guide future direct detection and mapping experiments with the Atacama Large Millimeter Array (ALMA). ALMA will achieve the sensitivities required to detect the dust continuum in typical individual LBGs in order to fully characterize these sources which are contributing to the submm background.

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REFERENCES

- Adelberger K. L., Steidel C. C., 2000, *ApJ*, 544, 218
 Adelberger K. L., Steidel C. C., Pettini M., Shapley A. E., Reddy N. A., Erb D. K., 2005, *ApJ*, 619, 697
 Baker A. J., Lutz D., Genzel R., Tacconi L. J., Lehnert M. D., 2001, *A&A*, 372, L37
 Baugh C. M., Lacey C. G., Frenk C. S., Granato G. L., Silva L., Bressan A., Benson A. J., Cole S., 2005, *MNRAS*, 356, 1191
 Bertin E., Arnouts S., 1996, *A&AS*, 317, 393
 Béthermin M., Dole H., Lagache G., Le Borgne D., Penin A., 2011, *A&A*, 529, 4
 Blain A. W., Smail I., Ivison R. J., Frayer D. T., 2002, *Phys. Rep.*, 369, 111
 Bouwens R. J., Illingworth G. D., Franx M., Ford H., 2007, *ApJ*, 670, 928
 Brammer G. B., van Dokkum P. G., Coppi P., 2008, *ApJ*, 686, 1503
 Bruzual G., 2007, in Vallenari A., Tantaló R., Portinari L., Moretti A., eds, *ASP Conf. Ser. Vol. 374, From Stars to Galaxies: Building the Pieces to Build up the Universe*. Astron. Soc. Pac., San Francisco, p. 303
 Carilli C. L., Hodge J., Walter F., Riechers D., Daddi E., Dannerbauer H., Morrison G. E., 2011, *ApJ*, 739, L33
 Chapin E. L., Berry D. S., Gibb A. G., Jenness T., Scott D., Tilanus R. P. J., Economou F., Holland W. S., 2013, *MNRAS*, 430, 2545
 Chapman S. C., Casey C. C., 2009, *MNRAS*, 398, 1615
 Chapman S. C. et al., 2000, *MNRAS*, 319, 318
 Chapman S. C., Blain A. W., Smail I., Ivison R. J., 2005, *ApJ*, 622, 772
 Chary R., Elbaz D., 2001, *ApJ*, 556, 562
 Chary R., Pope A., 2010, preprint ([arXiv:1003.1731](https://arxiv.org/abs/1003.1731))
 Coppin K. E. K. et al., 2006, *MNRAS*, 372, 1621
 Daddi E., Dannerbauer H., Elbaz D., Dickinson M., Morrison G., Stern D., Ravindranath S., 2008, *ApJ*, 673, L21
 Davies L. J. M., Bremer M. N., Stanway E. R., Mannering E., Lehnert M. D., Omont A., 2012, *MNRAS*, 425, 153
 Davies L. J. M., Bremer M. N., Stanway E. R., Lehnert M. D., 2013, *MNRAS*, 433, 2588
 Dempsey J. T. et al., 2013, *MNRAS*, 430, 2534
 Devlin M. J. et al., 2009, *Nature*, 458, 737
 Draine B. T. et al., 2007, *ApJ*, 663, 866
 Elbaz D. et al., 2007, *A&A*, 468, 33
 Fazio G. G. et al., 2004, *ApJS*, 154, 10
 Fixsen D. J., Dwek E., Mather J. C., Bennel C. L., Shafer R. A., 1998, *ApJ*, 508, 123
 Furusawa H. et al., 2008, *ApJS*, 176, 1
 Geach J. E. et al., 2013, *MNRAS*, 432, 53
 Hartley W. G. et al., 2013, *MNRAS*, 431, 3045
 Hauser M. G., Dwek E., 2001, *ARA&A*, 39, 249
 Heinis S. et al., 2013, *MNRAS*, 429, 1113
 Holland W. S. et al., 2013, *MNRAS*, 430, 2513
 Huang J.-S. et al., 2005, *ApJ*, 634, 137
 Huang J.-S. et al., 2007, *ApJ*, 660, L69
 Ivison R. J. et al., 2010, *A&A*, 518, L35
 Jenness T., Cavanagh B., Economou F., Berry D. S., 2008, in Argyle R. W., Bunclark P. S., Lewis J. R., eds, *ASP Conf. Ser. Vol. 394, Astronomical Data Analysis Software and Systems*. Astron. Soc. Pac., San Francisco, p. 565

⁶ This is so because of an effect called the negative k -correction, which applies for a galaxy of a given luminosity from $z \approx 1\text{--}8$ (Blain et al. 2002).

- Kennicutt R. C., Jr, 1998, *ARA&A*, 36, 189
- Kriek M. et al., 2006, *ApJ*, 645, 44
- Kurczynski P., Gawiser E., 2010, *AJ*, 139, 1592
- Lagache G., Puget J.-L., Dole H., 2005, *ARA&A*, 43, 727
- Laird E. S., Nandra K., Hobbs A., Steidel C. C., 2006, *MNRAS*, 373, 217
- Lawrence A. et al., 2007, *MNRAS*, 379, 1599
- Lee K., Alberts S., Atlee D., Dey A., Pope A., Jannuzi B. T., Reddy N., Brown M. J. I., 2012, *ApJ*, 758, L31
- McLure R. J., Cirasuolo M., Dunlop J. S., Foucaud S., Almaini O., 2009, *MNRAS*, 395, 2196
- Magdis G. E., Rigopoulou D., Huang J.-S., Fazio G. G., Willner S. P., Ashby M. L. N., 2008, *MNRAS*, 386, 11
- Magdis G. E., Rigopoulou D., Huang J.-S., Fazio G. G., 2010a, *MNRAS*, 401, 1521
- Magdis G. E. et al., 2010b, *ApJ*, 720, L185
- Meurer G. R., Heckman T. M., Calzetti D., 1999, *ApJ*, 521, 6
- Mortlock A. et al., 2013, *MNRAS*, 433, 1185
- Noeske K. G. et al., 2007, *ApJ*, 660, L43
- Oliver S. J. et al., 2012, *MNRAS*, 424, 1614
- Oteo I. et al., 2013, *A&A*, 554, L3
- Ouchi M. et al., 2004, *ApJ*, 611, 660
- Pascale E. et al., 2009, *ApJ*, 707, 1740
- Peacock J. et al., 2000, *MNRAS*, 318, 535
- Puget J. L., Abergel A., Bernard J. P., Boulanger F., Burton W. B., Desert F. X., Hatmann D., 1996, *A&A*, 308, L5
- Reddy N. A., Steidel C. C., 2009, *ApJ*, 692, 778
- Reddy N. A., Steidel C. C., Fadda D., Yan L., Pettini M., Shapley A. E., Erb D. K., Adelberger K. L., 2006, *ApJ*, 644, 792
- Reddy N. A. et al., 2012, *ApJ*, 744, 154
- Rieke G. H. et al., 2004, *ApJS*, 154, 25
- Rieke G. H., Alonso-Herrero A., Weiner B. J., Pérez-González P. G., Blaylock M., Donley J. L., Marcillac D., 2009, *ApJ*, 518, L25
- Rigopoulou D. et al., 2006, *ApJ*, 648, 81
- Rigopoulou D. et al., 2010, *MNRAS*, 409, L7
- Rodighiero G. et al., 2010, *A&A*, 515, A8
- Roseboom I. G. et al., 2010, *MNRAS*, 409, 48
- Salpeter E. E., 1955, *ApJ*, 121, 161
- Sawicki M., Yee H. K. C., 1998, *AJ*, 115, 1329
- Serjeant S. et al., 2010, *A&A*, 518, L7
- Shapley A. E., Steidel C. C., Adelberger K. L., Dickinson M., Giavalisco M., Pettini M., 2001, *ApJ*, 562, 95
- Shapley A. E., Steidel C. C., Erb D. K., Reddy N. A., Adelberger K. L., Pettini M., Barmby P., Huang J., 2005, *ApJ*, 626, 698
- Simpson C. J. et al., 2006, *MNRAS*, 372, 741
- Simpson C. J. et al., 2012, *MNRAS*, 421, 3060
- Simpson J. M. et al., 2014, *ApJ*, 788, 125
- Somerville R. S., Primack J. R., Faber S. M., 2001, *MNRAS*, 320, 504
- Spergel D. N. et al., 2003, *ApJS*, 148, 175
- Stanway E. R., Bremer M. N., Davies L. J. M., Lehnert M. D., 2010, *MNRAS*, 407, L94
- Stark D. P., Ellis R. S., Bunker A., Bundy K., Targett T., Benson A., Lacy M., 2009, *ApJ*, 697, 1493
- Steidel C. C., Hamilton D., 1993, *AJ*, 105, 2017
- Steidel C. C., Giavalisco M., Pettini M., Dickinson M., Adelberger K. L., 1996, *ApJ*, 462, L17
- Steidel C. C. M. P. Hunt, Shapley A. E., Adelberger K. L., Pettini M., Dickinson M., Giavalisco M., 2002, *ApJ*, 576, 653
- Swinbank A. M. et al., 2014, *MNRAS*, 438, 1267
- Ueda Y. et al., 2008, *ApJS*, 179, 124
- van Dokkum P. G. et al., 2006, *ApJ*, 638, L59
- Verma A., Lehnert M. D., Förster Schreiber N. M., Bremer M. N., Douglas L., 2007, *MNRAS*, 377, 1024
- Viero M. P. et al., 2013, *ApJ*, 779, 32
- Wardlow J. L. et al., 2011, *MNRAS*, 415, 1479
- Webb T. M. A. et al., 2003, *ApJ*, 582, 6
- Weiss A. et al., 2009, *ApJ*, 707, 1201
- Zemcov M., Blain A., Halpern M., Levenson L., 2010, *ApJ*, 721, 424
- ¹Centre for Astrophysics Research, Science & Technology Research Institute, University of Hertfordshire, College Lane, Hatfield AL10 9AB, UK
- ²The School of Physics and Astronomy, University of Nottingham, University Park, Nottingham NG7 2RD, UK
- ³SUPA, Institute for Astronomy, University of Edinburgh, Royal Observatory, Edinburgh EH9 3HJ, UK
- ⁴European Southern Observatory, Karl-Schwarzschild-Strasse 2, D-85748 Garching, Germany
- ⁵ETH Zürich, Institut für Astronomie, HIT J 11.3, Wolfgang-Pauli-Str. 27, CH-8093 Zürich, Switzerland
- ⁶Astrophysics Research Institute, Liverpool John Moores University, Twelve Quays House, Egerton Wharf, Birkenhead CH41 1LD, UK
- ⁷Institute for Computational Cosmology, Durham University, South Road, Durham DH1 3LE, UK
- ⁸Physics and Astronomy, University of Leicester, 1 University Road, Leicester LE1 7RH, UK
- ⁹Department of Physics, University of Bristol, H.H. Wills Physics Laboratory, Tyndall Avenue, Bristol BS8 1TL, UK
- ¹⁰Department of Physics and Atmospheric Science, Dalhousie University, Coburg Road, Halifax, NS B3H 1A6, Canada
- ¹¹ICRAR, The University of Western Australia, 35 Stirling Highway, Crawley, WA 6009, Australia
- ¹²Department of Physics, Virginia Tech, Blacksburg, VA 24061, USA
- ¹³Department of Physics and Astronomy, University of British Columbia, 6224 Agricultural Road, Vancouver, BC V6T 1Z1, Canada
- ¹⁴Department of Astronomy, Cornell University, Ithaca, NY 14853, USA
- ¹⁵Joint Astronomy Centre, 660 N. A'Ohoku Place, University Park, Hilo, HI 96720, USA
- ¹⁶Argelander-Institut für Astronomie, Universität Bonn, Auf dem Hügel 71, D-53121 Bonn, Germany
- ¹⁷Department of Earth and Space Sciences, Chalmers University of Technology, Onsala Space Observatory, SE-439 92 Onsala, Sweden
- ¹⁸Instituto de Física Astronomía, Universidad de Valparaíso, Avda, Gran Bretaña 111, Valparaíso, Chile
- ¹⁹Department of Physics, Oxford University, Denys Wilkinson Building, Keble Road, Oxford OX1 3RH, UK
- ²⁰RAL Space, Science & Technology Facilities Council, Rutherford Appleton Laboratory, Didcot, Oxfordshire OX11 0QX, UK
- ²¹UK Astronomy Technology Centre, Royal Observatory, Blackford Hill, Edinburgh EH9 3HJ, UK
- ²²Leiden Observatory, Leiden University, PO Box 9513, 2300 RA Leiden, the Netherlands

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