

On mirror maps for manifolds of exceptional holonomy

Andreas P. Braun,^a Suvajit Majumder^a and Alexander Otto^{b,c}

^a*Mathematical Institute, University of Oxford Woodstock Road,
Oxford, OX2 6GG, U.K.*

^b*Perimeter Institute for Theoretical Physics,
31 Caroline St N, Waterloo, ON N2L 2Y5, Canada*

^c*Department of Physics, University of Waterloo,
Waterloo, ON N2L 3G1, Canada*

E-mail: braun@maths.ox.ac.uk, majumder.suvajit95@gmail.com,
aotto@perimeterinstitute.ca

ABSTRACT: We study mirror symmetry of type II strings on manifolds with the exceptional holonomy groups G_2 and $\text{Spin}(7)$. Our central result is a construction of mirrors of $\text{Spin}(7)$ manifolds realized as generalized connected sums. In parallel to twisted connected sum G_2 manifolds, mirrors of such $\text{Spin}(7)$ manifolds can be found by applying mirror symmetry to the pair of non-compact manifolds they are glued from. To provide non-trivial checks for such geometric mirror constructions, we give a CFT analysis of mirror maps for Joyce orbifolds in several new instances for both the $\text{Spin}(7)$ and the G_2 case. For all of these models we find possible assignments of discrete torsion phases, work out the action of mirror symmetry, and confirm the consistency with the geometrical construction. A novel feature appearing in the examples we analyse is the possibility of frozen singularities.

KEYWORDS: String Duality, Conformal Field Models in String Theory, Superstring Vacua

ARXIV EPRINT: [1905.01474](https://arxiv.org/abs/1905.01474)

Contents

1	Introduction	1
2	A G_2 example	4
2.1	Smoothing of the orbifold	4
2.2	Constraints on discrete torsion	5
2.3	Cohomology of smoothings	7
2.4	Mirror symmetry	9
2.5	Realization as a TCS	11
2.6	TCS mirror map	12
3	Spin(7) examples	14
3.1	Smoothing of the orbifold	14
3.2	Discrete torsion and cohomology	15
3.2.1	Example 1	16
3.2.2	Example 2	18
3.3	Mirror symmetry	19
4	Spin(7) mirror maps for connected sums	20
4.1	Constructing Spin(7) manifolds as generalized connected sums	21
4.2	A mirror map for GCS Spin(7) manifolds	22
4.3	Examples	24
4.3.1	Example 1	24
4.3.2	Example 2	25
A	Discrete torsion and modular invariance	26
B	Discrete torsion analysis for the G_2 orbifold	27
C	Discrete torsion analysis for the Spin(7) orbifolds	31

1 Introduction

The study of string theory on manifolds with the exceptional holonomy groups G_2 and Spin(7) from the worldsheet perspective goes back to [1]. In particular, [1] (see also [2]) determined the extension of the worldsheet superconformal algebra for strings propagating on manifolds of exceptional holonomy and pointed out that type II string theories propagating on different G_2 (or Spin(7) manifolds) manifolds may result in equivalent physical theories in a phenomenon called G_2 (or Spin(7)) mirror symmetry. A necessary condition for any pair of manifolds M and M^\vee to be mirror is that the dimensions of the spaces of

exactly marginal operators of the extended $\mathcal{N} = (1, 1)$ worldsheet CFTs agree. As shown in [1], the number of exactly marginal operators simply equals the number of geometric moduli together with the degrees of freedom associated with the B-field. In the G_2 context, this implies the equality of the sum of the second and third Betti numbers for mirrors:

$$b^2(M) + b^3(M) = b^2(M^\vee) + b^3(M^\vee). \tag{1.1}$$

Likewise, a pair of Spin(7) mirrors M, M^\vee must satisfy

$$b^2(M) + b_-^4(M) + 1 = b^2(M^\vee) + b_-^4(M^\vee) + 1, \tag{1.2}$$

where $b_-^4(M)$ denotes the dimension of the space of anti self-dual four-forms.

This was made explicit for the first time in the context of Joyce orbifolds [3–5] in [6, 7]. In parallel to the well-studied case of mirror symmetry for type II strings on Calabi-Yau manifolds, where the mirror map can be understood from T-dualities along a calibrated T^3 fibration [8], mirror maps for type II strings on G_2 and Spin(7) manifolds were shown to arise from T-dualities along calibrated T^3 or T^4 fibrations for some of Joyce’s examples in [6, 7, 9, 10], see also [11].

More recently, [12–14] have given a construction of G_2 manifolds as twisted connected sums (TCS) by gluing appropriate pairs of asymptotically cylindrical (acyl) Calabi-Yau threefolds X_\pm (times a circle \mathbb{S}_\pm^1 on each side). Mirror maps for TCS G_2 manifolds were found in [15, 16], where it was shown that applying a mirror map to both of the acyl Calabi-Yau threefolds, together with T-dualities on the product circles, leads to another TCS G_2 manifold which satisfies (1.1). This map can be described as being the result of performing four T-dualities along a calibrated T^4 fibration, so that it maps type IIA (IIB) strings to IIA (IIB) strings. The T^4 fibre in question is understood as the product of the Strominger-Yau-Zaslow (SYZ) fibres of the acyl Calabi-Yau threefolds times the circles \mathbb{S}_\pm^1 .

Interestingly, TCS G_2 manifolds allow a second class of mirror maps satisfying (1.1), in which only one of the two acyl Calabi-Yau threefolds X_\pm is exchanged for its mirror [16]. In case one of the acyl Calabi-Yau threefolds carries an elliptic fibration, this mirror map can likewise be understood from three T-dualities along a T^3 fibre,¹ i.e. this duality must map type IIA strings to type IIB or vice versa. For compact Calabi-Yau manifolds which are hypersurfaces or complete intersections in toric varieties, mirror families have an elegant combinatorial construction using pairs of reflexive polytopes [18]. As shown in [15, 19], a completely analogous construction exists for acyl Calabi-Yau threefolds in terms of pairs of dual tops, which makes it possible to give large numbers of concrete examples of G_2 mirrors. An intriguing feature of both classes of mirror maps is that they sometimes map smooth geometries to singular ones [16]. This is analogous to mirror symmetry in the context of K3 surfaces [20], where the presence of the B-field prevents the occurrence of extra massless states.

Calibrated torus fibrations are not just interesting in the study of mirror symmetry of Calabi-Yau, G_2 or Spin(7) manifolds, but also feature in the duality between M-Theory

¹As discussed in [17], it is not expected on general grounds for G_2 manifolds to have such fibrations, so it might be better to speak of such G_2 manifolds as admitting a limit in their moduli space in which an associative T^3 collapses.

and heterotic string theory. The duality in seven dimensions between M-Theory on K3 surfaces and heterotic string theory on a three-torus T^3 with a flat connection can be used fibrewise to find examples of lower-dimensional dualities. This was exploited for the SYZ fibration of Calabi-Yau threefolds to find examples of heterotic duals of M-Theory on TCS G_2 manifolds in [21] (see also [22–24] for explorations of M-Theory on TCS G_2 manifolds). Furthermore, T^4 fibrations also play a crucial role for the determination of the superpotential of M-Theory on TCS G_2 manifolds. As argued in [25, 26], there is a large class of associative submanifolds of TCS G_2 manifolds which appear as sections of precisely the coassociative T^4 fibration relevant for mirror maps.

Applying the M-Theory duality map to heterotic strings on TCS G_2 manifolds yields M-Theory duals on Spin(7) manifolds with a decomposition similar to TCS, called generalized connected sum (GCS) in [27]. GCS Spin(7) manifolds are glued from two non-compact manifolds with the holonomy groups SU(4) and G_2 , so that it is tempting to exploit this structure to construct candidates of mirror manifolds by using mirrors for the building blocks. This is an analogue of the strategy used in [15, 16] for TCS G_2 manifolds and we will follow a similar path to define Spin(7) mirrors in the present work, and show that they indeed satisfy (1.2).

Given geometric constructions for mirrors of G_2 and Spin(7) manifolds, it becomes an interesting question if these can be recovered using worldsheet methods. This has been accomplished for a few of the examples of Joyce. These are resolutions of orbifolds, so that they can be treated from first principles in string theory [6, 9]. Furthermore, these geometries can also be decomposed as twisted connected sums. As shown in [16], the two complimentary approaches result in the same mirror maps and identify the same T^3 and T^4 fibrations. Furthermore, it has subsequently been shown [28] that the G_2 mirror maps of [15, 16] are associated with non-trivial automorphisms of the extended superconformal algebra of the worldsheet theory for TCS G_2 manifolds.

It is a central motivation of the present work to enlarge the class of models where a geometrical construction of mirrors for G_2 and Spin(7) manifolds can be compared with results obtained from the worldsheet, and we complete this task for several new examples.

Treating new examples of orbifolds, and linking the associated mirror maps with a geometry requires several steps to be completed. Crucially, the definition of string theory on the orbifolds we are considering involves an assignment of discrete torsion phases [29]. Possible assignments of discrete torsion phases can be constrained by the requirement of modular invariance for the partition function. A short introduction to how such constraints arise and can be analysed is presented in appendix A. Different assignments of discrete torsion will in turn lead to a different spectrum of RR ground states, which via the map to cohomology [30] signals the correspondence to topologically different smoothings of the orbifold geometry. Mirror symmetry can change the discrete torsion phases, so that it also associates CFTs on different smooth geometries. The action of mirror maps can be found by providing a free-field realization of the extended superconformal algebra, and finding automorphisms of this algebra which are induced by a sequence of T-dualities. Finally, to establish a link to the TCS mirror maps the orbifolds in question must be described as twisted connected sums.

In the G_2 examples studied in [9], the possible assignments of discrete torsion precisely match the different resolutions found in [4] (‘example 4’). Furthermore, the description of such models as TCS G_2 manifolds in [16] allowed for a straightforward determination of torsion in the homology group $H^3(M_k, \mathbb{Z})$, which precisely matches with the discrete torsion phases in the orbifold string theory, as expected from [31, 32].

In section 2, we will complete these tasks for a set of models which are free quotients of the orbifolds considered in [9, 16], they have been first presented as ‘example 5’ in [4]. As we will see, these orbifolds are an example the ‘extra-twisted’ connected sums described in [33]. Curiously, not all resolutions which have been constructed in [4] are realized by the set of consistent assignments of discrete torsion phases.

In section 3, we consider Joyce orbifolds which can be smoothed to manifolds with holonomy $\text{Spin}(7)$. Some Joyce orbifolds and their mirror maps have been previously from the worldsheet perspective in [6, 10]. As only submanifolds of real dimension four can be calibrated for $\text{Spin}(7)$ manifolds, such mirror maps must be associated with four T-dualities along a calibrated T^4 fibration. In section 3 of the present work we study two such examples, which first appeared as ‘example 1’ and ‘example 2’ in [5], as well the action of mirror symmetry on such geometries.

Finally, section 4 presents a general exposition of how the GCS construction can be used to define mirror maps for $\text{Spin}(7)$ manifolds. While our construction is motivated by the identification of a fibration by T^4 , the resulting check of (1.2) for the mirror geometries holds independently. We then proceed to describe how the examples studied in section 3 are decomposed as a GCS and verify that our geometric construction of GCS $\text{Spin}(7)$ mirrors precisely agrees with the worldsheet results.

The appendices contain a brief introduction to discrete torsion phases and modular invariance for strings on orbifolds, as well technical details of the examples we are treating.

2 A G_2 example

In this section we consider the following example from [4] (‘example 5’). It is based on a quotient of T^7 under the group $\Gamma = \mathbb{Z}_2^4$, with an action on the coordinates

	X^1	X^2	X^3	X^4	X^5	X^6	X^7
α	+	+	+	-	-	-	-
β	+	-	-	+	+	$-\frac{1}{2}$	-
γ	-	+	-	+	-	+	$-\frac{1}{2}$
σ_2	+	$\frac{1}{2}$	+	$\frac{1}{2}$	+	+	+

(2.1)

Here $X^i \sim X^i + 1$ giving T^7 , and $-$ indicates that the corresponding coordinate is sent to minus itself, while $\frac{1}{2}$ indicates a shift $X \rightarrow X + \frac{1}{2}$, and $-\frac{1}{2}$ is a shorthand for $X \rightarrow -X + \frac{1}{2}$.

2.1 Smoothing of the orbifold

Let us first analyse the topological properties of the smooth limit(s) obtained from this orbifold by resolving the singularities as in [4] (see also [17]). In order to do that, we need to first evaluate the fixed point set under the action of the orbifold group.

The only group elements of Γ which act non-freely and hence give rise to singularities are α , β , and γ . Each of these three elements fixes 16 T^3 s, which are then further identified under the action of the rest of the orbifold group. In case of the α sector, the action of $\langle \beta, \gamma, \sigma_2 \rangle$ is free, resulting in two orbits of 8 T^3 each. For the β sector, we have the same result as the α sector as these two sectors are isomorphic, up to the permutation of indices and fixed point labels. The analysis is a bit different for the γ case. Here we have a further \mathbb{Z}_2 identification under the action of $\alpha\beta$. This leads to 8 orbits consisting of 2 T^3/\mathbb{Z}_2 each. The action of σ_2 just reduces the fundamental domain for X^2 and X^4 (two of the extended directions in the γ fixed T^3 s). Hence, the γ action has a singular set consisting of 8 T^3/\mathbb{Z}_2 .

	singular set	elements in orbit under Γ	singular set in quotient
α	16 T^3	8 T^3	2 T^3
β	16 T^3	8 T^3	2 T^3
γ	16 T^3	2 T^3/\mathbb{Z}_2	8 T^3/\mathbb{Z}_2

The net contribution to the singular set is obtained by adding up the separate contributions from the different orbifold actions, giving us a total of 4 T^3 , and 8 T^3/\mathbb{Z}_2 . With the singular set at hand, we are now ready to compute the Betti numbers for the different smooth limits obtained by resolving the singularities present in our orbifold. First of all, the Betti numbers for the orbifold are given by $(b^2, b^3)(T^7/\Gamma) = (0, 7)$. Each of the singularities is locally modelled on $T^3 \times \mathbb{C}^2/\{\pm 1\}$ or $(T^3 \times \mathbb{C}^2/\{\pm 1\})/\langle \alpha\beta \rangle$. The contribution to the Betti numbers from resolving these singularities is:

$$\begin{aligned} \delta(b^2, b^3)(T^3) &= (1, 3) \\ \delta(b^2, b^3)(T^3/\mathbb{Z}_2) &= (1, 1) \text{ or } (0, 2) \end{aligned} \tag{2.2}$$

The two choices for the latter case come from different possible smoothings. Denoting the compact smooth G_2 manifold obtained by making the first choice k times by \hat{M}_k we find

$$\begin{aligned} b^2(\hat{M}_k) &= 4 + k \\ b^3(\hat{M}_k) &= 35 - k \end{aligned} \tag{2.3}$$

for $k = 0 \dots 8$.

2.2 Constraints on discrete torsion

Let us redo the analysis by studying the string theory living on the orbifold T^7/\mathbb{Z}_2^4 . For such string theories, we have the additional degree of freedom to switch on discrete torsion phases consistent with modular invariance constraints [29]. We then exploit the isomorphism between the Ramond-Ramond sector ground states and the target space cohomology to find the Betti numbers for the resulting resolution(s).

We shall begin by studying the twisted sectors corresponding to various elements of the orbifold group. If the orbifold element has fixed points in its action on the parent manifold, the corresponding sector can then be further decomposed into sub-sectors localised at those fixed points. In general, one would start by figuring out which of the sectors can actually have discrete torsion signs. For this, we need to write down the representation matrices for the orbifold elements in a basis of highest weight states in a particular sector. However,

as our main aim is to obtain the cohomology of the resulting orbifold smoothings, we can get away by working with a few of the twisted sectors only: ones that contribute to the ground state spectrum (and hence the cohomology), and ones that are needed to constrain the discrete torsion signs for the ground state contributing sectors.

In our case, the ground state contributions come from the α , β , and γ sectors, as all the remaining group elements act on at least one direction as $+\frac{1}{2}$, implying that the oscillator modes are half-integer in that direction. Aside from these sectors, we need to consider $\alpha\beta$, σ_2 , and $\sigma_2\alpha\beta$ sectors in order to figure out the modular trace constraints in the α , β , and γ sectors. Let us then find which of the above mentioned sectors have discrete torsion signs. The details of this analysis can be found in appendix B.

α - and β -sector: there are no discrete torsion choices in these sectors.

γ -sector: in the γ sector, we have the first instance of discrete torsion. We find that there are two orbifold generators with possible discrete torsion signs in their irreducible representations: $\alpha\beta$ and σ_2 . The corresponding phases are denoted by $\epsilon_{f_\gamma}(\alpha\beta)$, $\epsilon_{f_\gamma}(\sigma_2)$ respectively, where $f_\gamma = 1, 2, \dots, 8$ labels the different irreducible representations/orbits.

As reviewed in appendix A, one can constrain the discrete torsion signs by implementing modular invariance of the partition functions. For that, we utilize the orbifold elements with non-zero traces in the above representations and probe the twisted sector corresponding to those elements. These are given by the group elements $\alpha\beta$, σ_2 , and $\sigma_2\alpha\beta$.

$\alpha\beta$ -sector: in this sector, we have discrete torsion signs appearing in the representation matrix corresponding to the γ generator:

$$\gamma|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \begin{pmatrix} \epsilon_{f_{\alpha\beta}}^1(\gamma)\mathbb{I}_{4\times 4} & 0 \\ 0 & \epsilon_{f_{\alpha\beta}}^2(\gamma)\mathbb{I}_{4\times 4} \end{pmatrix} \quad (2.4)$$

Invariance of the partition function under the S transformation implies

$$\sum_{f_{\alpha\beta}, i} \epsilon_{f_{\alpha\beta}}^i(\gamma) = \sum_{f_\gamma=1}^8 \epsilon_{f_\gamma}(\alpha\beta) \quad (2.5)$$

where $f_{\alpha\beta} = 1, 2, 3, 4$ labels the irreducible representations, and $i=1,2$ labels the two signs in a given irreducible representation.

σ_2 -sector: as shown in appendix B, discrete torsion arises in the representation of the element γ , $\gamma|_{\mathcal{H}_{\sigma_2}} = \epsilon_{\sigma_2}(\gamma)\mathbb{I}_{4\times 4}$. We can now use the S transform relations to connect this discrete torsion sign with $\epsilon_{f_\gamma}(\sigma_2)$ from the γ twisted sector as follows:

$$8\epsilon_{\sigma_2}(\gamma) = \sum_{f_\gamma=1}^8 \epsilon_{f_\gamma}(\sigma_2) \quad (2.6)$$

$\sigma_2\alpha\beta$ -sector: analogous to the σ_2 sector, discrete torsion arises in the representation of the element $\gamma|_{\mathcal{H}_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}} = \epsilon_{f_{\sigma_2\alpha\beta}}(\gamma)\mathbb{I}_{8\times 8}$, as is expected from modular constraint requirements.

Using modular invariance of partition function under the S-transformation gives

$$2 \sum_{f=1}^4 \epsilon_{f_{\sigma_2\alpha\beta}}(\gamma) = \sum_{f=1}^8 \epsilon_{f_\gamma}(\sigma_2)\epsilon_{f_\gamma}(\alpha\beta) \quad (2.7)$$

2.3 Cohomology of smoothings

We are ready to derive the cohomology of the possible resolutions of the orbifold by exploiting its isomorphism with the RR ground states. Let us first consider the untwisted sector, \mathcal{H}_e . We have Majorana-Weyl spinors ψ^i corresponding to the 7 bosonic coordinates X^i , $i = 1, \dots, 7$. We can generate the RR ground states by acting on the vacuum $|0\rangle$ with the creation operators built out of the zero modes, $\psi_+^i = (\psi_0^i + i\psi_0^i)/2$. The orbifold invariant set of RR ground states is given by:

$$|0\rangle; \quad \psi_+^i \psi_+^j \psi_+^k |0\rangle; \quad \psi_+^a \psi_+^b \psi_+^c \psi_+^d |0\rangle; \quad \psi_+^1 \dots \psi_+^7 |0\rangle \quad (2.8)$$

with the following the triples and the 4-tuples of indices:

$$(i, j, k) \in \{(1, 2, 3), (1, 4, 5), (1, 6, 7), (2, 4, 6), (2, 5, 7), (3, 5, 7), (3, 4, 6)\} \quad (2.9)$$

$$(a, b, c, d) \in \{(4, 5, 6, 7), (2, 3, 6, 7), (2, 3, 4, 5), (1, 3, 5, 7), (1, 3, 4, 6), (1, 2, 4, 6), (1, 2, 5, 7)\} \quad (2.10)$$

Now the isomorphism between RR ground states and the target space cohomology allows us to use the following identification:

$$\psi_+^{i_1} \dots \psi_+^{i_n} |0\rangle \simeq dX^{i_1} \wedge \dots \wedge dX^{i_n} \quad (2.11)$$

As such, we can use the list of invariant states from (2.8) to get the Betti number contribution from the untwisted sector:

$$\begin{aligned} b_u^0 &= b_u^7 = 1 \\ b_u^3 &= b_u^4 = 7 \end{aligned} \quad (2.12)$$

For G_2 manifolds, the only unfixed, non-trivial Betti numbers are the second and third ones. Now on, we will only mention contributions to those.

Aside from the untwisted sector, the smoothings for the orbifold will also get contribution to the Betti numbers from the twisted sectors with RR ground states, i.e. α , β , and γ . Now, we shall evaluate such twisted sector contributions.

Let us look at the contribution from α sector. Firstly, following in the same line of argument as in [9], we should identify the vacuum state in this sector with the exceptional divisor Σ resolving the particular singularity (a 2-form in this case). Then, we list the invariant states in the α sector and use a similar isomorphism statement as in the untwisted sector case (in (2.11)) to get the Betti numbers contribution.

The list of Γ invariant RR ground states in the α twisted sector is built on the highest weight states $|0\rangle_\alpha^{f_\alpha}$ with vanishing momentum and winding modes, by acting with the creation operators built from the Majorana Weyl fermions as in the untwisted sector. The only difference is that we only have zero modes along X^1 , X^2 , and X^3 . Here the label $f_\alpha = 1, 2$ enumerates the two irreducible representations corresponding to $X^5 = 0, 1/2$ as in appendix B. They are given by:

$$|0\rangle_\alpha^{f_\alpha}; \quad \psi_+^i |0\rangle_\alpha^{f_\alpha}; \quad \psi_+^{i_1} \psi_+^{i_2} |0\rangle_\alpha^{f_\alpha}; \quad \psi_+^{k_1} \psi_+^{k_2} \psi_+^{k_3} |0\rangle_\alpha^{f_\alpha} \quad (2.13)$$

where $i \in \{1, 2, 3\}$; $(i_1, i_2) \in \{(1, 2), (1, 3), (2, 3)\}$; $(k_1, k_2, k_3) \in \{(1, 2, 3)\}$. The isomorphism statement can be again used to get the following identification:

$$\psi_+^{i_1} \dots \psi_+^{i_n} |0\rangle_\alpha^{f_\alpha} \simeq dX^{i_1} \wedge \dots \wedge dX^{i_n} \wedge \Sigma_{f_\alpha} \quad (2.14)$$

where Σ_{f_α} represents the exceptional divisor arising due to the resolution of the singularity corresponding to the orbit (\sim irreducible representation) labelled by f_α . Now, we can read off the Betti numbers δb_α^i from the above list as:

$$\begin{aligned} \delta b_\alpha^2 &= 2 \cdot 1 = 2 \\ \delta b_\alpha^3 &= 2 \cdot 3 = 6 \end{aligned} \quad (2.15)$$

where the factor of 2 comes from the index f_α . As the β sector is isomorphic to the α sector, we find the same contributions there.

For the γ -sector, things are different as there are discrete torsion signs present. Firstly, we need to solve for the trace relations in (2.5), (2.6), (2.7), that constrain the discrete torsion phases arising in this sector. There are two distinct cases corresponding to the sign of ϵ_{σ_2} . For $\epsilon_{\sigma_2} = 1$, (2.6) implies that the only possible solution is $\epsilon_{f_\gamma}(\sigma_2) = 1$ for all f_γ . Also from (2.7) we get:

$$2 \sum_{f_{\sigma_2 \alpha \beta} = 1}^4 \epsilon_{f_{\sigma_2 \alpha \beta}}(\gamma) = \sum_{f_\gamma = 1}^8 \epsilon_{f_\gamma}(\alpha \beta). \quad (2.16)$$

Note that we are forced by this relation to have an even number of positive discrete torsion signs $\epsilon_{f_\gamma}(\alpha \beta)$. Thus, we can re-label the fixed point indices such that:

$$\epsilon_{f_\gamma}(\alpha \beta) = \epsilon_{8-f_\gamma}(\alpha \beta), \quad f_\gamma = 1, 2, 3, 4 \quad (2.17)$$

Now let us list the orbifold invariant states as before. For $\epsilon_{f_\gamma}(\alpha \beta) = \epsilon_{f_\gamma} = 1$, they are:

$$|0\rangle_\gamma^{f_\gamma}; \quad \psi_+^6 |0\rangle_\gamma^{f_\gamma}; \quad \psi_+^2 \psi_+^4 |0\rangle_\gamma^{f_\gamma}; \quad \psi_+^2 \psi_+^4 \psi_+^6 |0\rangle_\gamma^{f_\gamma} \quad (2.18)$$

where $|0\rangle_\gamma^{f_\gamma}$ is the highest weight state representing the irreducible basis for the different orbits indexed by f_γ . On the other hand, for $\epsilon_{f_\gamma}(\alpha \beta) = -\epsilon_{f_\gamma} = -1$, the orbifold invariant states are:

$$\psi_+^2 |0\rangle_\gamma^{f_\gamma}; \quad \psi_+^4 |0\rangle_\gamma^{f_\gamma}; \quad \psi_+^2 \psi_+^6 |0\rangle_\gamma^{f_\gamma}; \quad \psi_+^4 \psi_+^6 |0\rangle_\gamma^{f_\gamma}; \quad (2.19)$$

For the Betti number contributions, we need the state-cohomology identification statement as before. In this case, the highest weight state is again mapped to a 2-form for either case of the discrete torsion sign of $\epsilon_{f_\gamma}(\alpha \beta)$. This is because the exceptional divisor in both cases corresponds to a 2-form. Using the same identification as in the α and β sectors, we can read off the Betti numbers from the list of invariant states in (2.18), and (2.19) as:

$$\begin{aligned} \delta b_\gamma^2 &= 1; \quad \delta b_\gamma^3 = 1, \quad \text{iff } \epsilon_{f_\gamma}(\alpha \beta) = 1 \\ \delta b_\gamma^2 &= 0; \quad \delta b_\gamma^3 = 2, \quad \text{iff } \epsilon_{f_\gamma}(\alpha \beta) = -1 \end{aligned} \quad (2.20)$$

Thus, the resulting Betti number contribution for $2l$ positive signs of $\epsilon_{f_\gamma}(\alpha \beta)$ is:

$$\begin{aligned} \delta b_\gamma^2 &= 2l \cdot 1 + (8 - 2l) \cdot 0 = 2l \\ \delta b_\gamma^3 &= 2l \cdot 1 + (8 - 2l) \cdot 2 = 16 - 2l \end{aligned} \quad (2.21)$$

for $l = 0 \dots 4$. These are all the contributions from the twisted sectors as the rest of the sectors all have at least one extended direction with half-integer modes. Let the resultant smoothing for $2l$ positive signs of $\epsilon_{f_\gamma}(\alpha\beta)$ be called \hat{M}_{2l} . The net Betti number for \hat{M}_{2l} can be obtained by adding up the contributions from the three twisted sectors α, β, γ along with that coming from the untwisted sector. Using (2.12), (2.15) and (2.21), we get:

$$\begin{aligned} b^2(\hat{M}_{2l}) &= b_u^2 + \delta b_\alpha^2 + \delta b_\beta^2 + \delta b_\gamma^2 = 0 + 2 + 2 + 2l = & 4 + 2l \\ b^3(\hat{M}_{2l}) &= b_u^3 + \delta b_\alpha^3 + \delta b_\beta^3 + \delta b_\gamma^3 = 7 + 6 + 6 + (16 - 2l) = & 35 - 2l \end{aligned} \tag{2.22}$$

for $l = 0 \dots 4$. Although every single one of these models corresponds to one of the resolutions discussed in section 2.1, not all of the possible geometries are realized. This ultimately stems from the constraint that the number of discrete torsion sign in (2.17) must be even. It would be interesting to understand this mismatch better, but as our main interest is in the action of mirror maps on these models we leave such an investigation to future work.

Now let us move onto the other case where we have $\epsilon_{\sigma_2} = -1$. Here, we don't get any states from the γ -twisted sector as no invariant states can be constructed with $\epsilon_{f_\gamma}(\sigma_2) = -1$. The Betti numbers for the partially smoothed solution, say \hat{N} , is then obtained by omitting the contribution from the γ sector in our previous computation in (2.22):

$$\begin{aligned} b^2(\hat{N}) &= b_u^2 + \delta b_\alpha^2 + \delta b_\beta^2 = 0 + 2 + 2 = 4 \\ b^3(\hat{N}) &= b_u^3 + \delta b_\alpha^3 + \delta b_\beta^3 = 7 + 6 + 6 = 19 \end{aligned} \tag{2.23}$$

What we have here is a scenario where the orbifold singularities could only partially be smoothed, but some are frozen. In particular, all of the singularities located at the γ fixed points must be left intact. The freezing of singularities is a well-known phenomenon for strings on orbifolds in the presence of discrete torsion [34], but it is, to the authors knowledge, the first time it has been observed for G_2 orbifolds.

2.4 Mirror symmetry

Let us now have a look at the realisation of mirror symmetry in our G_2 example through T-duality transformations on multiple suitably chosen coordinates.

The extended chiral algebra for a string moving on a compact G_2 manifold consists of a $\mathcal{N} = 1$ superconformal algebra generated by the bosonic stress tensor T along with its fermionic counterpart G , extended by currents (ϕ, X) of conformal dimensions $(h_\phi, h_X) = (3/2, 2)$ and their superpartners K, M . Their free field representation is

$$\begin{aligned} T_{G_2} &= 1/2 \sum_{j=1}^7 : \partial X^j \partial X^j : - 1/2 \sum_{j=1}^7 : \psi^j \partial \psi^j : \\ G_{G_2} &= \sum_{j=1}^7 \psi^j \partial X^j \\ X_{G_2} &= -\psi^4 \psi^5 \psi^6 \psi^7 - \psi^2 \psi^3 \psi^6 \psi^7 - \psi^2 \psi^3 \psi^4 \psi^5 \\ &\quad - \psi^1 \psi^3 \psi^5 \psi^7 + \psi^1 \psi^3 \psi^4 \psi^6 + \psi^1 \psi^2 \psi^5 \psi^6 + \psi^1 \psi^2 \psi^4 \psi^7 \\ \Phi_{G_2} &= \psi^1 \psi^2 \psi^3 + \psi^1 \psi^4 \psi^5 + \psi^1 \psi^6 \psi^7 + \psi^2 \psi^4 \psi^6 - \psi^2 \psi^5 \psi^7 - \psi^3 \psi^4 \psi^7 - \psi^3 \psi^5 \psi^6 \\ M_{G_2} &= [G, X] \\ K_{G_2} &= \{G, \Phi\} \end{aligned} \tag{2.24}$$

A mirror automorphism of this algebra is given by:

	T_{G_2}	G_{G_2}	Φ_{G_2}	X_{G_2}	K_{G_2}	M_{G_2}	
mirror $_{G_2}$	+	+	-	+	-	+	(2.25)

Defining two sets of triples of coordinate indices as follows:

$$\begin{aligned} \mathcal{I}_3^+ &= \{(1, 6, 7), (2, 4, 6), (3, 5, 6)\} \\ \mathcal{I}_3^- &= \{(1, 2, 3), (1, 4, 5), (2, 5, 7), (3, 4, 7)\} \end{aligned} \quad (2.26)$$

one can check that the elements \mathcal{I}_3^\pm are precisely those T-duality triples that generate the mirror automorphism of eq. (2.25) on the right-movers (left chiral algebra stays invariant). As the present model is a quotient of the one considered in [9], it should come as no surprise that these are the same triples which were found there (after an appropriate relabelling of coordinates).

The resultant action of these composite T-dualities on the discrete torsion phases can be understood through the following line of argument. If we focus on the representation of the $\alpha\beta$ generator in the γ -twisted sector (knowing that it has discrete torsion phases $\epsilon_{f_\gamma}(\alpha\beta)$ arising in its representation), we observe a need for flipping all or none of the discrete torsion signs upon action of these transformations. On the other hand, the signs corresponding to the σ_2 generator, i.e. $\epsilon_{f_\gamma}(\sigma_2)$ do not change. In the γ -twisted sector, the RR zero modes surviving are labelled the indices (2,4,6). Now $\alpha\beta$ flips the signs of coordinate labels (2,4) while σ_2 does not flip any of the three labels. As a result, we can write the elements in terms of the RR zero modes:

$$\begin{aligned} \alpha\beta &= \psi_0^2 \psi_0^4 \tilde{\psi}_0^2 \tilde{\psi}_0^4 \epsilon_{f_\gamma}(\alpha\beta) \\ \sigma_2 &= \epsilon_{f_\gamma}(\sigma_2) \end{aligned} \quad (2.27)$$

The above expressions imply that under the set of T-transformations in \mathcal{I}_3^+ , $\epsilon_{f_\gamma}(\alpha\beta)$ should remain the same; while for \mathcal{I}_3^- , $\epsilon_{f_\gamma}(\alpha\beta)$ should flip signs irrespective of the fixed point label f . On the other hand, the σ_2 representation matrix remaining invariant under the triples of T-dualities implies that $\epsilon_{f_\gamma}(\sigma_2)$ has to remain invariant as well. This implies that

$$\begin{aligned} \mathcal{I}_3^+ : \hat{M}_{2l} &\rightarrow \hat{M}_{2l} \\ \mathcal{I}_3^- : \hat{M}_{2l} &\rightarrow \hat{M}_{8-2l} \end{aligned} \quad (2.28)$$

Furthermore, the case with frozen singularities, represented by (2.23), should be considered self-mirror. Note that all of these maps take type IIA strings to type IIB strings and vice versa because of the odd number of fermionic modes being T-dualized.

We can combine any two of the above transformations to get a new transformation which T-dualizes along 4 of the 7 coordinates. These transformations split into \mathcal{I}_4^+ , and \mathcal{I}_4^- , with \mathcal{I}_4^+ keeping the discrete torsion signs $\epsilon_{f_\gamma}(\alpha\beta)$ intact and \mathcal{I}_4^- inverting all of them

$$\begin{aligned} \mathcal{I}_4^+ &= \{(1, 2, 4, 7), (1, 3, 5, 7), (2, 3, 4, 5)\} \\ \mathcal{I}_4^- &= \{(2, 3, 6, 7), (4, 5, 6, 7), (1, 2, 5, 6), (1, 3, 4, 6)\} \end{aligned} \quad (2.29)$$

In particular, these act as

$$\begin{aligned} \mathcal{I}_4^+ : \hat{M}_{2l} &\rightarrow \hat{M}_{2l} \\ \mathcal{I}_4^- : \hat{M}_{2l} &\rightarrow \hat{M}_{8-2l} \end{aligned} \quad (2.30)$$

Now as there are even number of fermionic modes, the mirror maps corresponding to these transformations do not alter the sign of GSO projection, mapping type IIA and IIB strings to IIA and IIB respectively.

2.5 Realization as a TCS

Let us now discuss how the orbifold treated above and its smoothings \hat{M}_k can be described as a extra twisted connected sum. Twisted connected sum G_2 manifolds are constructed as [12, 14]

$$M = (X_+ \times \mathbb{S}_+^1) \# (X_- \times \mathbb{S}_-^1), \quad (2.31)$$

for a pair of asymptotically cylindrical Calabi-Yau threefolds X_\pm which enjoy fibration by K3 surfaces S_\pm . In their asymptotic regions X_\pm are approximated (metrically) by the product of S_\pm and a cylinder, i.e. $S_\pm \times \mathbb{S}_\mp^1 \times I$ for an interval I . The gluing to M is then done by identifying those asymptotic regions by a diffeomorphism which induces an appropriate hyper-Kähler rotation on the K3 surfaces S_\pm .

The example \hat{M}_k discussed above is not a TCS G_2 manifold. However, it can be constructed as a quotient of another G_2 manifold M_k by the free involution generated by σ_2 . The realization of M_k as a TCS has been discussed in detail in [16]. The upshot of this analysis is that the K3 fibres S_\pm of both acyl Calabi-Yau threefolds X_\pm are given as a (smoothing) of T^4/γ with the coordinates (X^1, X^3, X^5, X^7) . The acyl Calabi-Yau threefold X_+ can be described as $T^5 \times \mathbb{R}/\langle \gamma, \alpha \rangle$, with X^6 parametrizing the non-compact direction and $(X^1, X^3, X^5, X^7, X^4)$ parametrizing the T^5 . Hence the base of the K3 fibration on X_+ has coordinates (X^4, X^6) and the coordinate along \mathbb{S}_+^1 is X^2 . Likewise, X_- can be described as $T^5 \times \mathbb{R}/\langle \gamma, \beta \rangle$, also with X^6 parametrizing the non-compact direction. Now, however the base of the K3 fibration on X_- has coordinates (X^2, X^6) and the coordinate along \mathbb{S}_-^1 is X^4 , as appropriate for gluing these two acyl CYs to a TCS G_2 manifold.

The freely acting involution σ_2 hence only acts on the coordinates X^2, X^4, X^6 along the base S^3 of the K3 fibration apparent in the TCS decomposition of M_k . In particular, it acts by shifting the coordinates on both of the S^1 factors in the decomposition of S^3 appearing in the TCS construction, where it is glued from two solid tori.² The further quotient by σ_2 turns this into an example of an *extra-twisted connected sum* as defined

²This is in fact the same freely-acting involution of S^3 which is naturally found by representing S^3 as

$$|z_1|^2 + |z_2|^2 = 1 \quad (2.32)$$

and acting with

$$\sigma_2 : \begin{aligned} z_1 &\rightarrow -z_1 \\ z_2 &\rightarrow -z_2 \end{aligned} \quad (2.33)$$

In this presentation, the decomposition into two solid tori can be seen by solving (2.32) using the parametrization

$$\begin{aligned} z_1 &= e^{i\phi} \cos(\eta) \\ z_2 &= e^{i\psi} \sin(\eta) \end{aligned} \quad (2.34)$$

in [33]. This construction differs from the usual TCS construction [12] in two regards: first of all, the G_2 manifold M is glued from two pieces

$$M = V_+ \#_{\vartheta} V_-, \quad (2.36)$$

where V_{\pm} are free quotients of $X_{\pm} \times \mathbb{S}_{\pm}^1$ by \mathbb{Z}_2^s for acyl Calabi-Yau threefolds X_{\pm} . In the neck region in which X_{\pm} asymptote to $S_{\pm} \times \mathbb{S}_{\mp}^1 \times I$ for K3 surfaces S_{\pm} , the quotient must purely act by shifts on the \mathbb{S}^1 factors, so that

$$V_{\pm} \setminus \kappa_{\pm} = S_{\pm} \times I \times (\mathbb{S}_+^1 \times \mathbb{S}_-^1) / \mathbb{Z}_2^s \quad (2.37)$$

for compact subsets κ_{\pm} . Note that $(\mathbb{S}_+^1 \times \mathbb{S}_-^1) / \mathbb{Z}_2^s$ is still a two-torus. The second difference is that the gluing now involves an angle ϑ with which these tori are identified, together with an appropriate altering of the hyper-Kähler rotation acting on the K3 surfaces S_{\pm} to keep the canonically defined G_2 forms Φ_{\pm} invariant. In the example discussed here, the ‘trivial’ choice $\vartheta = \pi/2$ appears, which means that the two \mathbb{S}^1 s are simply swapped, as in the standard TCS construction.

2.6 TCS mirror map

We can now describe a mirror map in this context, which is found by a slight generalization of the approach of [15, 16]. There, the central idea was to construct a mirror of a TCS G_2 manifold by applying mirror symmetry to either both, or to one of the two acyl Calabi-Yau threefolds in the TCS decomposition (2.31). The present example of an extra-twisted connected sum is constructed as a free quotient \hat{M}_k of the TCS G_2 manifolds M_k , which can also be described as acting separately on the acyl Calabi-Yau threefold X_{\pm} .

Omitting the action of σ_2 in (2.1), the TCS decomposition and the action of the TCS mirror map on the resulting G_2 manifolds M_k was described in [16]. The result is that both X_{\pm} are such that

$$\begin{aligned} |K_{\pm}| &= 4 \\ h^{2,1}(Z_{\pm}) &= 4 \\ |N_{\pm}| &= 10, \\ |N_+ \cap N_-| &= k \end{aligned} \quad (2.38)$$

where $K = \ker(H^{1,1}(X_{\pm}) \rightarrow H^{1,1}(S_{\pm}))$, $N_{\pm} = \text{Im}(H^{1,1}(X_{\pm}) \rightarrow H^{1,1}(S_{\pm}))$ and Z_{\pm} is the compactification of X_{\pm} found by gluing in a single K3 fibre (see [13] for more details). The topology of the resulting G_2 manifold is then determined from

$$\begin{aligned} b^2 &= |K_+| + |K_-| + |N_+ \cap N_-| \\ b^2 + b^3 &= 23 + 2(|K_+| + |K_-|) + 2(h^{2,1}(Z_+) + h^{2,1}(Z_-)). \end{aligned} \quad (2.39)$$

Here, ϕ and ψ have the range $0 \dots 2\pi$ and η parametrizes the interval $0 \dots \pi$, so that we can see S^3 decomposed as a being glued from the two solid tori described by the above for $0 \leq \eta \leq \pi/2$ and $\pi/2 \leq \eta \leq \pi$. In this parametrization, (2.33) is simply the map

$$\sigma_2 : \begin{aligned} \phi &\rightarrow \phi + \pi \\ \psi &\rightarrow \psi + \pi \end{aligned} \quad (2.35)$$

which is nothing but the half-shift induced by σ_2 in (2.1) again.

Orbifolding $X_{\pm} \times \mathbb{S}_{\pm}^1$ by the freely acting \mathbb{Z}_2^s produces a 7-manifold V_{\pm} with the holonomy group $SU(3) \times \mathbb{Z}_2^s$ ('barely G_2 '). In the present case, both the elliptic fibrations on X_{\pm} and the SYZ fibrations on X_{\pm} become fibrations by T^2 and T^3 on V_{\pm} . Repeating the analysis of [16] then motivates to consider two types of mirrors of \hat{M}_k :

$$\hat{M}^{\vee} = V_+^{\vee} \# V_-^{\vee} \tag{2.40}$$

associated with applying four T-dualities, as well as

$$\hat{M}^{\wedge} = V_+ \# V_-^{\vee} \tag{2.41}$$

associated with applying three T-dualities.

Given the data of X_{\pm} , the Betti numbers of the free quotients M_k/\mathbb{Z}_2^s are given by

$$\begin{aligned} b^2 &= |K_+^e| + |K_-^e| + |N_+ \cap N_-| \\ b^2 + b^3 &= 23 + 2(|K_+^e| + |K_-^e|) + 2(h_e^{2,1}(Z_+) + h_e^{2,1}(Z_-)) \end{aligned} \tag{2.42}$$

where e and $_e$ indicates taking the even subspace under the involution \mathbb{Z}_2^s . Note that the group \mathbb{Z}_2^s does not act on the K3 fibres of X_{\pm} , so that N_{\pm} are unchanged. In the present case, we have that

$$\begin{aligned} |K_{\pm}^e| &= 2 \\ h_e^{2,1}(Z_{\pm}) &= 2 \\ |N_{\pm}| &= 10 \\ |N_+ \cap N_-| &= k \end{aligned} \tag{2.43}$$

so that we recover

$$\begin{aligned} b^2(\hat{M}) &= 4 + k \\ b^3(\hat{M}) &= 35 - k \end{aligned} \tag{2.44}$$

We are now ready to discuss the action of the TCS mirror maps. The fact that X_{\pm} are self-mirror indicates that the same is true for V_{\pm} . As in the analysis of the Joyce orbifold in [16], the only non-trivial ingredient in the mirror construction is given by $N_+ \cap N_-$. As σ_2 does not act on the K3 fibres at all, we can just quote the result of the analysis of [16]: whereas $|N_+ \cap N_-| = k$ for \hat{M}^{\vee} , $|N_+ \cap N_-| = 8 - k$ for \hat{M}^{\wedge} . We hence find that

$$\begin{aligned} b^2(\hat{M}_k^{\vee}) &= 4 + k \\ b^3(\hat{M}_k^{\vee}) &= 35 - k \end{aligned} \tag{2.45}$$

whereas

$$\begin{aligned} b^2(\hat{M}_k^{\wedge}) &= 12 - k \\ b^3(\hat{M}_k^{\wedge}) &= 27 + k \end{aligned} \tag{2.46}$$

As in [16], we can associate \hat{M}^{\vee} with the image under \mathcal{I}_4^+ and \hat{M}^{\wedge} with the image under \mathcal{I}_3^- .

3 Spin(7) examples

In this section we will analyse two T^8/\mathbb{Z}_2^4 orbifolds in which the generators α, β, γ and δ of $\Gamma = \mathbb{Z}_2^4$ act on the T^8 coordinates as [5]:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7	X^8
α	-	-	-	-	+	+	+	+
β	+	+	+	+	-	-	-	-
γ	$-\frac{c_1}{2}$	$-\frac{c_2}{2}$	+	+	$-\frac{c_5}{2}$	$-\frac{c_6}{2}$	+	+
δ	$-\frac{d_1}{2}$	+	$-\frac{d_3}{2}$	+	$-\frac{d_5}{2}$	+	$-\frac{d_7}{2}$	+

(3.1)

where $c_i, d_i \in \{0, 1\}$. The two cases we will study have $c_i = (1, 1, 1, 1)$ and $d_i = (0, 1, 1, 1)$ ('Example 1'), and $c_i = (1, 0, 1, 0)$ and $d_i = (0, 1, 1, 1)$ ('Example 2'). Once again, $-1/2$ is shorthand for $X^i \rightarrow 1/2 - X^i$

3.1 Smoothing of the orbifold

First, let us review the resolution of the T^8/\mathbb{Z}_2^4 orbifolds above into a family of compact Spin(7) manifolds as described in [5]. In total, the Betti numbers will receive a contribution from the orbifold itself, as well as contributions from the resolutions of the singularities at the fixed points. Before resolution, the cohomology consists of the classes of T^8 which are invariant under the group \mathbb{Z}_2^4 . In either example, a simple calculation shows $b^0 = b^8 = 1$ and $b^4 = 14$. For example, $dX^1 \wedge dX^2 \wedge dX^3$ is not invariant, while $dX^1 \wedge dX^2 \wedge dX^3 \wedge dX^4$ is. In particular, the 14 4-forms are precisely the 14 elementary ones appearing in the 4-form on \mathbb{R}^8 that Spin(7) is defined to leave invariant (see e.g. equation (1) of [5]).

Now consider example 1, we would like to understand its resolutions and so must first understand its singular set. In this case, it is clear that only $\alpha, \beta, \gamma, \delta$ and $\alpha\beta$ have fixed points, as all other combinations involve $X^i \rightarrow X^i + 1/2$ for some i , which has no fixed points. So, the singular set consists of 5 sectors, $S_\alpha, S_\beta, S_\gamma, S_\delta$ and $S_{\alpha\beta}$.

For each single generator, the fixed points correspond to 16 T^4 's for the directions unchanged by the group action (For example, α has fixed points $X^i \in \{0, 1/2\}$ for $i = 1 \dots 4$ and X^i free for $i = 5 \dots 8$). For $\alpha\beta$, the only composite element with fixed points, they are the 256 points $X^i \in \{0, 1/2\}$ for $i = 1 \dots 8$. However, we must also consider the action of the other group generators on a given set of fixed points. For the case of α, β acts on the fixed T^4 's by -1 , giving 16 $T^4/\{\pm 1\}$'s. The subgroup $\langle \gamma, \delta \rangle$ acts freely, and so groups S_α into 4 orbits of $T^4/\{\pm 1\}$'s. In a similar way, we see that the set S_β also consists of 4 copies of $T^4/\{\pm 1\}$. For the γ -fixed points, $\langle \alpha, \beta, \delta \rangle$ acts freely and so we find that the singular set gets reduced to 2 T^4 's, and we apply the same reasoning to find that S_δ is also 2 sets of T^4 's. Lastly, $\langle \gamma, \delta \rangle$ acts freely on the $\alpha\beta$ fixed points, grouping the 256 points into 64

(here, $\langle \alpha\beta \rangle$ acts freely). In summary, the fixed set for example 1 looks like:

	singular set	elements in orbit under Γ	singular set in quotient
α	16 T^4 's	4 $T^4/\{\pm 1\}$'s	4 $T^4/\{\pm 1\}$'s
β	16 T^4 's	4 $T^4/\{\pm 1\}$'s	4 $T^4/\{\pm 1\}$'s
γ	16 T^4 's	8 T^4 's	2 T^4 's
δ	16 T^4 's	8 T^4 's	2 T^4 's
$\alpha\beta$	256 points	4 points	64 points

Lastly, we consider the neighbourhood of each singular point. For S_α and S_β , the neighbourhoods of the fixed loci are a total of 8 copies of $(T^4/\{\pm 1\}) \times (B_\zeta^4/\{\pm 1\})$ for a 4-ball B_ζ^4 . In the language of Proposition 3.1.1 in [5], these come from type (ii) singularities, which upon resolution increases b^1 by 1 and b^4 by 6. For S_γ and S_δ , the neighbourhoods consist of 4 $T^4 \times (B_\zeta^4/\{\pm 1\})$'s, which arise from type (i) singularities, and increase b^2 by 1, b^3 by 4 and b^4 by 6. Lastly, the neighbourhoods of $S_{\alpha\beta}$ consist of 64 $(B_\zeta^4/\{\pm 1\}) \times (B_\zeta^4/\{\pm 1\})$'s, which are of type (iii) and increase b^4 by 1. In total, we find that upon resolution the non-trivial Betti numbers of the manifold are:

$$(b^2, b^3, b^4) = (12, 16, 150) . \quad (3.2)$$

Let us now consider example 2. In this case, all sectors except the γ -fixed points contribute the same singularities, and thus Betti numbers, as example 1. However, for γ , $\alpha\delta$ now acts trivially. Starting from a neighbourhood $T^4 \times B_\zeta^4$, the γ action converts this to $T^4 \times (B_\zeta^4/\{\pm 1\})$, but in fact the action of $\alpha\delta$ turns this into the neighbourhood of a singularity of type (iv) [5]. The action of the rest of the group then orders these 16 type (iv) singularities into 4 orbits. In summary, the singular set of example 2 is given by

	singular set	elements in orbit under Γ	singular set in quotient
α	16 T^4 's	4 $T^4/\{\pm 1\}$'s	4 $T^4/\{\pm 1\}$'s
β	16 T^4 's	4 $T^4/\{\pm 1\}$'s	4 $T^4/\{\pm 1\}$'s
γ	16 T^4 's	4 T^4 's	4 T^4 's
δ	16 T^4 's	8 T^4 's	2 T^4 's
$\alpha\beta$	256 points	4 points	64 points

In this case, the singularities induced by the action of γ may be resolved in two inequivalent ways, one increasing b^2 by 1, b^3 by 2 and b^4 by 2, and the other increasing b^3 by 2 and b^4 by 4. As a result, including the contribution from the unresolved orbifold and letting $j \in \{0 \dots 4\}$ be the number of type (iv)'s we resolve in the first way, we find a family of 5 Spin(7) manifolds with Betti numbers given by:

$$(b^2, b^3, b^4) = (10 + j, 16, 154 - 2j) , \quad j \in \{0 \dots 4\} .$$

3.2 Discrete torsion and cohomology

In this section we will recover the resolution of the orbifold by studying the associated CFT. Once again, following a similar analysis used in [9] we will relate the distribution of the allowed discrete torsion signs (which will mostly be derived in the appendix) to the cohomology of the orbifold resolutions.

3.2.1 Example 1

We begin with Example 1, where $c_i = (1, 1, 1, 1)$ and $d_i = (0, 1, 1, 1)$. The cohomology of the smooth limit of the orbifold must be related to the ground states of RR ground states in each twisted sector. In particular, this means we are studying zero momentum states, which only arise in twisted sectors associated with group elements acting non-freely. Thus we only need to consider contributions from the generators $\alpha, \beta, \gamma, \delta$, as well as $\alpha\beta$. Using the analysis in appendix C, we find the following results for the discrete torsion phases:

α and β sectors: in each case the 16 twisted sectors get ordered into 4 sets of 4 labelled by numbers f_α and f_β , both running from 1 to 4 and thus contributing 4 ground states each. In each sector, we have a discrete torsion sign $\epsilon_{f_\alpha}(\beta)$ for the action of β on the α states and a sign $\epsilon_{f_\beta}(\alpha)$ in the other direction. In this case modular invariance requires the distribution of these signs within the α and β sectors to be the same:

$$\sum_{f_\alpha=1}^4 \epsilon_{f_\alpha}(\beta) = \sum_{f_\beta=1}^4 \epsilon_{f_\beta}(\alpha) . \tag{3.3}$$

γ and δ sectors: in these sectors, the 16 fixed points get organized into 2 sets of 8, however there is no discrete torsion sign available and so we only have 2 ground states to work with.

$\alpha\beta$ sector: in this sector, we have 256 fixed points organized into 64 sets of 4 labelled by a number $f_{\alpha\beta}$ running from 1 to 64. We have a discrete torsion degree of freedom for both the action of α and β on these states, which in fact must be equal and we call it $\epsilon_{f_{\alpha\beta}}(\alpha, \beta)$. We can then label the fixed points such that $f_{\alpha\beta} = f_\alpha + 4k$ for $k = 0, \dots, 15$ and get $\epsilon_{f_\alpha+4k}(\alpha, \beta) = \epsilon_{f_\alpha}$. In other words, the 64 possible signs get grouped into 4 sets of 16 signs which must all be equal, and how we distribute such signs across these sets of 16 must correspond to how we distribute the signs in the α and β sectors:

$$\sum_{f_{\alpha\beta}=1}^{64} \epsilon_{f_{\alpha\beta}}(\beta) = 16 \sum_{f_\alpha=1}^4 \epsilon_{f_\alpha} \tag{3.4}$$

With this in mind, we can now determine the contribution to the cohomology from the different twisted sectors. We begin first with the γ and δ sectors, for which there are no discrete torsion phases. In either case, the 16 fixed point/twist fields get grouped into 2 sets of 8 states. Following the methods in [9], we need to find linear combinations of these 8 states in each of the two sectors (we call such a combination $|0, 0; f_g\rangle_g$, for $f_g = 1, 2$ and $g = \gamma$ or δ) such that when acted on by the raising operators ψ_+^i (where $i \in I_\gamma = \{3, 4, 7, 8\}$ for γ and $i \in I_\delta = \{2, 4, 6, 8\}$ for δ) we obtain \mathbb{Z}_2^4 invariant states. In doing so we find states of the following form:

$$\begin{aligned} &|0, 0; f_g\rangle_g , & \psi_+^i |0, 0; f_g\rangle_g , & \psi_+^i \psi_+^j |0, 0; f_g\rangle_g , \\ &\psi_+^i \psi_+^j \psi_+^k |0, 0; f_g\rangle_g , & \psi_+^i \psi_+^j \psi_+^k \psi_+^l |0, 0; f_g\rangle_g . \end{aligned}$$

Where $i, j, k, l \in I_g$ and $g = \gamma$ or δ . In total, there is one state of the first form, 4 of the second, 6 of the 3rd, 4 of the 4th and one of the 5th for each of $g = \gamma, \delta$. Thus, identifying $|0, 0; f\rangle_g \simeq$ a 2 form as in [9] and $\psi_+^i \simeq dX^i$ we find the contribution (recalling that the distribution of signs within each sector must be the same):

$$\begin{aligned} \delta b_g^2 &= 2 \\ \delta b_g^3 &= 8 \\ \delta b_g^4 &= 12 \end{aligned} \tag{3.5}$$

for $g = \gamma$ or δ .

The α , β and $\alpha\beta$ sectors are more interesting, as they involve the choice of a discrete torsion sign. Beginning with α (the β sector is identical), we recall that our 16 twist fields were grouped into 4 sets of 4 and follow a similar procedure to above. Provided we choose the appropriate signs between states in one orbit, we find that when we choose $\epsilon_{f_\alpha}(\beta)$ or $\epsilon_{f_\beta}(\alpha)$ to be +1 we get states:

$$|0, 0; f_g\rangle_g, \quad \psi_+^i \psi_+^j |0, 0; f_g\rangle_g, \quad \psi_+^i \psi_+^j \psi_+^k \psi_+^l |0, 0; f_g\rangle_g,$$

Where $i, j, k, l \in I_\alpha = \{5, 6, 7, 8\}$ or $I_\beta = \{1, 2, 3, 4\}$ and $g = \alpha$ or β . So, we have 1 state with no oscillators, 6 with 2, and 1 with 4. When we take it to be -1 , we get states with 1 and 3 ψ_+^i 's with i taking values in the same possible index set. So, we would find 1 state with no oscillators, 4 with 1, and 4 with 3. If k is the number of signs we take to be +1, and we make the same identifications of states with forms as before, we find the contribution:

$$\begin{aligned} \delta b_g^2 &= k \\ \delta b_g^3 &= 16 - 4k \\ \delta b_g^4 &= 6k \end{aligned} \tag{3.6}$$

for $g = \alpha$ or β and $k = 0, \dots, 4$.

For $\alpha\beta$, the 256 fixed points were organized into 64 sets of 4. There are no oscillators to use as raising operators, so we only have states $|0, 0; f_{\alpha\beta}\rangle_{\alpha\beta}$. However, there are 2 interesting things to note here. Firstly, this state should not be interpreted as a 2 form, but rather a 4 form (loosely, we can think of these as products of α and β singularities, and so the associated form as a product of 2 forms, giving a 4 form). Secondly, recall that choosing one of the α or β discrete torsion phases to be -1 meant choosing 16 of the $\epsilon_{f_{\alpha\beta}}(\alpha, \beta)$'s to be -1 for consistency. When we do so, because α and β act diagonally on the states within the 64 orbits, we find we cannot create any invariant linear combination. Thus, if k is again the number of positive signs we find that this sector contributes:

$$\delta b_{\alpha\beta}^4 = 16k \tag{3.7}$$

We must also include the contribution from the untwisted sectors, which arises from the states $\prod_i (\psi_+^i)^{l_i} |0\rangle$ where $|0\rangle$ is the untwisted ground state and each $l_i \in \{0, 1\}$. In particular, we identify $|0\rangle$ with the constant 0-form, each ψ_+^i with dX^i , and take \mathbb{Z}_2^4 invariant states. As expected, we find that they are in one to one correspondence with the invariant

classes in $H^i(T^8/\mathbb{Z}_2^4)$, and so the total Betti numbers are, after adding up each δ_g^p for $g = \alpha, \beta, \gamma, \delta$ and $\alpha\beta$:

$$(b^2, b^3, b^4)(M_k) = (10, 2k + 4, 48 - 8k), \quad k \in \{0 \dots 4\}. \quad (3.8)$$

Rather interestingly, this does not actually completely agree with the result in [5]! Rather, it only reduces to the expected result when $k = 4$. At first, we may naively believe we have found a set of other new resolutions, but on returning to the previous analysis we find that this is not quite the case. Recall that when $k < 4$, some of the 64 $\alpha\beta$ -twisted sectors could not provide us with states, as they were not invariant. We then interpret this as saying we are *not* resolving 16 of the associated $\alpha\beta$ singularities. Thus, when $k < 4$, we do not actually have a complete resolution — rather, we have only a partial resolution of the orbifold and the remaining singularities are frozen due to the presence of discrete torsion.

3.2.2 Example 2

Now we analyze example 2, where $c_i = (1, 0, 1, 0)$ and $d_i = (0, 1, 1, 1)$. This model is very similar to example 1, the full analysis for this sector is done in appendix C. The only crucial difference being that the $\gamma, \alpha\delta$ and $\alpha\delta\gamma$ sectors all receive a discrete torsion sign — however, modular invariance constrains them so that we may set the distribution of these signs must be the same within each sector:

$$4 \sum_{f_\gamma=1}^4 \epsilon_{f_\gamma}(\delta) = 4 \sum_{f_{\alpha\delta\gamma}=1}^4 \epsilon_{f_{\alpha\delta\gamma}}(\delta) = 2 \sum_{f_{\alpha\delta}=1}^8 \epsilon_{f_{\alpha\delta}}(\gamma). \quad (3.9)$$

Once again, of these 3 only the γ sector will have zero momentum states to contribute to the cohomology, and the contributions δb_g^p for $g = \alpha, \beta, \delta$ and $\alpha\beta$ are the same as in example 1.

The 16 fixed points of the γ sector are organized into 4 sets of 4 labelled by a number $f_\gamma \in \{1, \dots, 4\}$, and so we build states from linear combinations of states within the orbits, which we call $|0, 0; f_\gamma\rangle_\gamma$. We then fill out the ground states with ψ_+^i for $i = 3, 4, 7, 8$. The discrete torsion phases come from the actions of $\delta, \alpha\delta$ and $\alpha\delta\gamma$, but consistency of the representation requires $\epsilon_{f_\gamma}(\delta) = \epsilon_{f_\gamma}(\alpha\delta) = \epsilon_{f_\gamma}(\alpha\delta\gamma) \equiv \epsilon_{f_\gamma}$. When we choose $\epsilon_{f_\gamma} = +1$, we find that the only invariant states are:

$$|0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^3 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^8 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^3 \psi_+^8 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^4 \psi_+^7 |0, 0; f_\gamma\rangle_\gamma,$$

whereas when $\epsilon_{f_\gamma} = -1$ we find the states:

$$\begin{aligned} \psi_+^4 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^7 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^3 \psi_+^4 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^3 \psi_+^7 |0, 0; f_\gamma\rangle_\gamma, \\ \psi_+^4 \psi_+^8 |0, 0; f_\gamma\rangle_\gamma, \quad \psi_+^7 \psi_+^8 |0, 0; f_\gamma\rangle_\gamma. \end{aligned}$$

Letting j be the number of ϵ_{f_γ} which are equal to 1, we find that the contribution from this sector to the cohomology of the associated resolution is:

$$\begin{aligned} \delta b_\gamma^2 &= j \\ \delta b_\gamma^3 &= 8 \\ \delta b_\gamma^4 &= 16 - 2j \end{aligned} \quad (3.10)$$

and thus the total cohomology is given by:

$$(b^2, b^3, b^4)(M_{j,k}) = (2 + 2k + j, 48 - 8k, 42 + 28k - 2j) , \quad (3.11)$$

where $M_{j,k}$ represents the space we obtain after the different resolutions of the orbifold, parametrized by j and k . When $k = 4$ we get:

$$(b^2, b^3, b^4)(M_{j,4}) = (10 + j, 16, 154 - 2j) . \quad (3.12)$$

which are the Betti numbers found in [5]. Once again, a choice of $k < 4$ means we cannot create certain states in the $\alpha\beta$ sector and some of the singularities are frozen by discrete torsion.

3.3 Mirror symmetry

In this section, we will briefly discuss mirror symmetry for the smooth Spin(7) manifolds obtained in example 2. Example 1 can be treated analogously, with the result that all of the (partially or fully resolved) models obtained there are self-mirror.

The starting point is the Spin(7) superconformal algebra, the generators of which can be obtained from the generators $(T_{G_2}, G_{G_2}, \Phi_{G_2}, X_{G_2}, K_{G_2}, M_{G_2})$, (2.24), of the G_2 algebra as follows:

$$\begin{aligned} T &= T_{G_2} + \frac{1}{2} : \partial X^8 \partial X^8 : - \frac{1}{2} : \psi^8 \partial \psi^8 : , \\ G &= G_{G_2} + : \psi^8 \partial X^8 : , \\ X &= X_{G_2} + \Phi_{G_2} \psi^8 + \frac{1}{2} \psi^8 \partial \psi^8 , \\ M &= [G, X] = \partial X^8 \Phi_{G_2} - K_{G_2} - M_{G_2} + \frac{1}{2} \partial^2 X^8 \psi^8 - \frac{1}{2} \partial X^8 \partial \psi^8 . \end{aligned}$$

The algebra these operators satisfy is worked out in [1]. Written this way, it is easy to work out the analogue of the G_2 mirror automorphism by combining the G_2 map:

$$(T_{G_2}, G_{G_2}, \Phi_{G_2}, X_{G_2}, K_{G_2}, M_{G_2}) \rightarrow (T_{G_2}, G_{G_2}, -\Phi_{G_2}, X_{G_2}, -K_{G_2}, M_{G_2}) \quad (3.13)$$

together with a T-duality along X^8 . Doing so, we see that this combination maps the algebra directly back on to itself. Using combinations of the 3-direction T-dualities found in [9] combined with T-duality along X^8 , we can then explicitly realize this automorphism as such a duality. We find that the following 7 combinations generate our Spin(7) automorphism:

$$\{(2, 4, 6, 8), (2, 3, 5, 8), (1, 2, 7, 8), (1, 3, 6, 8), (1, 4, 5, 8), (3, 4, 7, 8), (5, 6, 7, 8)\} .$$

Another option is to use an automorphism that leaves the G_2 algebra invariant, without affecting the terms involving X^8 or ψ^8 . This can be done using the combinations of 4 T-dualities found in [9], which leave the G_2 algebra invariant and so map the Spin(7) algebra

on to itself once again. The following 7 index sets generate an automorphism of the Spin(7) algebra in this way:

$$\{(1, 2, 5, 7), (1, 4, 6, 7), (3, 4, 5, 6), (2, 4, 5, 7), (2, 3, 6, 7), (1, 2, 5, 6), (1, 2, 3, 4)\} .$$

We would now like to see how these maps act on any discrete torsion signs. To do so, we construct representations of the operators in terms of the ψ_0^i and $\tilde{\psi}_0^i$'s. Our first focus is on the γ sector, for which it is the $\alpha\delta$ parity that we are interested in. A quick calculation gives:

$$\alpha\delta|_{\mathcal{H}^{f_\gamma}} = \frac{1}{4}\psi_0^4\psi_0^7\tilde{\psi}_0^4\tilde{\psi}_0^7\epsilon_{f_\gamma}(\delta) .$$

Next, we want to consider the representation of α in the β sector, β in the α sector, and α and β in the $\alpha\beta$ sector. We find, for β in the α sector:

$$\beta|_{\mathcal{H}^{f_\alpha}} = \frac{1}{16}\psi_0^5\psi_0^6\psi_0^7\psi_0^8\tilde{\psi}_0^5\tilde{\psi}_0^6\tilde{\psi}_0^7\tilde{\psi}_0^7 \cdot \epsilon_{f_\alpha}(\beta) . \tag{3.14}$$

and for α acting in the β sector:

$$\alpha|_{\mathcal{H}^{f_\beta}} = \frac{1}{16}\psi_0^1\psi_0^2\psi_0^3\psi_0^4\tilde{\psi}_0^1\tilde{\psi}_0^2\tilde{\psi}_0^3\tilde{\psi}_0^4 \cdot \epsilon_{f_\beta}(\alpha) . \tag{3.15}$$

In the $\alpha\beta$ sector we have no zero modes, and so there is no representation of α or β in terms of the ψ 's (i.e. it can be represented by $\alpha|_{\mathcal{H}^{f_{\alpha\beta}}} = \beta|_{\mathcal{H}^{f_{\alpha\beta}}} = \epsilon_{f_{\alpha\beta}}(\alpha, \beta)$). Applying the T-dualities, we find that the 14 possible combinations group into two sets:

$$\begin{aligned} I_+ &= \{(2, 3, 5, 8), (1, 3, 6, 8), (3, 4, 7, 8), (1, 4, 6, 7), (2, 4, 5, 7), (1, 2, 5, 6)\} . \\ I_- &= \{(2, 4, 6, 8), (1, 2, 7, 8), (1, 4, 5, 8), (5, 6, 7, 8), (1, 2, 5, 7), (3, 4, 5, 6), \\ &\quad (2, 3, 6, 7), (1, 2, 3, 4)\} . \end{aligned}$$

Those in I_- effectively swap the discrete torsion signs $\epsilon_{f_\gamma}(\delta)$, while those in I_+ leave them alone. Interestingly, none of these 14 combinations change the signs in the α , β or $\alpha\beta$ sectors. Thus we have the set of dualities:

$$\begin{aligned} I_- : M_{j,k} &\rightarrow M_{4-j,k} \\ I_+ : M_{j,k} &\rightarrow M_{j,k} \end{aligned} . \tag{3.16}$$

Both of these maps take type IIA string theory to type IIA and IIB to IIB. When $k = 4$, these are the dualities found in [10]. However, for $k = 0, 1, 2, 3$ these are dualities between singular manifolds, which were not found in their analysis. Note there is no combination of T-dualities which change k .

4 Spin(7) mirror maps for connected sums

In this section we consider mirror maps for Spin(7) manifolds realized as generalized connected sums (GCS) [27] and show that these agree with our results obtained above.

4.1 Constructing Spin(7) manifolds as generalized connected sums

As a preparation, let us briefly review the construction of GCS Spin(7) manifolds of [27]. The building blocks from which such Spin(7) are formed are a asymptotically cylindrical Calabi-Yau fourfold Z_+ with asymptotic neck region $X_3 \times S^1 \times I$ for a Calabi-Yau threefold X_3 and an interval I , and an asymptotically cylindrical G_2 manifold Z_- with neck region $X_3 \times I$. Taking $Z_- \times S^1$ and identifying the isomorphic neck regions $X_3 \times S^1 \times I$, we may then form a compact eight-dimensional manifold Z as the generalized connected sum

$$Z = Z_+ \# [Z_- \times S^1] . \tag{4.1}$$

Based on a number of observations, it has been conjectured in [27] that there exists a Ricci flat metric of holonomy Spin(7) on such manifolds. The evidence for this is as follows. First of all, the examples of Spin(7) manifolds realized as resolutions of T^8/Γ for Γ a finite subgroup of Spin(7) given in [5] allow precisely such a decomposition. We reviewed a decomposition such as (4.1) below in section 4.3. Second, compactifications of heterotic string theory on TCS G_2 manifolds should have a lift to M-Theory on a Spin(7) manifold which can also be decomposed into two pieces. By applying an appropriate fibrewise duality map to a TCS G_2 manifold, the authors of [27] argued that one finds a decomposition such as (4.1) on the M-Theory side and checked the equivalence of the spectra of light fields in a few examples. Finally, acyl G_2 manifolds Z_- can be realized as (a resolution of) a quotient $(X_3 \times \mathbb{R})/\mathbb{Z}_2$ in which the \mathbb{Z}_2 acts as an anti-holomorphic quotient on X_3 and as $t \rightarrow -t$ on \mathbb{R} . In this case, Z can be globally described as a resolution of an anti-holomorphic quotient of a suitably chosen Calabi-Yau fourfold Y , recovering the construction of [35].

For an acyl Calabi-Yau fourfold Z_+ and an acyl G_2 manifold Z_- given as a resolution of $(X_3 \times \mathbb{R})/\mathbb{Z}_2$, the Betti numbers of Z are found to be [27]

$$\begin{aligned} b^1(Z) &= 0 \\ b^2(Z) &= n_+^2 + n_-^2 + b_e^2 \\ b^3(Z) &= n_-^2 + n_-^3 + n_+^3 \\ b^4(Z) &= n_-^3 + n_-^4 + n_+^4 + b_o^2 + b_o^3 + b_e^3 + b_e^4 \end{aligned} \tag{4.2}$$

Here n_{\pm}^i are the kernels of the restriction maps

$$\begin{aligned} \beta_+^i &: H^i(Z_+, \mathbb{Z}) \rightarrow H^i(X_3 \times \mathbb{S}^1, \mathbb{Z}) \\ \beta_-^i &: H^i(Z_-, \mathbb{Z}) \rightarrow H^i(X_3, \mathbb{Z}) . \end{aligned} \tag{4.3}$$

and b_o^i and b_e^i are the dimensions of the odd/even subspaces of the i -th cohomology group of X_3 under the action \mathbb{Z}_2 . Furthermore, we have assumed that the images of β_+^2 and β_+^4 are surjective and that $H^3(Z_+) = \ker \beta_+^3$ holds.³

By using the fact that there exists a single covariantly constant spinor on Z , it follows that

$$b_-^4(Z) + 1 = -8 + \frac{1}{3} (2 - b^2(Z) + b^3(Z) + b^4(Z)) , \tag{4.4}$$

³This last assumptions is slightly weaker than the assumptions made for technical simplicity in [27]. By following the same analysis presented there, it is straightforward to see that (4.2) holds in the present case.

and we can compute

$$b^2(Z) + b^4(Z) + 1 = \frac{2}{3} - 8 + \frac{1}{3} (n_+^4 + 2n_+^2 + n_+^3) + \frac{1}{3} (3n_-^2 + 3n_-^3 + 2(b_o^2 + b_e^2) + 2b_o^3) . \quad (4.5)$$

Here, we have used that for anti-holomorphic involutions $b_o^3 = b_e^3$ and $b_e^4 = b_o^2$ holds.

4.2 A mirror map for GCS Spin(7) manifolds

As shown in [1], exactly marginal deformations of Spin(7) sigma models are counted by (4.5), so a mirror map for a Spin(7) manifold Z must produce another manifold Z^\vee such that

$$b^2(Z) + b^4(Z) + 1 = b^2(Z^\vee) + b^4(Z^\vee) + 1 . \quad (4.6)$$

Furthermore, such a map can be the result of an application of four T-dualities along a calibrated T^4 fibration [6].⁴ The GCS decomposition (4.1) suggests how such a structure might be realized. The acyl Calabi-Yau fourfold Z_+ has a SYZ fibration by T^4 which becomes the T^3 SYZ fibre of X_3 times a circle S^1 in the neck region. On $Z_- \times S^1$, the circle simply becomes the product S^1 while the T^3 SYZ fibre of X_3 sits inside $Z_- = (X_3 \times \mathbb{R})/\mathbb{Z}_2$. We hence expect to find a Spin(7) mirror by performing four T-dualities along this T^4 .

This motivates the following construction: for a Spin(7) manifold realized as a GCS as in (4.1), a mirror is given by⁵

$$Z^\vee = Z_+^\vee \# [Z_-^\vee \times S^1] . \quad (4.7)$$

which are glued along a neck region which is isomorphic to $X_3^\vee \times S^1 \times I$. In particular, Z_-^\vee is constructed from an antiholomorphic involution of X_3^\vee as $Z_-^\vee = (X_3^\vee \times \mathbb{R})/\mathbb{Z}_2$.

In the following, we will collect some evidence for this proposal by showing that (4.6) indeed holds for this construction. In order to prove this, we will stick to the same simplifying assumptions under which (4.2) holds. Our main task is to work out how the topology of Z_\pm^\vee is related to that of Z_\pm . This can be done as follows. There is a compact Calabi-Yau fourfold Y realized

$$Y = Z_+ \# Z_+ \quad (4.8)$$

realized by gluing two copies of Z_+ along $X_3 \times S^1 \times I$, and a G_2 manifold M realized as

$$M = Z_- \# Z_- = (X_3 \times S^1)/\mathbb{Z}_2 \quad (4.9)$$

by gluing two copies of Z_- along $X_3 \times I$. For both of these compact geometries, there are mirror maps which act in the usual way, i.e.

$$\begin{aligned} h^{1,1}(Y) &= h^{3,1}(Y^\vee) \\ h^{2,1}(Y) &= h^{2,1}(Y^\vee) \\ h^{3,1}(Y) &= h^{1,1}(Y^\vee) \end{aligned} \quad (4.10)$$

⁴The moduli space of a Cayley (calibrated) four-cycle N inside a Spin(7) manifold has dimension $-\frac{1}{2}N \cdot N$, so that we can at best hope to approximate such a fibration in a collapsed limit.

⁵We would like to thank Michele del Zotto for suggesting this construction.

and

$$b^2(M) + b^3(M) = b^2(M^\vee) + b^3(M^\vee). \tag{4.11}$$

Furthermore, Y^\vee and M^\vee now have the decompositions

$$Y^\vee = Z_+^\vee \# Z_+^\vee \tag{4.12}$$

glued along $X_3^\vee \times S^1 \times I$, and

$$M^\vee = Z_-^\vee \# Z_-^\vee = (X_3^\vee \times S^1) / \mathbb{Z}_2. \tag{4.13}$$

glued along $X_3^\vee \times I$. Using these relations is the key to find the topology of Z_\pm^\vee in terms of Z_\pm .

Let us now work out the resulting relations in detail. Starting with Y , the Mayer-Vietoris sequence for the decomposition (4.8) gives

$$\begin{aligned} b^2(Y) &= 2n_+^2 + h^{1,1}(X_3) + 1 \\ b^3(Y) &= 2n_+^3 \\ b^4(Y) &= 2n_+^4 + 2h^{1,1}(X_3) + 4h^{2,1}(X_3) + 4 \end{aligned} \tag{4.14}$$

As Y is a Calabi-Yau fourfold and $h^{2,1}(Y) = 0$ there is the relation

$$h^{3,1}(Y) = \frac{1}{6}b^4(Y) - \frac{2}{3}b^2(Y) - \frac{23}{3} \tag{4.15}$$

so that

$$h^{1,1}(Y^\vee) + h^{3,1}(Y^\vee) = \frac{1}{3}(n_+^4 + 2n_+^2) + \frac{2}{3}(h^{1,1}(X_3) + h^{2,1}(X_3)) - \frac{23}{3} \tag{4.16}$$

The mirror map acting on Y must leave the above expression invariant. As this mirror map also maps X_3 to X_3^\vee , so that $h^{1,1}(X_3) + h^{2,1}(X_3)$ is preserved, it follows that $n_+^4 + 2n_+^2$ must also be invariant under the mirror map acting on Z_+ . Furthermore, $h^{2,1}(Y) = h^{2,1}(Y^\vee)$ implies that n_+^3 is the same for Z_+ and Z_+^\vee .

Let us now discuss M . Here, the Mayer-Vietoris sequence yields

$$\begin{aligned} b^2(M) &= b_e^2 + 2n_-^2 \\ b^3(M) &= b_o^2 + b_e^3 + 2n_-^3 \end{aligned} \tag{4.17}$$

Under the mirror map acting on the G_2 manifold M , $b^2(M) + b^3(M) = b^2(M^\vee) + b^3(M^\vee)$. As $b_e^2 + b_o^2 + b_e^3 = h^{1,1}(X_3) + h^{2,1}(X_3) + 1$ for anti-holomorphic involutions, this expression is preserved by the mirror map. It hence follows that $n_-^2 + n_-^3$ must also be left invariant under the mirror map acting on Z_- .

Altogether, we have shown that the expressions $n_+^4 + 2n_+^2$, $n_-^2 + n_-^3$ and $b_e^2 + b_o^2 + b_e^3$ are all left invariant under an application of the mirror map acting on Z_+ , Z_- and X_3 . It then follows that the expression (4.5) for GCS Spin(7) manifolds is preserved under the mirror map, i.e. (4.6) holds.

Note that we have not provided an explicit construction of mirrors for Z_+ and Z_- , but only used the topological constraints they have to satisfy to arrive at this conclusion.

It should be possible to give a construction of acyl Calabi-Yau fourfolds from projecting five-dimensional tops as has been done for acyl Calabi-Yau threefolds in [19]. This would in turn allow to derive combinatorial formulae for the topological invariants of Z_+ and Z_+^\vee , which in turn must imply that $n_+^4 + 2n_+^2$ does not change under the mirror map.

It is of course straightforward to describe mirrors of X_3 , but the definition of Z_- furthermore involves specifying an antiholomorphic action of \mathbb{Z}_2 on X_3 and a resolution of the orbifold singularities of $(X_3 \times \mathbb{R})/\mathbb{Z}_2$. Clearly, $b_e^2 + b_o^2 + b_e^3 = h^{1,1}(X_3) + h^{2,1}(X_3) + 1$ does not depend on the details of the antiholomorphic involution chosen. Furthermore, in case there exists a resolution of $(X_3 \times S^1)/\mathbb{Z}_2$ we have [36]

$$b^i(M) = b^i((X_3 \times S^1)/\mathbb{Z}_2) + b^{i-2}(L, \zeta) \tag{4.18}$$

where L is the (real three-dimensional or empty) fixed locus of the involution and ζ is a possible twist. This potentially constrains which antiholomorphic involutions and which resolution can be chosen to construct M^\vee and hence Z_-^\vee .

As shown in [27], the GCS construction of Spin(7) manifolds is closely related to the work of [35], in which Spin(7) manifolds are found by resolving anti-holomorphic quotients of Calabi-Yau fourfolds. This offers another possible perspective on mirror maps of Spin(7) manifolds in general, and the ones considered here in particular.

4.3 Examples

In this section we revisit the two examples of Spin(7) manifolds studies in section 3 and show that the mirror map found there agrees with the GCS mirror map described above in section 4.2.

4.3.1 Example 1

Let us first study the example of section 3.2.1, which has $c = (1, 1, 1, 1)$ and $d = (0, 1, 1, 1)$, and start by describing its GCS decomposition. Such a decomposition can be found by cutting the orbifold along $X_7 = \frac{1}{8}$. At $X_7 = \frac{1}{8}$, only two generators α and γ act non-trivially on $X_1 \dots X_6$, so that we can identify the neck region as $\tilde{X}_3 \times S^1 \times I$, where X_7 is a coordinate on I and X_8 a coordinate on the S^1 . Resolving the orbifold $\tilde{X}_3 = T^6/\langle\alpha, \gamma\rangle$ produces the Calabi-Yau threefold X_3 with

$$\begin{aligned} h^{1,1}(X_3) &= 19 \\ h^{2,1}(X_3) &= 19 \end{aligned} \tag{4.19}$$

Restricting $X_7 \leq \frac{1}{8}$, we find an acyl Calabi-Yau fourfold $\tilde{Z}_+ = T^7 \times \mathbb{R}/\langle\alpha, \beta, \gamma\rangle$. Two copies of \tilde{Z}_+ can be glued to form an Calabi-Yau orbifold \tilde{Y} , which has already been studied in [37]. In their terminology, this case is the fourfold ‘model B’, in which σ is the Nikulin involution with invariants $(r, a, \delta) = (10, 8, 0)$. It can be described as $(K3 \times K3)/\mathbb{Z}_2$ with the \mathbb{Z}_2 acting as the Nikulin involution with invariants $(r, a, \delta) = (10, 8, 0)$ simultaneously on both K3 surfaces. The topology of the resolution Y of \tilde{Y} is given by

$$\begin{aligned} h^{1,1}(Y) &= 24 \\ h^{2,1}(Y) &= 8 \\ h^{3,1}(Y) &= 24 \end{aligned} \tag{4.20}$$

and $\chi(Y) = 288$. The relevant data of the decomposition $Y = Z_+ \# Z_+$ is

$$\begin{aligned} n_+^2 &= 2 \\ n_+^3 &= 8 \\ n_+^4 &= 76 \end{aligned} \tag{4.21}$$

Restricting $X_7 \geq \frac{1}{8}$, we find the product of an S^1 and an acyl G_2 manifold $\tilde{Z}_- = T^6 \times \mathbb{R} / \langle \alpha, \gamma, \delta \rangle$, with X_8 being a coordinate on the product S^1 . Two copies of this orbifold can be glued to the compact G_2 orbifold \tilde{M} which has a unique resolution to a G_2 manifold M with (see [4])

$$\begin{aligned} b^2(M) &= 12 \\ b^3(M) &= 43 \end{aligned} \tag{4.22}$$

Furthermore, the action of δ at $X_7 = \frac{1}{4}$ on X_3 is such that $b_c^2 = 8$ and $b_o^2 = 11$. This determines that

$$\begin{aligned} n_-^2 &= 2 \\ n_-^3 &= 6 \\ n_-^4 &= 6 \end{aligned} \tag{4.23}$$

As a check, one can now use (4.2) to recover the Betti numbers of Z given in (3.8) (note that the complete resolution corresponds to setting $k = 4$ in (3.8)). As we have seen from the CFT analysis this Spin(7) manifold should be considered self-mirror. The same conclusion is reached by applying the GCS mirror map: both Y and M are self-mirror, so that

$$Z_{\pm} = Z_{\pm}^{\vee} \tag{4.24}$$

and our Spin(7) mirror map gives $Z^{\vee} = Z$.

4.3.2 Example 2

Let us now study the example of section 3.2.2, which has $c = (1, 0, 1, 0)$ and $d = (0, 1, 1, 1)$. We can proceed in the same way as for the first example and cut along $X_7 = \frac{1}{8}$. The neck region is again formed as $\tilde{X}_3 = T^6 / \langle \alpha, \gamma \rangle$ with

$$\begin{aligned} h^{1,1}(X_3) &= 19 \\ h^{2,1}(X_3) &= 19 \end{aligned} \tag{4.25}$$

The acyl Calabi-Yau fourfolds \tilde{Z}_+ and its resolution Z_+ found by setting $x_7 \leq \frac{1}{8}$ are the same as in the first example, so that we already know their topological data, (4.20) and (4.21).

The acyl G_2 orbifolds \tilde{Z}_- and \tilde{M} are different in this example, but \tilde{M} is again one of the elementary examples of [4]. Its resolution is not unique and produces nine distinct G_2 manifolds M_l with Betti numbers

$$\begin{aligned} b^2(M_n) &= 8 + l \\ b^3(M_n) &= 47 - l \end{aligned} \tag{4.26}$$

for $l = 0 \dots 8$. The action of δ at $X_7 = \frac{1}{4}$ on X_3 is again such that $b_e^2 = 8$ and $b_o^2 = 11$. We now find that l must be even and that

$$\begin{aligned} n_-^2 &= l/2 \\ n_-^3 &= 8 - l/2. \\ n_-^4 &= 8 - l/2 \end{aligned} \tag{4.27}$$

This data again reproduces (3.8) from (4.2) setting $l = 2j$ and $k = 4$ (again, only the case $k = 4$ corresponds to a complete resolution).

We are now ready to discuss the GCS mirror map for Z . We have $Z_+ = Z_+^\vee$ as before and $M_l^\vee = M_{8-l}$. This means that the GCS mirror map replaces $l \rightarrow 8 - l$ in (4.27), so that it reproduces the CFT results $Z_j^\vee = Z_{4-j}$.

Acknowledgments

We would like to thank Michele del Zotto, Marc-Antoine Fiset, Sakura Schäfer-Nameki and Ashoke Sen for discussions and inspiration related to this project. The work of APB is supported by the ERC Consolidator Grant 682608 ‘‘Higgs bundles: Supersymmetric Gauge Theories and Geometry’’ (HIGGSBNDL). The research of AO was supported by the Perimeter Institute for Theoretical Physics. Research at Perimeter Institute is supported by the Government of Canada through the Department of Innovation, Science and Economic Development Canada and by the Province of Ontario through the Ministry of Economic Development, Job Creation and Trade. AO also acknowledges the support of NSERC through the PGS D scholarship.

A Discrete torsion and modular invariance

To set the stage and outline our strategy, let us review a few basic facts about (generalized) discrete torsion for strings on orbifolds following [9]. Crucially, the definition of string theory on orbifolds in general involves an assignment of discrete torsion phases [29]. String theory on a orbifold of T^n by a group Γ is built from the untwisted sector \mathcal{H}_e composed of Γ -invariant states, as well as a twisted sector \mathcal{H}_g for every non-trivial group element g of Γ . To find the states in the twisted sectors \mathcal{H}_g , we need to study the action of other group elements $h \neq g$ on \mathcal{H}_g . This action in general involves the assignment of phases,

$$h|_{\mathcal{H}_g} = \epsilon_g(h)h^0|_{\mathcal{H}_g}, \tag{A.1}$$

where h^0 refers to the usual action of h in the g -twisted sector as expected from the orbifold group action on the coordinates. These discrete torsion phases must form a representation of Γ and furthermore must satisfy [29]

$$\epsilon_g(h) = \epsilon_{h^c g^d}(h^a g^b) \quad \text{for } ad - bc = 1 \tag{A.2}$$

to guarantee modular invariance. As the twisted sector associated with a group element g typically decomposes as

$$\mathcal{H}_g = \oplus_f \mathcal{H}_{g,f}, \tag{A.3}$$

e.g. in case g has several fixed points labelled by f_g , a different assignment of phases $\epsilon_{f_g}(h)$ for each f_g is possible [9]. Of course, these still have to form a representation of Γ .

If we choose to include such ‘generalized’ discrete torsion phases $\epsilon_{f_g}(h)$ in our model, modular invariance must be reconsidered. The partition function for our models can be written as

$$Z(q, \bar{q}) = \frac{1}{|\Gamma|} \sum_{h, g \in \Gamma} \text{Tr}_{\mathcal{H}_h} \left(g q^{L_0 - c/24} \bar{q}^{L_0 - c/24} \right) \equiv \frac{1}{|\Gamma|} \sum_{h, g \in \Gamma} Z_{h;g}. \tag{A.4}$$

where $Z_{h;g}$ refers to the partition function component restricted to the h twisted sector, as in the summation in the middle. Modular invariance then implies

$$Z(\tau + 1) = Z(\tau) \rightarrow Z_{g;e}(\tau + 1) = Z_{g;g} \tag{A.5}$$

from which $\epsilon_{f_g}(g) = 1$ follows, and

$$Z(-1/\tau) = Z(\tau) \rightarrow Z_{g;h}(-1/\tau) = Z_{h;g}(\tau) \tag{A.6}$$

which constrains possible assignments of the $\epsilon_{f_g}(h)$ by linking them to the phases $\epsilon_{f_h}(g)$.

Although modular invariance for bosonic strings at one loop is sufficient to guarantee modular invariance at higher genus if $\epsilon_{f_g}(h) = \epsilon_{f'_g}(h)$ for all f_g, f'_g , this is not the case for more general assignments. However, studying solutions to the above constraints at least provide us with necessary conditions, which will be enough for our purposes. Furthermore, we are only going to study partition functions of bosonic strings. Although it is generally believed that modular invariance of the bosonic string partition function is necessary and sufficient for modular invariance of the full superstring theory, higher genus amplitudes again present a caveat to this analysis, see [9] for a more detailed discussion. For the examples we are presenting, these subtleties are alleviated by the fact that we can match them to known smooth geometries obtained by a smoothing of the orbifolds in question.

For the examples discussed in this paper, the computation of partition functions is significantly simplified by the fact that all of the elements of the orbifold group act diagonally on $T^7 = (S^1)^7$ or $T^8 = (S^1)^8$. For the sake of brevity, we have omitted the details of these computations.

B Discrete torsion analysis for the G_2 orbifold

In this appendix, we derive the necessary conditions on discrete torsion phases for the G_2 model introduced in section 2. We work out the representation matrices for the orbifold elements in the highest weight states of the different twisted sectors. These matrices will have discrete torsion signs showing up, which are then constrained by trace relations coming from the S-transformation.

The orbifold we are interested in is defined by

	X^1	X^2	X^3	X^4	X^5	X^6	X^7
α	+	+	+	-	-	-	-
β	+	-	-	+	+	$-\frac{1}{2}$	-
γ	-	+	-	+	-	+	$-\frac{1}{2}$
σ_2	+	$+\frac{1}{2}$	+	$+\frac{1}{2}$	+	+	+

As discussed in the section 2, we do not need to analyse all the different twisted sectors in this orbifold. Instead, we would focus on the particular sectors twisted under the action of α , β , γ , $\alpha\beta$, σ_2 and $\sigma_2\alpha\beta$. The α , β , and γ sectors are relevant because they are the only ones that contribute to the ground state spectrum of the orbifold string theory, and the others are needed only in order to fix the discrete torsion phases.

α sector. The α -twisted sector can be decomposed into 16 smaller sectors corresponding to the fixed points of the action of α . They can be labelled by the different values of $\{(X^4, X^5, X^6, X^7) : X^i \in \{0, \frac{1}{2}\}\}$. Now each of these twisted sectors localised at the fixed points, have a highest weight state of zero momentum and zero winding. We want to find the representation matrices for the orbifold elements in the basis of these highest weight states. Under the free action of $\langle\beta, \gamma, \sigma_2\rangle$, we get two 8D irreducible representations corresponding to the two X^5 choices.

We can then assign coordinate labels for the basis states $|j\rangle_\alpha^{f_\alpha}$, where index $f_\alpha = \{1, 2\}$ corresponds to the two choices of X^5 and $j=1, 2, \dots, 8$ enumerates the different choices of the other 3 fixed-point coordinates (X^4, X^6, X^7) :

$$\begin{aligned}
 |1\rangle_\alpha^{f_\alpha} &\sim (0, 0, 0); & |2\rangle_\alpha^{f_\alpha} &\sim \left(0, 0, \frac{1}{2}\right); & |3\rangle_\alpha^{f_\alpha} &\sim \left(0, \frac{1}{2}, 0\right); & |4\rangle_\alpha^{f_\alpha} &\sim \left(0, \frac{1}{2}, \frac{1}{2}\right); \\
 |5\rangle_\alpha^{f_\alpha} &\sim \left(\frac{1}{2}, 0, 0\right); & |6\rangle_\alpha^{f_\alpha} &\sim \left(\frac{1}{2}, 0, \frac{1}{2}\right); & |7\rangle_\alpha^{f_\alpha} &\sim \left(\frac{1}{2}, \frac{1}{2}, 0\right); & |8\rangle_\alpha^{f_\alpha} &\sim \left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right);
 \end{aligned}
 \tag{B.1}$$

In this basis, the orbifold generators act as follows:

$$\begin{aligned}
 \alpha|_{\mathcal{H}_\alpha^{f_\alpha}} &= \mathbb{I}_{8 \times 8} & \beta|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} |1\rangle_\alpha^{f_\alpha} \leftrightarrow |3\rangle_\alpha^{f_\alpha} \\ |2\rangle_\alpha^{f_\alpha} \leftrightarrow |4\rangle_\alpha^{f_\alpha} \\ |5\rangle_\alpha^{f_\alpha} \leftrightarrow |7\rangle_\alpha^{f_\alpha} \\ |6\rangle_\alpha^{f_\alpha} \leftrightarrow |8\rangle_\alpha^{f_\alpha} \end{pmatrix} \\
 \gamma|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} |1\rangle_\alpha^{f_\alpha} \leftrightarrow |2\rangle_\alpha^{f_\alpha} \\ |3\rangle_\alpha^{f_\alpha} \leftrightarrow |4\rangle_\alpha^{f_\alpha} \\ |5\rangle_\alpha^{f_\alpha} \leftrightarrow |6\rangle_\alpha^{f_\alpha} \\ |7\rangle_\alpha^{f_\alpha} \leftrightarrow |8\rangle_\alpha^{f_\alpha} \end{pmatrix} & \sigma_2|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} |1\rangle_\alpha^{f_\alpha} \leftrightarrow |5\rangle_\alpha^{f_\alpha} \\ |2\rangle_\alpha^{f_\alpha} \leftrightarrow |6\rangle_\alpha^{f_\alpha} \\ |3\rangle_\alpha^{f_\alpha} \leftrightarrow |7\rangle_\alpha^{f_\alpha} \\ |4\rangle_\alpha^{f_\alpha} \leftrightarrow |8\rangle_\alpha^{f_\alpha} \end{pmatrix},
 \end{aligned}
 \tag{B.2}$$

where $\mathcal{H}_\alpha = \bigoplus_{f_\alpha} \mathcal{H}_\alpha^{f_\alpha}$ and $\mathcal{H}_\alpha^{f_\alpha}$ is the space spanned by the highest weight states $|j\rangle_\alpha^{f_\alpha}$. The representation matrices for the generators after removing spurious phases are

$$\begin{aligned}
 \alpha|_{\mathcal{H}_\alpha^{f_\alpha}} &= \mathbb{I}_{8 \times 8} & \beta|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} 0 & \mathbb{I}_{2 \times 2} & 0 & 0 \\ \mathbb{I}_{2 \times 2} & 0 & 0 & 0 \\ 0 & 0 & 0 & \mathbb{I}_{2 \times 2} \\ 0 & 0 & \mathbb{I}_{2 \times 2} & 0 \end{pmatrix} \\
 \gamma|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} H & 0 & 0 & 0 \\ 0 & H & 0 & 0 \\ 0 & 0 & H & 0 \\ 0 & 0 & 0 & H \end{pmatrix} & \sigma_2|_{\mathcal{H}_\alpha^{f_\alpha}} &= \begin{pmatrix} 0 & \mathbb{I}_{4 \times 4} \\ \mathbb{I}_{4 \times 4} & 0 \end{pmatrix}
 \end{aligned}
 \tag{B.3}$$

where $H = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$.

β sector. The β sector only differs from the α case in the fixed point coordinate labels. Here they are given by two choices each for the set of coordinates (X^2, X^3, X^6, X^7) . α behaves the same way as β did in the α -sector, γ mixes the 2 X^7 choices as before, and σ_2 mixes the choices for X^2 . So, no discrete torsion phase arises in this sector either.

γ sector. Analogous to the α sector, there are 16 highest weight states in the γ sector which can be identified by their fixed point coordinate labels (X^1, X^3, X^5, X^7) . Under the free action of $\langle \alpha, \beta, \sigma_2 \rangle$, we get eight 2D irreducible representations corresponding to the two choices for X^1 , X^3 , and X^5 each.

Let us now assign coordinate labels for the basis states $|j\rangle_{\mathcal{H}_\gamma}^{f_\gamma}$, where index $f_\gamma = \{1, 2, \dots, 8\}$ corresponds to the 8 choices of (X^1, X^3, X^5) and $j=1,2$ enumerates the two choices for X^7 :

$$|1\rangle_{\mathcal{H}_\gamma}^{f_\gamma} \sim \left(X^7 = \frac{1}{4} \right); \quad |2\rangle_{\mathcal{H}_\gamma}^{f_\gamma} \sim \left(X^7 = \frac{3}{4} \right) \quad (\text{B.4})$$

Then the action of the orbifold generators can be obtained from their action on the coordinate labels, as follows:

$$\alpha|_{\mathcal{H}_\gamma}^{f_\gamma} = \left(|1\rangle_{\mathcal{H}_\gamma}^{f_\gamma} \leftrightarrow |2\rangle_{\mathcal{H}_\gamma}^{f_\gamma} \right) = \beta|_{\mathcal{H}_\gamma}^{f_\gamma}; \quad \gamma|_{\mathcal{H}_\gamma}^{f_\gamma} = \text{id.} = \sigma_2|_{\mathcal{H}_\gamma}^{f_\gamma} \quad (\text{B.5})$$

Removing spurious phases by exploiting commutation relations of the representation matrices, we get:

$$\alpha|_{\mathcal{H}_\gamma}^{f_\gamma} = H; \quad \beta|_{\mathcal{H}_\gamma}^{f_\gamma} = \epsilon_{f_\gamma}(\alpha\beta)H; \quad \gamma|_{\mathcal{H}_\gamma}^{f_\gamma} = \mathbb{I}_{2 \times 2}; \quad \sigma_2|_{\mathcal{H}_\gamma}^{f_\gamma} = \begin{pmatrix} \epsilon_1 & 0 \\ 0 & \epsilon_2 \end{pmatrix} \quad (\text{B.6})$$

where $\epsilon_{f_\gamma}(\alpha\beta) = \pm 1$. The relation $\sigma_2\alpha = \alpha\sigma_2$ yields $\epsilon_1 = \epsilon_2 = \epsilon_{f_\gamma}(\sigma_2) = \pm 1$:

$$\sigma_2 = \epsilon_{f_\gamma}(\sigma_2)\mathbb{I}_{2 \times 2} \quad (\text{B.7})$$

So there are two choices of discrete torsion signs available in each irreducible representation.

$\alpha\beta$ sector. The action of $\alpha\beta$ on the coordinates is given by:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7
$\alpha\beta$	+	-	-	-	-	$+\frac{1}{2}$	+

In the $\alpha\beta$ twisted sector, lowest energy states are labelled by the half-integer mode n_6 taking values in $\pm\frac{1}{2}$, and coordinate labels (X^2, X^3, X^4, X^5) . Now the irreducible representations are 8D, and are spanned by the (n_6, X^2, X^4) coordinate labels. Let us assign the basis states $|j\rangle_{\alpha\beta}^{f_{\alpha\beta}}$, $f_{\alpha\beta} = 1, 2, \dots, 4$ labels the different irreducible representations corresponding to the choices for (X^3, X^5) :

$$\begin{aligned} |1\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(\frac{1}{2}, 0, 0 \right); & |2\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(-\frac{1}{2}, 0, 0 \right); & |3\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(\frac{1}{2}, 0, \frac{1}{2} \right); \\ |4\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(-\frac{1}{2}, 0, \frac{1}{2} \right); & |5\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{1}{2}, 0 \right); & |6\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{1}{2}, 0 \right); \\ |7\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right); & |8\rangle_{\alpha\beta}^{f_{\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{1}{2}, \frac{1}{2} \right) \end{aligned} \quad (\text{B.8})$$

In this basis, we can write the representation matrices for the orbifold elements by looking at their action on the basis states just as we did before for the previous sectors:

$$\alpha|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \beta|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \begin{pmatrix} H & 0 & 0 & 0 \\ 0 & H & 0 & 0 \\ 0 & 0 & H & 0 \\ 0 & 0 & 0 & H \end{pmatrix} \tag{B.9}$$

$$\gamma|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \begin{pmatrix} \epsilon_{f_{\alpha\beta}}^1(\gamma)\mathbb{I}_{4\times 4} & 0 \\ 0 & \epsilon_{f_{\alpha\beta}}^2\mathbb{I}_{4\times 4}(\gamma) \end{pmatrix}; \quad \sigma_2|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \begin{pmatrix} 0 & 0 & 0 & \mathbb{I}_{2\times 2} \\ 0 & 0 & \mathbb{I}_{2\times 2} & 0 \\ 0 & \mathbb{I}_{2\times 2} & 0 & 0 \\ \mathbb{I}_{2\times 2} & 0 & 0 & 0 \end{pmatrix}$$

Here, the discrete torsion sign shows up in the γ matrix: $\epsilon_{f_{\alpha\beta}}^i(\gamma) = \pm 1, i=1,2$; as is expected from our γ sector analysis and the S-transform relations.

σ_2 sector. The action of σ_2 on the coordinates is given by:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7
σ_2	+	$+\frac{1}{2}$	+	$+\frac{1}{2}$	+	+	+

In the σ_2 -twisted sector, the lowest energy states are labelled by the two half-integer winding numbers (n_2, n_4) each taking values in $\pm\frac{1}{2}$. We can then have the following assignment of basis states $|j\rangle_{\sigma_2}$, where $j=1,2,3,4$:

$$|1\rangle_{\sigma_2} \sim \left(\frac{1}{2}, \frac{1}{2}\right); \quad |2\rangle_{\sigma_2} \sim \left(\frac{1}{2}, -\frac{1}{2}\right); \quad |3\rangle_{\sigma_2} \sim \left(-\frac{1}{2}, \frac{1}{2}\right); \quad |4\rangle_{\sigma_2} \sim \left(-\frac{1}{2}, -\frac{1}{2}\right) \tag{B.10}$$

Looking at the action of the different generators on the basis states as listed above, we get the following 4D representation matrices:

$$\alpha = \begin{pmatrix} H & 0 \\ 0 & H \end{pmatrix}; \quad \beta = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}; \quad \gamma = \epsilon_{\sigma_2}(\gamma)\mathbb{I}_{4\times 4}; \quad \sigma_2 = \mathbb{I}_{4\times 4} \tag{B.11}$$

Note that the only non-trivial trace involving a discrete torsion sign in this sector is for γ .

$\sigma_2\alpha\beta$ sector. The action of $\sigma_2\alpha\beta$ on the coordinates is given by:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7
$\sigma_2\alpha\beta$	+	$-\frac{1}{2}$	-	$-\frac{1}{2}$	-	$+\frac{1}{2}$	+

Just as the case with the $\alpha\beta$ sector, the lowest energy states are labelled by $(X^2, X^3, X^4, X^5, n_6)$, each of which takes two values. The irreducible representation, as deduced from the action of the orbifold generators, corresponds to the labels (n^6, X^2, X^4) ,

and can be organised in the basis $|j\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}$, where $j=1,2,\dots,8$, and $f_{\sigma_2\alpha\beta} = 1, 2, \dots, 4$:

$$\begin{aligned}
 |1\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{1}{4}, \frac{1}{4}\right); & |2\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{1}{4}, \frac{1}{4}\right); & |3\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{1}{4}, \frac{3}{4}\right); \\
 |4\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{1}{4}, \frac{3}{4}\right); & |5\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{3}{4}, \frac{1}{4}\right); & |6\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{3}{4}, \frac{1}{4}\right); \\
 |7\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(\frac{1}{2}, \frac{3}{4}, \frac{3}{4}\right); & |8\rangle_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}} &\sim \left(-\frac{1}{2}, \frac{3}{4}, \frac{3}{4}\right)
 \end{aligned}
 \tag{B.12}$$

After absorption of spurious phases via commutation relations and normalization of states, we get a discrete torsion sign arising in γ as follows:

$$\begin{aligned}
 \alpha|_{\mathcal{H}_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}} &= \begin{pmatrix} 0 & H & 0 & 0 \\ H & 0 & 0 & 0 \\ 0 & 0 & 0 & H \\ 0 & 0 & H & 0 \end{pmatrix}; & \beta|_{\mathcal{H}_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}} &= \begin{pmatrix} 0 & 0 & H & 0 \\ 0 & 0 & 0 & H \\ H & 0 & 0 & 0 \\ 0 & H & 0 & 0 \end{pmatrix} \\
 \gamma|_{\mathcal{H}_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}} &= \epsilon_{f_{\sigma_2\alpha\beta}}(\gamma)\mathbb{I}_{8\times 8}; & \sigma_2|_{\mathcal{H}_{\sigma_2\alpha\beta}^{f_{\sigma_2\alpha\beta}}} &= \begin{pmatrix} 0 & 0 & 0 & \mathbb{I}_{2\times 2} \\ 0 & 0 & \mathbb{I}_{2\times 2} & 0 \\ 0 & \mathbb{I}_{2\times 2} & 0 & 0 \\ \mathbb{I}_{2\times 2} & 0 & 0 & 0 \end{pmatrix}
 \end{aligned}
 \tag{B.13}$$

C Discrete torsion analysis for the Spin(7) orbifolds

In this appendix, we will explicitly determine the allowed discrete torsion phases and constraints for the Spin(7) orbifold in section 3 — the general structure follows the same logic as in appendix B.

First, let us recall the definition of the orbifold we are interested in — we focus on example 2 from section 3. This is a T^8/\mathbb{Z}_2^4 orbifold where the group generators α, β, γ and δ act as:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7	X^8
α	-	-	-	-	+	+	+	+
β	+	+	+	+	-	-	-	-
γ	$-\frac{1}{2}$	-	+	+	$-\frac{1}{2}$	-	+	+
δ	-	+	$-\frac{1}{2}$	+	$-\frac{1}{2}$	+	$-\frac{1}{2}$	+

Let us now determine the allowed discrete torsion phases.

α, β and $\alpha\beta$ sectors. We begin with α . Here we have 16 twist fields in one-to-one correspondence with the fixed points of the action of α , as outlined in section 3. These states are labelled by the choice of a coordinate set $\{(X^1, X^2, X^3, X^4) : X^i \in \{0, \frac{1}{2}\}\}$. Under the action of the rest of the group, in particular by $\langle \gamma, \delta \rangle$, these get grouped into 4 sets of 4, corresponding to 4 irreducible representations of the orbifold group. Each of these representations come with a highest weight state of zero momentum and winding, and our goal is to find the matrix representations of the group element in the basis of such highest weight states.

Explicitly, in this sector γ and δ permute the X^1 and X^3 fixed points, and so we can label each of the 4 representations by a number $f_\alpha \in \{1, \dots, 4\}$ corresponding to one of the 4 choices of X^2 and X^4 . In other words, the 16 dimensional space of these highest weight states decomposes further into a sum of 4 dimensional spaces as $\mathcal{H}_\alpha = \oplus_{f_\alpha} \mathcal{H}_\alpha^{f_\alpha}$. In each $\mathcal{H}_\alpha^{f_\alpha}$, the basis of states then consists of vectors $|i\rangle_\alpha^{f_\alpha}$, where the label $i \in \{1, \dots, 4\}$ represents one of the 2-tuples in $\{(X^1, X^3) : X^1, X^3 \in \{0, \frac{1}{2}\}\}$. Explicitly, we set:

$$\begin{aligned} |1\rangle_\alpha^{f_\alpha} &= (0, 0), & |2\rangle_\alpha^{f_\alpha} &= (1/2, 0), \\ |3\rangle_\alpha^{f_\alpha} &= (0, 1/2), & |4\rangle_\alpha^{f_\alpha} &= (1/2, 1/2). \end{aligned} \tag{C.1}$$

In this basis, α acts trivially, β acts diagonally and the action of γ and δ is:

$$\gamma|_{\mathcal{H}_\alpha^{f_\alpha}} : \begin{pmatrix} |1\rangle_\alpha^{f_\alpha} \leftrightarrow |2\rangle_\alpha^{f_\alpha} \\ |3\rangle_\alpha^{f_\alpha} \leftrightarrow |4\rangle_\alpha^{f_\alpha} \end{pmatrix}, \quad \text{and} \quad \delta|_{\mathcal{H}_\alpha^{f_\alpha}} : \begin{pmatrix} |1\rangle_\alpha^{f_\alpha} \leftrightarrow |3\rangle_\alpha^{f_\alpha} \\ |2\rangle_\alpha^{f_\alpha} \leftrightarrow |4\rangle_\alpha^{f_\alpha} \end{pmatrix}. \tag{C.2}$$

Introducing discrete torsion phases, by an appropriate choice of normalization of the basis vectors the matrix representations of γ , δ and β take the form:

$$\gamma|_{\mathcal{H}_\alpha^{f_\alpha}} = \begin{pmatrix} H & 0 \\ 0 & H \end{pmatrix}, \quad \delta|_{\mathcal{H}_\alpha^{f_\alpha}} = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & e^{i\theta} \\ 1 & 0 & 0 & 0 \\ 0 & e^{-i\theta} & 0 & 0 \end{pmatrix} \quad \text{and} \quad \beta|_{\mathcal{H}_\alpha^{f_\alpha}} = \begin{pmatrix} \epsilon_1 & 0 & 0 & 0 \\ 0 & \epsilon_2 & 0 & 0 \\ 0 & 0 & \epsilon_3 & 0 \\ 0 & 0 & 0 & \epsilon_4 \end{pmatrix}. \tag{C.3}$$

with $\epsilon_i^2 = 1$. The requirement that all group elements commute sets the phases in $\delta|_{\mathcal{H}_\alpha^{f_\alpha}}$ to 1, and forces all ϵ_i to be equal, $\epsilon_i \equiv \epsilon_{f_\alpha}(\beta)$ for all i , so that $\beta = \epsilon_{f_\alpha}(\beta) \cdot \mathbb{I}_{4 \times 4}$. So, the 16 signs we would expect in $Z_{\alpha;\beta}$ are identified in 4's, reducing to only 4 degrees of freedom $\epsilon_{f_\alpha}(\beta)$.

The β analysis is virtually identical, only differing by the coordinate labelling of states (e.g. the fixed points now correspond to the set $\{(X^5, X^6, X^7, X^8) : X^i \in \{0, \frac{1}{2}\}\}$) and the exchanging of the roles of α and β . Once again, the 16 signs get identified in 4's, and so we end up with the 4 sign degrees of freedom $\epsilon_{f_\beta}(\alpha)$ for $f_\beta = 1 \dots 4$.

Next, let's do $\alpha\beta$. In this case $\alpha\beta$ sends $X^i \rightarrow -X^i$ for all i , and so we have 256 fixed points corresponding to the choices $X^i \in \{0, \frac{1}{2}\}$. α and β clearly act diagonally, γ permutes X^1 and X^5 , and δ permutes X^3 , X^5 and X^7 . We can choose two of these as labels for our states, and in particular choose X^1 and X^3 . This groups our states into 64 sets of 4 — i.e. we have a decomposition $\mathcal{H}_{\alpha\beta} = \oplus_{f_{\alpha\beta}=1}^{64} \mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}$, with each $\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}$ 4 dimensional. It is then easy to see that the actions of $\gamma|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}}$ and $\delta|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}}$ are the same as they were in e.g. the α sector, and so they take the same form (also without discrete torsion phases). For α and β , we may set them both equal (in order to impose $\alpha\beta|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \mathbb{I}_{4 \times 4}$) to:

$$\alpha|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \beta|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \begin{pmatrix} \epsilon_1 & 0 & 0 & 0 \\ 0 & \epsilon_2 & 0 & 0 \\ 0 & 0 & \epsilon_3 & 0 \\ 0 & 0 & 0 & \epsilon_4 \end{pmatrix}, \tag{C.4}$$

so that $\alpha\beta = 1$ when $\epsilon_i^2 = 1$. The commutation constraints set all ϵ_i equal, and we call them $\epsilon_{f_{\alpha\beta}}(\alpha, \beta)$ so that $\alpha|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \beta|_{\mathcal{H}_{\alpha\beta}^{f_{\alpha\beta}}} = \epsilon_{f_{\alpha\beta}}(\alpha, \beta) \cdot \mathbb{I}_{4 \times 4}$ (we use the notation $\epsilon_{f_{\alpha\beta}}(\alpha, \beta)$ to denote the fact that both α and β have discrete torsion signs in the $\alpha\beta$ sector), and we end up with 64 sign degrees of freedom $\epsilon_{f_{\alpha\beta}}(\alpha, \beta)$.

δ sector. Now let's do the δ sector. Here we get 16 fixed points in (X^1, X^3, X^5, X^7) grouped into 2 sets of 8 by the rest of the group, where α permutes X^3 , β permutes X^5 and X^7 , and γ permutes X^1 . We choose to label each of the two 8 dimensional sets by $f_\delta \in \{1, 2\}$, corresponding to either the case $X^5 = X^7$ or $X^5 \neq X^7$. In each of these sectors, we have an 8 dimensional representation $\mathcal{H}_\delta^{f_\delta}$ with basis states labelled by $X^1 \in \{0, 1/2\}$, and $X^3, X^5 \in \{1/4, 3/4\}$. The 8 states $|i\rangle_\delta^{f_\delta} = (X^1, X^3, X^5)$ are:

$$\begin{aligned} |1\rangle_\delta^{f_\delta} &= (0, 1/4, 1/4), & |2\rangle_\delta^{f_\delta} &= (0, 1/4, 3/4), & |3\rangle_\delta^{f_\delta} &= (0, 3/4, 1/4), \\ |4\rangle_\delta^{f_\delta} &= (0, 3/4, 3/4), & |5\rangle_\delta^{f_\delta} &= (1/2, 1/4, 1/4), & |6\rangle_\delta^{f_\delta} &= (1/2, 1/4, 3/4), \\ |7\rangle_\delta^{f_\delta} &= (1/2, 3/4, 1/4), & |8\rangle_\delta^{f_\delta} &= (1/2, 3/4, 3/4). \end{aligned} \quad (\text{C.5})$$

and the group action is:

$$\alpha|_{\mathcal{H}_\delta^{f_\delta}} : \begin{pmatrix} |1\rangle_\delta^{f_\delta} \leftrightarrow |5\rangle_\delta^{f_\delta} \\ |2\rangle_\delta^{f_\delta} \leftrightarrow |6\rangle_\delta^{f_\delta} \\ |3\rangle_\delta^{f_\delta} \leftrightarrow |7\rangle_\delta^{f_\delta} \\ |4\rangle_\delta^{f_\delta} \leftrightarrow |8\rangle_\delta^{f_\delta} \end{pmatrix}, \quad \beta|_{\mathcal{H}_\delta^{f_\delta}} : \begin{pmatrix} |1\rangle_\delta^{f_\delta} \leftrightarrow |2\rangle_\delta^{f_\delta} \\ |3\rangle_\delta^{f_\delta} \leftrightarrow |4\rangle_\delta^{f_\delta} \\ |5\rangle_\delta^{f_\delta} \leftrightarrow |6\rangle_\delta^{f_\delta} \\ |7\rangle_\delta^{f_\delta} \leftrightarrow |8\rangle_\delta^{f_\delta} \end{pmatrix}, \quad \text{and} \quad \gamma|_{\mathcal{H}_\delta^{f_\delta}} : \begin{pmatrix} |1\rangle_\delta^{f_\delta} \leftrightarrow |3\rangle_\delta^{f_\delta} \\ |2\rangle_\delta^{f_\delta} \leftrightarrow |4\rangle_\delta^{f_\delta} \\ |5\rangle_\delta^{f_\delta} \leftrightarrow |7\rangle_\delta^{f_\delta} \\ |6\rangle_\delta^{f_\delta} \leftrightarrow |8\rangle_\delta^{f_\delta} \end{pmatrix}. \quad (\text{C.6})$$

The matrix representations are 8×8 and can easily be constructed, and we find after imposing any commutation relations that any discrete torsion signs vanish.

γ sector. Here we get 16 fixed points with $X^1, X^5 \in \{1/4, 3/4\}$ and $X^2, X^6 \in \{0, 1/2\}$. In this sector, α and δ both permute X^1 and β permutes X^5 — thus, the 16 states get grouped into 4 sets of 4 labelled by the choices of X^2 and X^6 . Choosing a label $f_\gamma \in \{1, \dots, 4\}$ to represent this choice, the 4 states within each sub sector correspond to the 2-tuples (X^1, X^5) :

$$\begin{aligned} |1\rangle_\gamma^{f_\gamma} &= (1/4, 1/4), & |2\rangle_\gamma^{f_\gamma} &= (3/4, 1/4), \\ |3\rangle_\gamma^{f_\gamma} &= (1/4, 3/4), & |4\rangle_\gamma^{f_\gamma} &= (3/4, 3/4), \end{aligned} \quad (\text{C.7})$$

with:

$$\alpha|_{\mathcal{H}_\gamma^{f_\gamma}}, \delta|_{\mathcal{H}_\gamma^{f_\gamma}} : \begin{pmatrix} |1\rangle_\gamma^{f_\gamma} \leftrightarrow |2\rangle_\gamma^{f_\gamma} \\ |3\rangle_\gamma^{f_\gamma} \leftrightarrow |4\rangle_\gamma^{f_\gamma} \end{pmatrix}, \quad \text{and} \quad \beta|_{\mathcal{H}_\gamma^{f_\gamma}} : \begin{pmatrix} |1\rangle_\gamma^{f_\gamma} \leftrightarrow |3\rangle_\gamma^{f_\gamma} \\ |2\rangle_\gamma^{f_\gamma} \leftrightarrow |4\rangle_\gamma^{f_\gamma} \end{pmatrix}. \quad (\text{C.8})$$

It is then clear that $\alpha\delta$ acts diagonally. We can then turn these into matrix representations, and after imposing any constraints we find:

$$\alpha|_{\mathcal{H}_\gamma^{f_\gamma}} = \begin{pmatrix} H & 0 \\ 0 & H \end{pmatrix}, \quad \delta|_{\mathcal{H}_\gamma^{f_\gamma}} = \epsilon_{f_\gamma}(\delta) \cdot \alpha|_{\mathcal{H}_\gamma^{f_\gamma}} \quad \text{and} \quad \beta|_{\mathcal{H}_\gamma^{f_\gamma}} = \begin{pmatrix} 0 & \mathbb{I}_{2 \times 2} \\ \mathbb{I}_{2 \times 2} & 0 \end{pmatrix}. \quad (\text{C.9})$$

So we get 4 sign degrees of freedom $\epsilon_{f_\gamma}(\delta)$. Only the elements $\alpha\delta|_{\mathcal{H}_\gamma^{f_\gamma}} = \alpha\delta\gamma|_{\mathcal{H}_\gamma^{f_\gamma}} = \epsilon_{f_\gamma}(\delta) \cdot \mathbb{I}_{4 \times 4}$ act diagonally, and so we should look at these sectors next.

$\alpha\delta$ and $\alpha\delta\gamma$. The actions of $\alpha\delta$ and $\alpha\delta\gamma$ are:

	X^1	X^2	X^3	X^4	X^5	X^6	X^7	X^8
$\alpha\delta$	+	−	$+\frac{1}{2}$	−	$-\frac{1}{2}$	+	$-\frac{1}{2}$	+
$\alpha\delta\gamma$	$-\frac{1}{2}$	+	$+\frac{1}{2}$	−	+	−	$-\frac{1}{2}$	+

So, for $\alpha\delta$ we have fixed points in X^2, X^4, X^5 and X^7 , and the action $X^3 \rightarrow X^3 + 1/2$ means the winding number n_3 takes values in $\mathbb{Z} + 1/2$. In the lowest energy state we must take $n_3 = \pm 1/2$, and here α and δ both permute n_3 , γ acts diagonally, and β permutes X^5 and X^7 (both $\in \{1/4, 3/4\}$). This action groups the 32 states (16 fixed points/twist fields with two possible winding numbers each) into 8 sets of 4, with a label $f_{\alpha\delta} \in 1 \dots 8$ representing these sectors. Within each, we label states by (X^5, n_3) :

$$\begin{aligned} |1\rangle_{\alpha\delta}^{f_{\alpha\delta}} &= (1/4, +), & |3\rangle_{\alpha\delta}^{f_{\alpha\delta}} &= (3/4, +), \\ |2\rangle_{\alpha\delta}^{f_{\alpha\delta}} &= (1/4, -), & |4\rangle_{\alpha\delta}^{f_{\alpha\delta}} &= (3/4, -). \end{aligned} \quad (\text{C.10})$$

Here, \pm is short for $n_3 = \pm 1/2$. So we find:

$$\alpha|_{\mathcal{H}_{\alpha\delta}^{f_{\alpha\delta}}}, \delta|_{\mathcal{H}_{\alpha\delta}^{f_{\alpha\delta}}} : \left(\begin{array}{l} |1\rangle_{\alpha\delta}^{f_{\alpha\delta}} \leftrightarrow |2\rangle_{\alpha\delta}^{f_{\alpha\delta}} \\ |3\rangle_{\alpha\delta}^{f_{\alpha\delta}} \leftrightarrow |4\rangle_{\alpha\delta}^{f_{\alpha\delta}} \end{array} \right), \quad \text{and} \quad \beta|_{\mathcal{H}_{\alpha\delta}^{f_{\alpha\delta}}} : \left(\begin{array}{l} |1\rangle_{\alpha\delta}^{f_{\alpha\delta}} \leftrightarrow |3\rangle_{\alpha\delta}^{f_{\alpha\delta}} \\ |2\rangle_{\alpha\delta}^{f_{\alpha\delta}} \leftrightarrow |4\rangle_{\alpha\delta}^{f_{\alpha\delta}} \end{array} \right), \quad (\text{C.11})$$

with γ diagonal. After imposing representation constraints, only the generator γ retains a phase and has a matrix representation of $\gamma = \epsilon_{f_{\alpha\delta}}(\gamma) \cdot \mathbb{I}_{4 \times 4}$. As required, the discrete torsion sign is independent of the winding number, and the 16 sign degrees of freedom get organized across the twist fields into 8 sets of 2.

For $\alpha\gamma\delta$, the twist fields have labels $(X^1, X^4, X^6, X^7, n_3)$ with $X^4, X^6 \in \{0, 1/2\}$, $X^1, X^7 \in \{1/4, 3/4\}$ and $n_3 \in \{1/2, -1/2\}$. α and δ permute X^1, X^3 and n_3 , while β permutes X^7 and γ acts diagonally. This orders the 32 states into 4 sets of 8, labelled by a number $f_{\alpha\delta\gamma} \in \{1, \dots, 4\}$ and by (X^1, X^7, n_3) within each set:

$$\begin{aligned} |1\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (1/4, 1/4, +), & |2\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (1/4, 3/4, +), & |3\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (3/4, 1/4, +), \\ |4\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (3/4, 3/4, +), & |5\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (1/4, 1/4, -), & |6\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (1/4, 3/4, -), \\ |7\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (3/4, 1/4, -), & |8\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} &= (3/4, 3/4, -), \end{aligned} \quad (\text{C.12})$$

The group action is:

$$\alpha|_{\mathcal{H}_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}}}, \delta|_{\mathcal{H}_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}}} : \left(\begin{array}{l} |1\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |7\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |2\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |8\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |3\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |5\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |4\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |6\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \end{array} \right), \quad \text{and} \quad \beta|_{\mathcal{H}_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}}} : \left(\begin{array}{l} |1\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |2\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |3\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |4\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |5\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |6\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \\ |7\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \leftrightarrow |8\rangle_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}} \end{array} \right). \quad (\text{C.13})$$

After organizing the phases, we find that $\delta = \epsilon_{f_{\alpha\gamma\delta}}(\delta, \gamma) \cdot \alpha$, $\gamma = \epsilon_{f_{\alpha\gamma\delta}}(\delta, \gamma) \cdot \mathbb{I}_{8 \times 8}$ and β has no phases (again, $\epsilon_{f_{\alpha\gamma\delta}}(\delta, \gamma)$ emphasizes that both γ and δ get signs in this sector). This way, $\alpha\delta\gamma|_{\mathcal{H}_{\alpha\delta\gamma}^{f_{\alpha\delta\gamma}}} = \mathbb{I}_{8 \times 8}$. Once again the discrete torsion phases are winding independent, and the 16 signs get organized into 8 sets of 2 across the twist fields labelled by $f_{\alpha\gamma\delta}$.

$\alpha\gamma, \beta\gamma, \delta\gamma, \beta\delta, \alpha\beta\gamma, \alpha\beta\delta$ and $\alpha\beta\delta\gamma$. For the rest of the composite group elements, the analysis follows the above structure and we find that imposing commutation constraints or equalities such as $\alpha\gamma|_{\mathcal{H}_{\alpha\gamma}^{f_{\alpha\gamma}}} \stackrel{!}{=} \mathbb{1}$ allow us to absorb all discrete torsion phases. So, no new constraints will arise here.

Constraints. In summary, we have found that there are discrete torsion phases in the $\alpha, \beta, \gamma, \alpha\beta, \alpha\delta$ and $\alpha\delta\gamma$ sectors. We would now like to understand how modular invariance relates these phases.

First, let us consider the $\gamma, \alpha\delta$ and $\alpha\delta\gamma$ twisted sectors. After taking into account the identification of discrete torsion phases, we find:

$$\begin{aligned} Z_{\gamma;\alpha\delta}(-1/\tau) = Z_{\alpha\delta;\gamma}(\tau) &\Rightarrow 4 \sum_{f_\gamma=1}^4 \epsilon_{f_\gamma}(\delta) = 2 \sum_{f_{\alpha\delta}=1}^8 \epsilon_{f_{\alpha\delta}}(\gamma) \\ Z_{\gamma;\alpha\delta\gamma}(-1/\tau) = Z_{\alpha\delta\gamma;\gamma}(\tau) &\Rightarrow 4 \sum_{f_\gamma=1}^4 \epsilon_{f_\gamma}(\delta) = 4 \sum_{f_{\alpha\delta\gamma}=1}^4 \epsilon_{f_{\alpha\delta\gamma}}(\delta, \gamma) \end{aligned} \quad . \quad (\text{C.14})$$

As a result, after a possible reordering of the labels f_g we may set $f_\gamma = f_{\alpha\delta\gamma} = f_{\alpha\delta}$ and get:

$$\epsilon_{f_\gamma} \equiv \epsilon_{f_\gamma}(\delta) = \epsilon_{f_\gamma}(\delta, \gamma) = \epsilon_{f_\gamma+4}(\gamma) . \quad (\text{C.15})$$

Now consider the case of α, β and $\alpha\beta$. For α and β we find:

$$4 \sum_{f_\alpha=1}^4 \epsilon_{f_\alpha}(\beta) = 4 \sum_{f_\beta=1}^4 \epsilon_{f_\beta}(\alpha) . \quad (\text{C.16})$$

So, after setting $f_\alpha = f_\beta$:

$$\epsilon_{f_\alpha} \equiv \epsilon_{f_\alpha}(\beta) = \epsilon_{f_\alpha}(\alpha) . \quad (\text{C.17})$$

However, consider the case of $\alpha\beta$. Taking into account the identification of phases across the twist fields, we find:

$$4 \sum_{f_{\alpha\beta}=1}^{64} \epsilon_{f_{\alpha\beta}}(\alpha, \beta) = 16 \left(4 \sum_{f_\alpha=1}^4 \epsilon_{f_\alpha} \right) \Rightarrow \sum_{f_{\alpha\beta}=1}^{64} \epsilon_{f_{\alpha\beta}}(\alpha, \beta) = 16 \sum_{f_\alpha=1}^4 \epsilon_{f_\alpha} ,$$

where we also used equation (C.17). Thus, we may set:

$$\epsilon_{f_\alpha} = \epsilon_{f_{\alpha\beta}+4k}(\alpha, \beta) , \quad (\text{C.18})$$

for $k = 0, \dots, 15$ and $f_\alpha = 1, \dots, 4$.

Open Access. This article is distributed under the terms of the Creative Commons Attribution License ([CC-BY 4.0](https://creativecommons.org/licenses/by/4.0/)), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.

References

- [1] S.L. Shatashvili and C. Vafa, *Superstrings and manifold of exceptional holonomy*, *Selecta Math.* **1** (1995) 347 [[hep-th/9407025](#)] [[INSPIRE](#)].
- [2] J.M. Figueroa-O'Farrill, *A note on the extended superconformal algebras associated with manifolds of exceptional holonomy*, *Phys. Lett.* **B 392** (1997) 77 [[hep-th/9609113](#)] [[INSPIRE](#)].
- [3] D.D. Joyce, *Compact Riemannian 7-manifolds with holonomy g_2 . I*, *J. Diff. Geom.* **43** (1996) 291.
- [4] D.D. Joyce, *Compact Riemannian 7-manifolds with holonomy g_2 . II*, *J. Diff. Geom.* **43** (1996) 329.
- [5] D.D. Joyce, *Compact 8-manifolds with holonomy $Spin(7)$* , *Inv. Math.* **123** (1996) 507.
- [6] B.S. Acharya, *On mirror symmetry for manifolds of exceptional holonomy*, *Nucl. Phys.* **B 524** (1998) 269 [[hep-th/9707186](#)] [[INSPIRE](#)].
- [7] B.S. Acharya, *Dirichlet Joyce manifolds, discrete torsion and duality*, *Nucl. Phys.* **B 492** (1997) 591 [[hep-th/9611036](#)] [[INSPIRE](#)].
- [8] A. Strominger, S.-T. Yau and E. Zaslow, *Mirror symmetry is T duality*, *Nucl. Phys.* **B 479** (1996) 243 [[hep-th/9606040](#)] [[INSPIRE](#)].
- [9] M.R. Gaberdiel and P. Kaste, *Generalized discrete torsion and mirror symmetry for G_2 manifolds*, *JHEP* **08** (2004) 001 [[hep-th/0401125](#)] [[INSPIRE](#)].
- [10] W.-y. Chuang, *A note on mirror symmetry for manifolds with $Spin(7)$ holonomy*, *J. Phys.* **A 43** (2010) 235403 [[hep-th/0406151](#)] [[INSPIRE](#)].
- [11] G. Papadopoulos and P.K. Townsend, *Compactification of $D = 11$ supergravity on spaces of exceptional holonomy*, *Phys. Lett.* **B 357** (1995) 300 [[hep-th/9506150](#)] [[INSPIRE](#)].
- [12] A. Kovalev, *Twisted connected sums and special Riemannian holonomy*, *J. Reine Angew. Math.* **565** (2003) 125.
- [13] A. Corti, M. Haskins, J. Nordström and T. Pacini, *Asymptotically cylindrical Calabi-Yau 3-folds from weak Fano 3-folds*, *Geom. Topol.* **17** (2013) 1955.
- [14] A. Corti, M. Haskins, J. Nordström and T. Pacini, *G_2 -manifolds and associative submanifolds via semi-Fano 3-folds*, *Duke Math. J.* **164** (2015) 1971 [[arXiv:1207.4470](#)] [[INSPIRE](#)].
- [15] A.P. Braun and M. Del Zotto, *Mirror symmetry for G_2 -manifolds: twisted connected sums and dual tops*, *JHEP* **05** (2017) 080 [[arXiv:1701.05202](#)] [[INSPIRE](#)].
- [16] A.P. Braun and M. Del Zotto, *Towards generalized mirror symmetry for twisted connected sum G_2 manifolds*, *JHEP* **03** (2018) 082 [[arXiv:1712.06571](#)] [[INSPIRE](#)].
- [17] D. Joyce, *Compact manifolds with special holonomy*, Oxford mathematical monographs, Oxford University Press, Oxford U.K. (2000).
- [18] V.V. Batyrev, *Dual polyhedra and mirror symmetry for Calabi-Yau hypersurfaces in toric varieties*, *J. Alg. Geom.* **3** (1994) 493 [[alg-geom/9310003](#)] [[INSPIRE](#)].
- [19] A.P. Braun, *Tops as building blocks for G_2 manifolds*, *JHEP* **10** (2017) 083 [[arXiv:1602.03521](#)] [[INSPIRE](#)].
- [20] P.S. Aspinwall and D.R. Morrison, *String theory on $K3$ surfaces*, [hep-th/9404151](#) [[INSPIRE](#)].

- [21] A.P. Braun and S. Schäfer-Nameki, *Compact, singular G_2 -holonomy manifolds and M /Heterotic/ F -theory duality*, *JHEP* **04** (2018) 126 [[arXiv:1708.07215](#)] [[INSPIRE](#)].
- [22] J. Halverson and D.R. Morrison, *The landscape of M -theory compactifications on seven-manifolds with G_2 holonomy*, *JHEP* **04** (2015) 047 [[arXiv:1412.4123](#)] [[INSPIRE](#)].
- [23] J. Halverson and D.R. Morrison, *On gauge enhancement and singular limits in G_2 compactifications of M -theory*, *JHEP* **04** (2016) 100 [[arXiv:1507.05965](#)] [[INSPIRE](#)].
- [24] T.C. da C. Guio, H. Jockers, A. Klemm and H.-Y. Yeh, *Effective Action from M -theory on Twisted Connected Sum G_2 -Manifolds*, *Commun. Math. Phys.* **359** (2018) 535 [[arXiv:1702.05435](#)] [[INSPIRE](#)].
- [25] A.P. Braun et al., *Infinitely many $M2$ -instanton corrections to M -theory on G_2 -manifolds*, *JHEP* **09** (2018) 077 [[arXiv:1803.02343](#)] [[INSPIRE](#)].
- [26] B.S. Acharya, A.P. Braun, E.E. Svanes and R. Valandro, *Counting associatives in compact G_2 orbifolds*, *JHEP* **03** (2019) 138 [[arXiv:1812.04008](#)] [[INSPIRE](#)].
- [27] A.P. Braun and S. Schäfer-Nameki, *$Spin(7)$ -manifolds as generalized connected sums and $3d$ $\mathcal{N} = 1$ theories*, *JHEP* **06** (2018) 103 [[arXiv:1803.10755](#)] [[INSPIRE](#)].
- [28] M.-A. Fiset, *Superconformal algebras for twisted connected sums and G_2 mirror symmetry*, *JHEP* **12** (2018) 011 [[arXiv:1809.06376](#)] [[INSPIRE](#)].
- [29] C. Vafa, *Modular invariance and discrete torsion on orbifolds*, *Nucl. Phys.* **B 273** (1986) 592.
- [30] E. Witten, *Supersymmetry and Morse theory*, *J. Diff. Geom.* **17** (1982) 661 [[INSPIRE](#)].
- [31] P.S. Aspinwall, D.R. Morrison and M. Gross, *Stable singularities in string theory*, *Commun. Math. Phys.* **178** (1996) 115 [[hep-th/9503208](#)] [[INSPIRE](#)].
- [32] M. Gross, *Special Lagrangian Fibrations II: geometry*, [math/9809072](#) (1998).
- [33] D. Crowley, S. Goette and J. Nordström, *An analytic invariant of G_2 manifolds*, [arXiv:1505.02734](#) [[INSPIRE](#)].
- [34] C. Vafa and E. Witten, *On orbifolds with discrete torsion*, *J. Geom. Phys.* **15** (1995) 189 [[hep-th/9409188](#)] [[INSPIRE](#)].
- [35] D. Joyce, *A new construction of compact 8-manifolds with holonomy $Spin(7)$* , *J. Diff. Geom.* **53** (1999) 89 [[math/9910002](#)] [[INSPIRE](#)].
- [36] D. Joyce and S. Karigiannis, *A new construction of compact torsion-free G_2 -manifolds by gluing families of Eguchi-Hanson spaces*, [arXiv:1707.09325](#).
- [37] R. Gopakumar and S. Mukhi, *Orbifold and orientifold compactifications of F -theory and M -theory to six-dimensions and four-dimensions*, *Nucl. Phys.* **B 479** (1996) 260 [[hep-th/9607057](#)] [[INSPIRE](#)].