

1 **Reassessment of hydrate destabilization mechanisms offshore west**

2 **Svalbard confirms link to recent ocean warming**

3 **Akash Trivedi^{1*}, Sudipta Sarkar¹, Hector Marin-Moreno², Timothy A. Minshull³, Pippa**
4 **L. Whitehouse⁴, Utpal Singh¹**

5 **Affiliations:**

6 ¹Department of Earth and Climate Science, Indian Institute of Science Education and
7 Research Pune, Pune, 411008, India

8 ²Norwegian Geotechnical Institute, PB 3930 Ullevål Stadion, NO-08906 Oslo, Norway

9 ³School of Ocean and Earth Science, University of Southampton, Waterfront Campus,
10 European Way, Southampton SO14 3ZH, UK

11 ⁴Department of Geography, Durham University, Durham, DH1 3LE, UK

12 *Correspondence to: Akash Trivedi (trivedi.akash@students.iiserpune.ac.in)

13 **Key Points:**

- 14
- 15 • We examine the effect of the past 11000 years' bottom water temperature and
16 relative sea level change on Arctic gas hydrate dynamics.
 - 17 • Relative sea-level fall over the past 8 ka thins the hydrate layer, but several
18 bottom water warming pulses also contribute to its dissociation.
 - 19 • Simulation confirms present-day gas seeps and observed chloride anomaly can
20 be explained by thermal dissociation of hydrate over 1978-2016.
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30 **Abstract**

31 The stability of methane hydrates at the feather edge of hydrate stability on the upper
32 continental slope (UCS) is prone to ocean warming and relative sea-level (RSL) change.
33 West of Svalbard, methane seeps on the UCS were initially proposed to result from hydrate
34 destabilization resulting from four decades of warming of Atlantic bottom water.
35 Alternatively, it has been proposed that hydrate dissociation was triggered by RSL fall due to
36 isostatic rebound over the past 8000 years rather than recent bottom water temperature
37 (BWT) rise. Here we address these two contrasting hypotheses by simulating the impact of
38 11000 years of BWT and RSL change on hydrates located at the UCS off west Svalbard. Our
39 numerical simulation considers multiphase fluid and heat flow coupled with hydrate
40 formation and dissociation. We used two reconstructions of local ice history (UiT and ICE-
41 6G_C) that predict contrasting results for the local sea-level history. Over the past 8000
42 years, the UiT model predicts a fall in RSL on the UCS, while the ICE-6G_C model, which
43 provides a better fit to nearby coastal RSL observations, predicts a continuous rise. Our
44 modeling shows that whilst long-term RSL fall would progressively thin the region of
45 hydrate stability, the abrupt rise in BWT enhances hydrate dissociation. Even in the model
46 with an RSL rise, the increase in BWT causes hydrate destabilization and pore water
47 freshening that matches observations. We conclude that recent ocean warming plays a critical
48 role in hydrate dissociation west of Svalbard regardless of the longer term sea-level history.

49 **Plain language summary:**

50 Methane hydrate is an ice-like substance that is stable at the low temperatures and high
51 pressures found in the mud and sand that builds up at the edges of oceans. Hydrate can melt
52 and vent methane gas into the ocean if the ocean warms or sea-level falls. Drilling in a region
53 of such venting west of the Arctic Archipelago of Svalbard in 2016 confirmed that it was
54 caused by the melting of hydrate, but the cause of this melting is debated. Proposed causes
55 are ocean warming during 1978–2016, or relative sea-level fall since 8000 years ago as the
56 ocean floor near Svalbard rose due to the removal of the weight of the melting Svalbard ice
57 sheet. We used a computer model to study the impact of changes in ocean temperature and
58 sea level on this hydrate since 11000 years ago. We found that sea-level fall causes the
59 hydrate to melt, but there are periods of colder ocean temperatures when it can re-form. Sea-
60 level rise stabilizes the hydrate, but ocean warming can override its effect, leading to the
61 hydrate melting that is inferred during 1978–2016. We conclude that ocean warming played a
62 critical role in hydrate melting.

63 **1. Introduction**

64 The effects of ocean warming are clearly visible across the Arctic (Hassol, 2004), as
65 evidenced by the disappearance of sea ice (Comiso et al., 2008; Piechura & Walczowski,
66 2009) and enhanced glacier melting (Osterkamp, 2005; Nagornov et al., 2006; Anthony et al.,
67 2012). Dissociation of methane hydrates has also been linked with past climate warming
68 (Nisbet, 1990; Dickens, 2003; Kennett et al., 2003), and there is concern about the stability of
69 methane hydrate deposits trapped in Arctic marine sediments (Kretschmer et al., 2015).
70 Hydrate at the feather edge of its stability zone, where the base of this zone outcrops on the
71 upper continental slope, are particularly vulnerable to thermal destabilization. The marine
72 sediments of the Fram Strait, the Arctic-Atlantic gateway located in the western Svalbard
73 region, host widespread methane hydrates (Figure 1a, Sarkar et al., 2012; Dumke et al., 2016;
74 Minshull et al., 2020). A marine geophysical expedition in 2008 discovered plumes of
75 methane bubbles emanating from the seafloor close to and further up-slope of the present-day

76 upper limit of methane hydrate stability on the eastern margin of the Fram Strait (Westbrook
77 et al., 2009). Here warm and saline Atlantic water brings heat into the Arctic Ocean through
78 the Fram Strait via the West Spitsbergen Current. This Atlantic water is inferred to have
79 warmed the bottom water temperature (BWT) from 2° to 3° C over the past four decades,
80 resulting in the migration of the feather edge of methane hydrate stability from 370 m to
81 410 m depth (Westbrook et al., 2009). Westbrook et al. (2009) proposed that these methane
82 bubble plumes indicate thermal degradation of the hydrates at the landward limit of hydrate
83 stability.

84 Several subsequent marine surveys consistently found methane seeps (Sahling et al., 2014;
85 Mau et al., 2017; Veloso-Alarcon et al., 2019) between 370 m and 410 m water depths on the
86 upper slope (Figure 1b). The gas flare density and methane release rates are reduced during
87 the winter when the temperatures are substantially lower than in summer (Veloso-Alarcon et
88 al., 2019; Ferré et al., 2020). This suggests that seasonal fluctuations in BWT can lead to
89 repeated episodes of hydrate formation and dissociation (Berndt et al., 2014; Veloso-Alarcon
90 et al., 2019; Ferré et al., 2020). However, the link between methane seepage and
91 anthropogenic ocean warming has been questioned because the presence of seep carbonates
92 in this area suggests that seepage has been active for >8,000 yr (Berndt et al., 2014).
93 Numerical simulation of the response of gas hydrates to BWT changes was previously carried
94 out considering past and future ocean temperature variations (Reagan et al., 2011; Thatcher et
95 al., 2013; Marín-Moreno et al., 2013; Marín-Moreno et al., 2015). Results from these studies
96 indicated that for hydrate shallower than a few metres beneath the seabed, bottom water
97 warming since the late 1970s has produced gas from hydrate dissociation. However, these
98 studies did not account for the effect of relative sea level (RSL) change on hydrate stability.

99 Hydrate dissociation off west Svalbard was first geochemically confirmed by Wallmann et al.
100 (2018). Pore water samples collected from sediments during a “MeBo” seafloor drilling
101 experiment conducted in 2016 at a water depth of 391 m (Figure 1b) showed a decline in
102 chloride concentration at 15–22 m below the seabed, indicating freshening caused by hydrate
103 dissociation. Based on the University of Tromsø (UiT) RSL model (Patton et al., 2017),
104 Wallmann et al. (2018) inferred that approximately 8000 years ago, the glacial isostatic
105 rebound rate exceeded the rate of eustatic sea-level rise. The resulting RSL fall and
106 consequent drop in hydrostatic pressure led to hydrate destabilization. In their model, hydrate
107 dissociation was completed 1000 years before the present (BP, where ‘Present’ is considered
108 as 2016). Hydrate could re-form in the model during colder BWT conditions in the late 1970s
109 (Ferré et al., 2012; Wallmann et al., 2018). Wallmann et al. (2018) concluded that a 1°C
110 linear rise in BWT over 1980–2016 would be insufficient to result in hydrate dissociation.
111 However, in their modeling they did not account for the decadal-scale temperature
112 fluctuations observed in the measured records. Such variations can impact hydrate dynamics
113 and methane flow at the seabed, and therefore, should be included in a realistic assessment of
114 whether recent bottom water warming could have caused the observed hydrate dissociation.

115 Glacial retreat across Svalbard took place about 13000 years BP (Landvik et al., 2008),
116 causing rapid seafloor uplift and a prominent relative sea-level (RSL) drop over 8–12 ka
117 along the western Svalbard coast (Forman et al., 2004). RSL change at the MeBo drill site is
118 not constrained by data and must be estimated by geodynamic modeling. Comparison of
119 observed and predicted RSL change based on various deglacial scenarios has revealed the
120 inadequacy of some models in reproducing observations of RSL change throughout the
121 Barents Sea region (Auriac et al., 2016). The ICE-6G_C scenario correctly predicts a fall in
122 RSL over 8–12 ka in west Svalbard (Auriac et al., 2016) but, in contrast to the UiT model

123 employed by Wallmann et al. (2018), it predicts RSL rise at the MeBo site during the past
124 8000 years (see Section 3.3). Ultimately it is not possible to confirm which scenario is
125 correct. Therefore, we consider both the UiT and ICE-6G_C RSL scenarios to examine the
126 role of pressure in hydrate dynamics west of Svalbard.

127 Here we address the limitations of previous studies and place further constraints on hydrate
128 dissociation processes through coupled heat and mass transfer numerical simulations of the
129 impact of 11000 years of ocean temperature and RSL change on hydrate located at the upper
130 continental slope off west Svalbard. We have attempted to reproduce the present-day chloride
131 anomaly observed in the sediment core collected at 391 m water depth (Wallmann et al.,
132 2018) by varying the initial hydrate saturation and considering realistic RSL and BWT
133 reconstructions, including decadal and seasonal temperature fluctuations. Based on this
134 modeling, we conclude that rising BWT plays a dominant role in hydrate destabilization,
135 potentially overprinting the effect of any ongoing RSL rise or fall.

136 **2. The West Svalbard Margin**

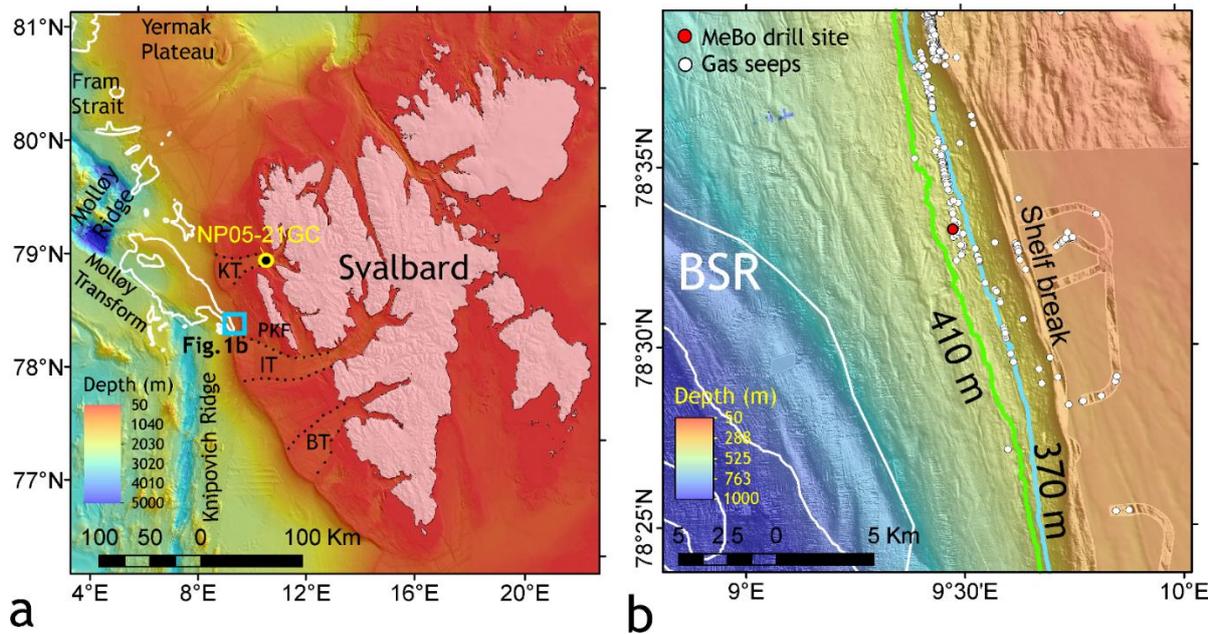
137 The west Svalbard continental margin has experienced repeated Pliocene–Pleistocene
138 glaciations and deglaciations that affected sea level, sedimentation, and erosion patterns
139 (Landvik et al., 2008). Fast-flowing ice streams carved out the Kongsfjorden, Isfjorden, and
140 Bellsund cross-shelf troughs (KT, IT, and BT, Figure 1a). The present-day shelf break
141 represents the maximum extent reached by the ice sheet and grounded ice streams (Solheim
142 et al., 1996; Knies et al., 2009). Our study area (Figure 1b) is located in the inter-fan region
143 between the Kongsfjorden and the Isfjorden cross-shelf troughs (Sarkar et al., 2011). Using
144 multi-channel reflection seismic data, Sarkar et al. (2012) identified a gas hydrate bottom-
145 simulating reflector (BSR) (Figures 1a and 1b) beneath the upper continental slope. The BSR
146 represents the boundary between gas hydrates above and free gas below (Shipley et al.,
147 1979). There is no evidence of a BSR in water depths shallower than ~600 m (Sarkar et al.,
148 2012); instead, there is seismic evidence of gas accumulations (400–700 m water depths), as
149 evidenced by low-velocity pockets and negative-polarity bright spots within the top 75 m
150 sediment column below the seabed (Chabert et al., 2011). In addition, regions of anomalously
151 high electrical resistivity suggest the presence of hydrate within its stability field (Goswami
152 et al., 2016).

153 The Pliocene–Pleistocene sediments on the shelf and upper continental slope are
154 predominantly glaciogenic in origin. The marine hemipelagic sediments are interbedded with
155 glacial debris flow units at the upper slope (Sarkar et al., 2011, 2012). The MeBo drilling
156 experiment drilled a mixture of hemipelagic sediments and poorly-sorted glacial debris
157 consisting of a broad range of grain sizes, from clay to sand with gravel to pebble-sized clasts
158 in variable amounts (Wallmann et al., 2018). During the deglaciations, meltwater discharge
159 on the upper continental slope resulted in turbidity currents cascading down the slope,
160 depositing sand and silt-rich sediments (Vorren et al., 1998). The lithological heterogeneity
161 strongly influences fluid flow below the methane seepage sites, for example, the flow of gas
162 is impeded by the glaciogenic debris. Lateral fluid migration occurs along the sand and silt-
163 rich permeable sediments, and near-vertical flow occurs through fractures within poorly
164 stratified, low-permeability glaciogenic sediments (Haacke & Westbrook, 2006; Haacke et
165 al., 2009; Sarkar et al., 2012; Thatcher et al., 2013).

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170 **Figure 1.** West Svalbard continental margin with bathymetry compiled from various sources,
 171 such as Norwegian Hydrographic Survey, swath-bathymetry from JR211 cruise (Sarkar et al.,
 172 2012), and International Bathymetric Chart of the Arctic Ocean (Jakobsson et al., 2020). (a)
 173 The Kongsfjorden (KT), Isfjorden (IT), and Bellsund (BT) cross-shelf troughs are the palaeo-
 174 ice stream pathways dissecting the shelf. The study area is located between the KT and IT
 175 and west of Prins Karl Forland (PKF). The solid white line represents the extent of the
 176 seismic bottom simulating reflector (BSR) (Sarkar et al., 2012; Dumke et al., 2016). The
 177 hydrate province is bounded by the Molløy Transform Fault that links the Knipovich Ridge
 178 and the Molløy Ridge. The sediment core NP05-21GC (yellow dot) is located at 327 m water
 179 depth on the Kongsfjorden shelf west of Svalbard (Rasmussen et al., 2014). (b) Gas seeps on
 180 the upper continental slope are located between 370 and 410 m water depths (Sarkar et al.,
 181 2012), which correspond to the theoretical limit of methane hydrate stability at 2°C and 3°C
 182 bottom water temperatures respectively. The MeBo drill site is located at 391 m water depth.

183 3. Data and Method

184 3.1. Modeling approach

185 We performed 1D numerical simulations of vertical fluid and heat flow coupled with hydrate
 186 formation and dissociation using the TOUGH+Hydrate code v1.5 (Moridis et al., 2014). We
 187 used an equilibrium hydrate formation and dissociation model with three mass components,
 188 i.e., methane, water, and salt (NaCl), partitioned into four possible phases (gas hydrate, ice,
 189 gas, and liquid). The model considers conductive and convective heat transport, pressure-
 190 driven transport of water and methane gas (Darcy's flow), advective and diffusive transport of
 191 dissolved methane and salt in water, and latent heat of hydrate formation and dissociation.

192 Model initialization at 11200 yr BP was done assuming a hydrostatic pressure gradient,
 193 constant heat flow at the base of the model of $10 \times 10^{-2} \text{ W m}^{-2}$ and water saturated sediment
 194 thermal conductivity of $2 \text{ W m}^{-1} \text{ K}^{-1}$ (Riedel et al., 2018), generating a geothermal gradient of
 195 50° C/km , 3.5 wt% pore water salinity and appropriate BWT and pressure in the top cell of
 196 the model (Figure 2a). We used the BWT and pressure values from the reconstructed curves
 197 discussed in sections 3.2 and 3.3. Model initialization is done based on various BWT and
 198 RSL scenarios (Supplementary Figure 1). The model has a constant cell thickness of 0.5 m

199 from 1 mm to a maximum depth of 750 m and assumes a present-day water depth of 391 m.
200 The topmost cell where the top boundary conditions are applied has a thickness of 1 mm and
201 represents the ocean bottom.

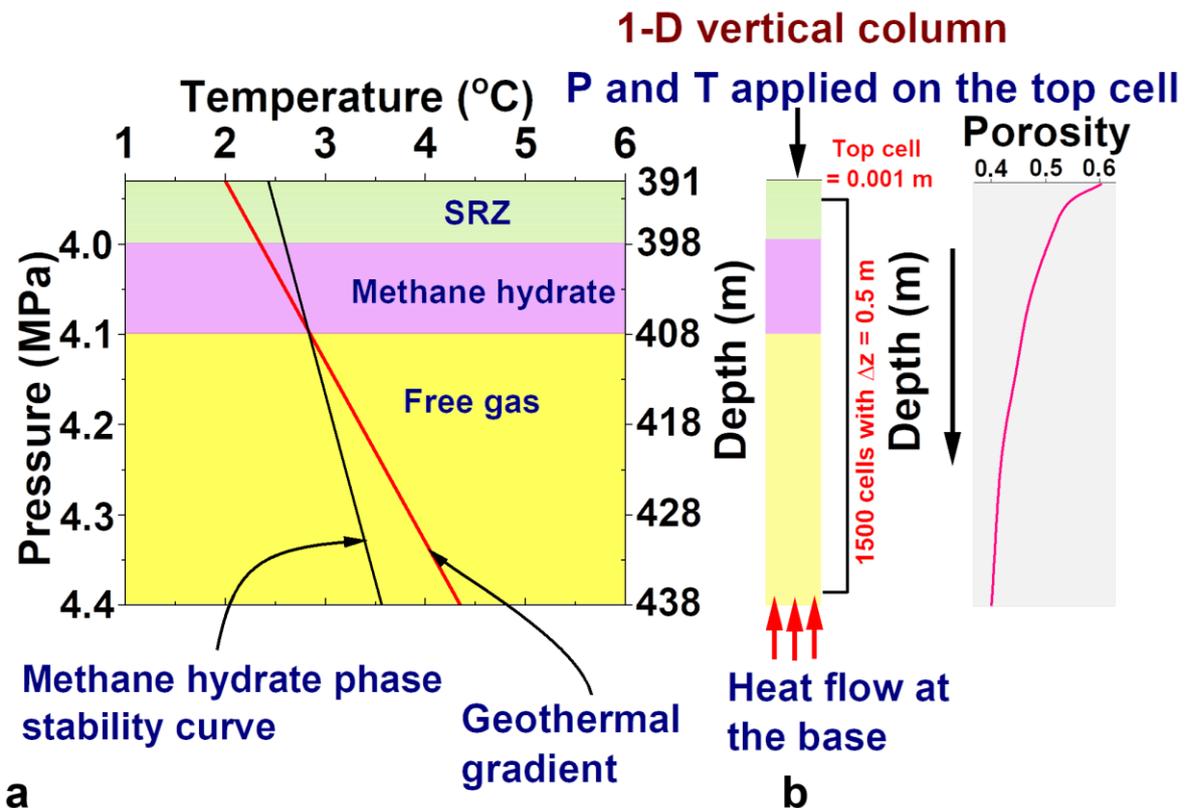
202 The base of the sulfate reduction zone at the MeBo site is 5–7 m (Thatcher et al., 2013;
203 Marín-Moreno et al., 2013; Marín-Moreno et al., 2015 and Wallmann et al., 2018). To
204 represent this, we assume the top 7 m of the sediment column is hydrate-free in the starting
205 model, but we do not consider that this zone was hydrate-free throughout the simulation. The
206 sulfate reduction zone comprises marine anoxic sediments within which there is a decline in
207 sulfate concentration, and dissolved methane is consumed by anaerobic oxidation (Boetius
208 and Wenzhöfer, 2013; Borowski et al., 1996). We have not modeled this process. The
209 Holocene sedimentation rate is negligible, only 4 cm kyr⁻¹ (Panieri et al., 2016); therefore,
210 sedimentation is not included in the modeling.

211 The initial hydrate layer in our model was placed up to the base of methane hydrate stability
212 zone (MHSZ) (Figure 2b). We considered a range of initial hydrate saturations (e.g., 9–60%)
213 for different scenarios. Since no hydrate was recovered during MeBo drilling experiment,
214 this initial hydrate should completely melt. We assume an initial value for gas saturation
215 below the MHSZ of 2–4% based on geophysical evidence of gas in the inter-fan region in a
216 depth range of 480–1285 m on the upper slope (Chabert et al., 2011). Goswami et al. (2016)
217 inferred a much larger gas saturation (>10%) around the landward edge of the GHSZ based
218 on high resistivity. However, the high resistivity values could also be influenced by the
219 presence of authigenic carbonates. We used 2–4% gas below the hydrate stability zone based
220 on earlier models (Marín-Moreno et al., 2013 and Marín-Moreno et al., 2015). The gas was
221 pure methane (James et al., 2011).

222 We assumed an initial hydrate-free uniform intrinsic (absolute) permeability of 10⁻¹³ m²,
223 which is 2–4 orders of magnitude higher than the permeability of hemipelagic glaciomarine
224 sediments (Table S1). We adopted the value from Thatcher et al. (2013), who argued that
225 shallow glacial sediments could rarely sustain a low intrinsic permeability during rapid gas
226 release from hydrate dissociation because the pore pressure will surpass the lithostatic load
227 within a few years of the onset of hydrate dissociation (Thatcher et al., 2013), subsequently
228 leading to the development of fractures that enhance permeability. Hence this high value of
229 intrinsic permeability simulates the effects of fracture permeability. In the scenario they
230 investigated, Thatcher et al. (2013) found that gas appeared at the seabed in response to
231 recent decadal-scale warming, consistent with observations, while if lower permeabilities
232 were used, gas release into the ocean was delayed. In the presence of hydrate in the pore
233 space, there is a reduction in the intrinsic permeability, which is accounted for in the model
234 by using the Evolving Porous Medium model (Moridis et al., 2014; Table S1). The
235 irreducible gas saturation and water saturation are 2 and 12%, respectively, consistent with
236 laboratory measurements (Liu and Flemings, 2007). We note that, to our knowledge, there is
237 no published laboratory data from sediment cores in our study area to better constrain the
238 irreducible gas saturation. We also incorporated multiphase molecular diffusion. Methane
239 flux by molecular diffusion is slow compared to the free gas flow and advection mechanisms,
240 but the simulation time is long; therefore, transport by molecular diffusion can be important.

241 Changes in BWT and RSL were applied to the model at yearly intervals by changing the
242 temperature and pressure in the top cell. Our model includes the instantaneous response of
243 the sediment column to RSL-induced pressure changes and addresses the process of hydrate
244 dissociation and reformation; for example, the model accounts for the elevated pore-pressure

245 and salinity drop that occur during hydrate dissociation, which can further inhibit melting, as
 246 suggested by Liu and Flemings (2009). A summary of the properties used in the models is
 247 provided in Table S1.



248

249 **Figure 2.** Schematic representation of the model domain at a water depth of 391 m. (a)
 250 Methane hydrate stability shown on the Temperature vs. Pressure (Depth) plot. The thickness
 251 of the methane hydrate stability is calculated using the geothermal gradient (red line) and the
 252 methane hydrate phase stability curve (black line). SRZ denotes the sulfate reduction zone
 253 (green area). Methane hydrate-bearing sediments are represented by the purple, and the free
 254 gas zone by the yellow colour. (b) Schematic representation of a 1-D vertical sediment
 255 column with the boundary conditions. Time-dependent pressure and temperature values are
 256 applied as boundary conditions on the top cell while a constant heat source is maintained at
 257 the bottom cell. The porosity depth profile is shown by the pink curve.

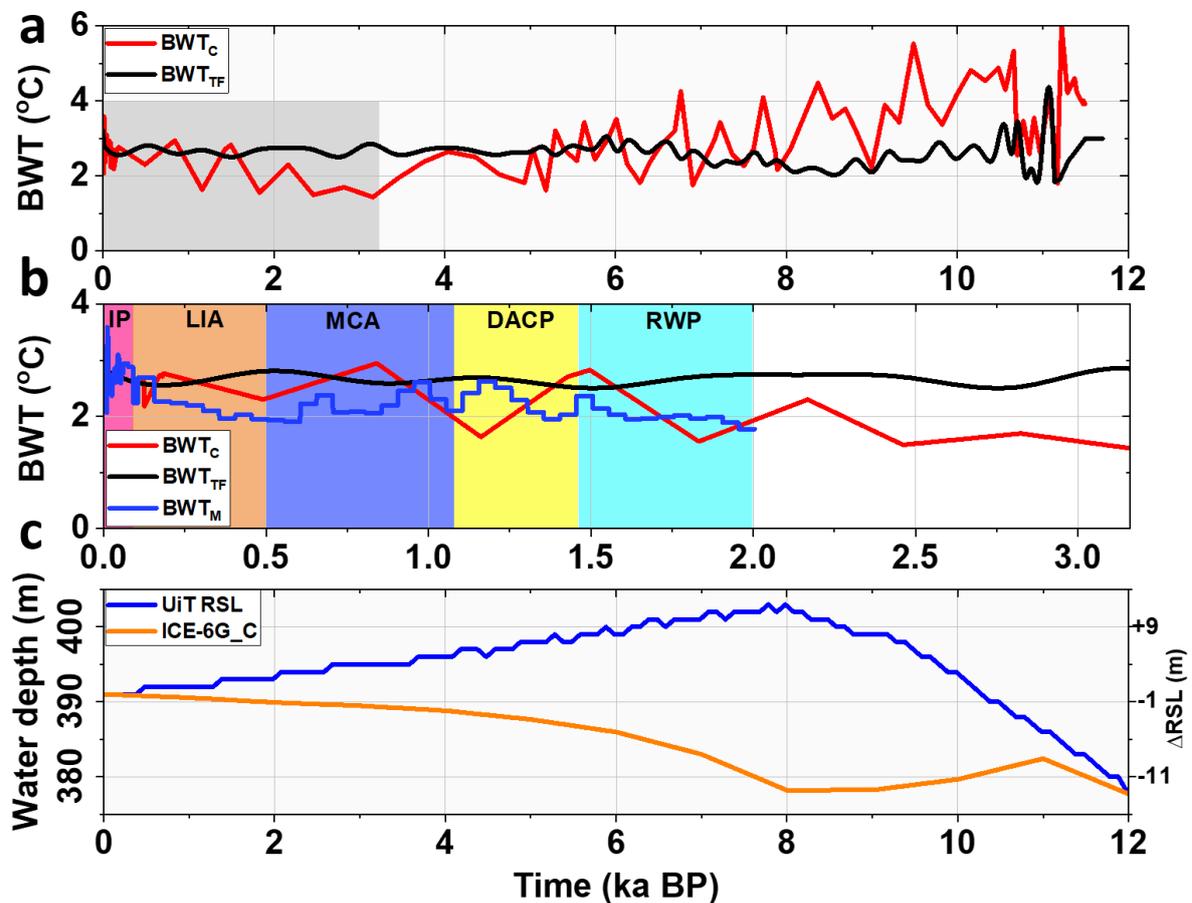
258 3.2. Sea bed temperature history

259 Rasmussen et al. (2014) reconstructed the BWT west of Svalbard over the past 11200 years
 260 based on two proxies – (1) benthic foraminiferal census counts (BWT_{TF}) and (2) benthic
 261 foraminifera δ₁₈O values (BWT_C) from the sediment core NP05-21GC taken at 327 m water
 262 depth on the Kongsfjorden shelf west of Svalbard. NP05-21GC is closest sediment core to the
 263 methane seepage site on the upper continental slope (Figure 1a). It is currently bathed by
 264 warm Atlantic water (Sarkar et al., 2015), and hence we assume that the reconstructed BWT
 265 values well represent past temperature variations of the Atlantic water mass.

266 The BWT_C temperatures, based on δ₁₈O values in *Cibicides lobatulus*, are influenced by
 267 summer conditions since it calcifies during the summer. Temperatures derived using this
 268 approach may therefore be biased warm. The BWT_C curve shows incursions of warmer

269 Atlantic water during the early Holocene (11500–6000 years BP) and a subsequent
 270 temperature dip during 5000–2000 years BP (Figure 3a). BWT_{TF} (also shown in Figure 3a)
 271 shows a smoother curve compared to BWT_C because its reconstruction relies on annual
 272 faunal counts (Rasmussen et al., 2014). Since we do not know which is the most
 273 representative temperature curve, we considered both the temperature series.

274 In addition, Marín-Moreno et al. (2013) compiled measured and estimated temperature
 275 records at 400 m water depth on the West Svalbard continental margin for the past two
 276 millennia. For the years A.D. 1–1900, they used 50-year-running summer temperature means
 277 at 50 m water depth, derived from planktic foraminifera proxy (Spielhagen et al., 2011), and
 278 applied appropriate scaling to predict the temperature at 400 m depth based on the ocean/sea-
 279 ice model NEMO at $1/12^\circ$ resolution (ORCA12) (Madec, 2012). We use the resulting BWT_M
 280 curve of Marín-Moreno et al. (2013) (Figure 3b) in a specific modeling scenario (refer to
 281 section 4.3).



282

283 **Figure 3.** BWT and RSL changes used as model boundary conditions. (a) Bottom water
 284 temperatures derived from shelf core NP05-21GC (327 m depth) based on a benthic
 285 foraminifera transfer function (BWT_{TF}) and oxygen isotopes (BWT_C) (Rasmussen et al.,
 286 2014) for the period 11.7–0 ka. (b) Reconstructed BWT_M values from Marín-Moreno et al.
 287 (2013) during the Roman Warm Period (RWP, 2005–1405 BP), the Dark Ages Cold Period
 288 (DACP, 1405–1105 BP), the Medieval Climate Anomaly (MCA, 1105–505 BP), the Little
 289 Ice Age (LIA, 505–100 years BP), and the Industrial period (IP, 100–present) (Spielhagen et
 290 al., 2011). For the IP, BWT values are from Polyakov (2004) for 1900–1950, Holliday et al.
 291 (2008) for 1950–1975, and CTD-based near-seabed temperature values from Westbrook et al.

292 (2009) and Thatcher et al. (2013) for 1975–2005. (c) Modeled RSL (relative to present) over
293 the past 12000 years at the MeBo site (Fig. 1b) based on the University of Tromsø (UiT) ice
294 load history and Earth rheology model (Auriac et al., 2016) (UiT RSL) and the ICE-6G_C
295 scenario (ICE-6G_C RSL).

296 **3.3. Relative sea-level change**

297 We considered two relative sea-level (RSL) change scenarios during the past 11200 years
298 (Figure 3c) – (1) the RSL values of Wallmann et al. (2018), which are based on the
299 University of Tromsø (UiT) model of Eurasian ice sheet change (Patton et al., 2017), and (2)
300 the RSL change predicted using the global ICE-6G_C ice history model (Peltier et al., 2015;
301 Auriac et al., 2016). The UiT model of ice sheet change was developed using a thermo-
302 mechanical ice model coupled to a climate model (Patton et al., 2017). The UiT RSL curve
303 used here is an extended version of the curve used by Wallmann et al. (2018), which only
304 extended back to 8 ka BP. It was derived by combining the isostatic response to ice load
305 change, calculated using an elastic lithosphere/relaxed asthenosphere model (Le Meur and
306 Huybrechts, 1996), with a eustatic sea-level curve (Waelbroeck et al. 2002).

307 The ICE-6G_C model is the updated version of its predecessor ICE-5G (Argus et al., 2014;
308 Peltier et al., 2015). ICE-6G_C was tuned and tested using a suite of relative sea-level data,
309 GPS observations of vertical motion, and time-dependent gravity change. Auriac et al. (2016)
310 found that ICE-6G_C provides the best fit to RSL data across the Barents Sea when
311 combined with an Earth model that has a relatively thin lithosphere (71 km), an upper mantle
312 viscosity of 0.2×10^{21} Pa s, and a lower mantle viscosity of 2×10^{21} Pa s. However, when
313 seeking to fit RSL data from Svalbard alone they identified an optimum Earth model with
314 96 km lithosphere, 0.3×10^{21} Pa s upper mantle viscosity, and 2×10^{21} Pa s lower mantle
315 viscosity (Auriac et al., 2016; their Table 4). We combined the ICE-6G_C ice history with
316 two lithosphere thickness values, 71 and 96 km, and a range of combinations of upper and
317 lower mantle viscosities (Supplementary Figure 2). A global glacial isostatic adjustment
318 (GIA) model was used to predict gravitationally self-consistent RSL for the ICE-6G_C model
319 at Prins Karl Forland – to compare with existing RSL data from Forman et al. (2004) – and
320 the MeBo site (Supplementary Figure 2 and Table S2). Differences between model
321 predictions are small (<5 m) during the period of interest, and we use the model characterized
322 by 71 km lithosphere thickness, 0.2×10^{21} Pa s upper mantle viscosity, and 5×10^{21} Pa s
323 lower mantle viscosity.

324 We use both the UiT and the ICE-6G_C ice history models to provide RSL boundary
325 conditions when modeling the past behaviour of methane hydrates because they predict
326 contrasting behaviour over the past 11000 years (Figure 3c). The UiT scenario predicts a rise
327 in RSL (increase in water depth) for the upper slope site over 11200–8000 years BP followed
328 by a fall. The ICE-6G_C model predicts steady RSL prior to 8000 years BP and a steep RSL
329 rise between 8000–4000 years BP, followed by a gentler rise during the last 4000 years.

330 **4. Results**

331 We ran the 1-D hydrate model with four different plausible combinations of BWT and RSL
332 changes, viz. BWT_{TF} & UiT RSL, BWT_{TF} & ICE-6G_C RSL, BWT_C & UiT RSL, BWT_C &
333 ICE-6G_C RSL (Figures 3a and 3c) to understand the past behavior of methane hydrates in
334 response to both temperature and pressure changes. At the start of each model, we assume
335 that any chloride anomaly due to the formation of the hydrate that is already present or due to
336 hydrate dissociation before the start of the model run has diffused away. We present results
337 showing the transient behavior of the hydrate layer that dissociates into methane and
338 freshwater. The MeBo drilling experiment detected a freshwater anomaly, but it did not

339 directly recover hydrate, i.e., we do not have direct constraints on hydrate saturation values.
340 We varied the initial hydrate saturations with the aim to identify the critical hydrate saturation
341 for which hydrates completely dissociated, and the final chloride values mimicked the
342 observed values (Wallmann et al., 2018).

343 **4.1. Response to BWT_{TF} variations and UiT RSL change (Scenario 1)**

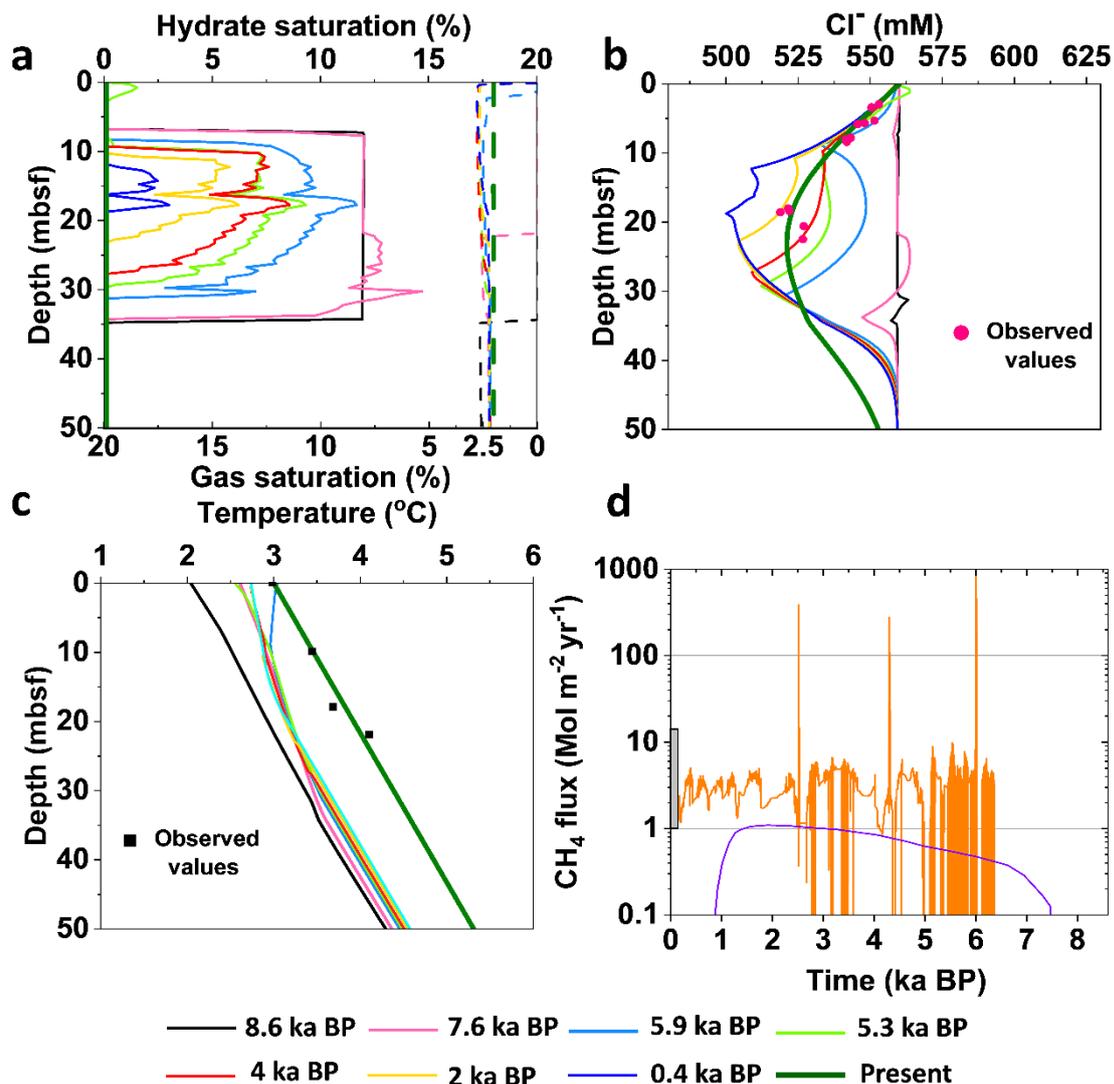
344 The evolution of the hydrate layer in response to the BWT_{TF} time series and a fall in RSL
345 (UiT), which is considered Scenario 1 (Figures 3a and 3c), is illustrated in Figure 4. The
346 model forcing incorporated in this scenario is similar to Wallmann et al. (2018). Table S3
347 shows the differences between our model and the depressurization model proposed by
348 Wallmann et al. (2018).

349 We initialized the model at 8.6 ka when the seabed temperature was at 2.05° C, and high sea
350 level favored a thicker initial hydrate column (Figure 4a and Supplementary Figure 1a)
351 compared to Wallmann et al. (2018). We assumed 12% initial hydrate saturation and a gas
352 layer beneath it with 2–4% saturation (Figure 4a), 560 mM Cl⁻ concentration (Figure 4b)
353 equivalent to pore water salinity of 3.5 wt%, and a steady geothermal gradient of 50° C/km
354 (Figure 4c) in the sediment column. For this scenario, the dissociation of 12% hydrate
355 resulted in the chloride anomaly that best matched the observations (Figures 4a–c).

356 During 8.6–7.6 ka, the seabed temperature increased from 2 to 2.6 °C, and RSL increased by
357 ~2 m (Figures 3a and 3c). During this time, a small amount of hydrate dissociated at the base
358 of the hydrate stability zone (Figure 4a). Subsequently, RSL dropped by ~3 m, and BWT
359 increased from 2.6 to 3 °C over 7.6–5.9 ka. During this time, hydrate dissociation was
360 governed by warming and depressurization (Figure 4a), freshening the pore water and
361 resulting in a chloride anomaly (Figure 4b). The methane released due to hydrate dissociation
362 flowed upward through the stability zone and vented into the ocean. When the BWT reduced
363 to 2.5° C, such as around 5.3 ka, methane was converted into hydrate and chloride
364 concentration increased at shallower depths (Figures 4a and 4b). In the Wallmann et al.
365 (2018) model, shallow hydrate reformation was not considered, and methane vented rapidly.
366 From 5 ka to 0.4 ka, RSL fall continued to erode the hydrate layer. Between 0.4–0.08 ka, the
367 remaining hydrate layer dissociated due to high BWT conditions (>2.5 °C), since the RSL
368 change was negligible, and the geothermal gradient shifted toward the steady geotherm
369 (Figure 4c).

370 According to the UiT model, isostatic rebound resulted in an ~8 m fall in RSL between 6 ka
371 and present at our study site (Wallmann et al., 2018). Similar to Wallmann et al. (2018), the
372 consequent drop in hydrostatic pressure was responsible for hydrate dissociation. Hydrate
373 saturation in our model is significantly less than their model; however, the initial hydrate
374 layer is six times thicker than their assumption. In the Wallmann et al. (2018) model, the
375 hydrate layer completely dissociated 1000 years BP producing a large chloride anomaly that
376 progressively diffused away, and methane flow ceased. However, in our case, hydrate
377 dissociation was not completed until 80 years BP, at which point it resulted in a chloride
378 minimum in the system. Subsequently, the chloride anomaly diffused away to match the
379 observed values (Figure 4b) and methane flow was maintained at a low rate (Figure 4d).

380 We also tested an initial hydrate saturation of 10% for this particular pressure-temperature
381 scenario. This model could not fit the observed chloride saturation profile since the hydrate
382 dissociated earlier (Supplementary Figure 3).



383

384 **Figure 4.** Simulation results for pressure-temperature driven (UiT RSL and BWT_{TF} as model
 385 forcing) methane hydrate dissociation. The results are plotted at a series of specific
 386 times. (a) Evolution of the hydrate layer thickness and saturation (solid line) and gas layer
 387 beneath it (dash line) over the simulation period. Hydrate completely dissociated 80 years
 388 BP. (b) Chloride depletion suggests pore water freshening as hydrates dissociate. The
 389 chloride concentration started to diffuse away after hydrate dissociation was complete. Pink
 390 dots indicate values in cores retrieved at 391 m water depth (Wallmann et al., 2018). (c)
 391 Temperature profile in the sediment column. Black squares indicate temperatures measured
 392 in drill holes at 391 m water depth. (d) Methane emission rate from sediments into the
 393 overlying water column. Modeled fluxes induced by hydrate dissociation over the past 8610
 394 years (orange curve) are compared to the area-averaged range of present-day methane gas
 395 fluxes measured at active seeps (grey vertical bar) in the study area (Sahling et al., 2014;
 396 Wallmann et al., 2018; Ferré et al., 2020) and methane flux (purple curve) from the
 397 depressurization-induced hydrate dissociation model of Wallmann et al. (2018).

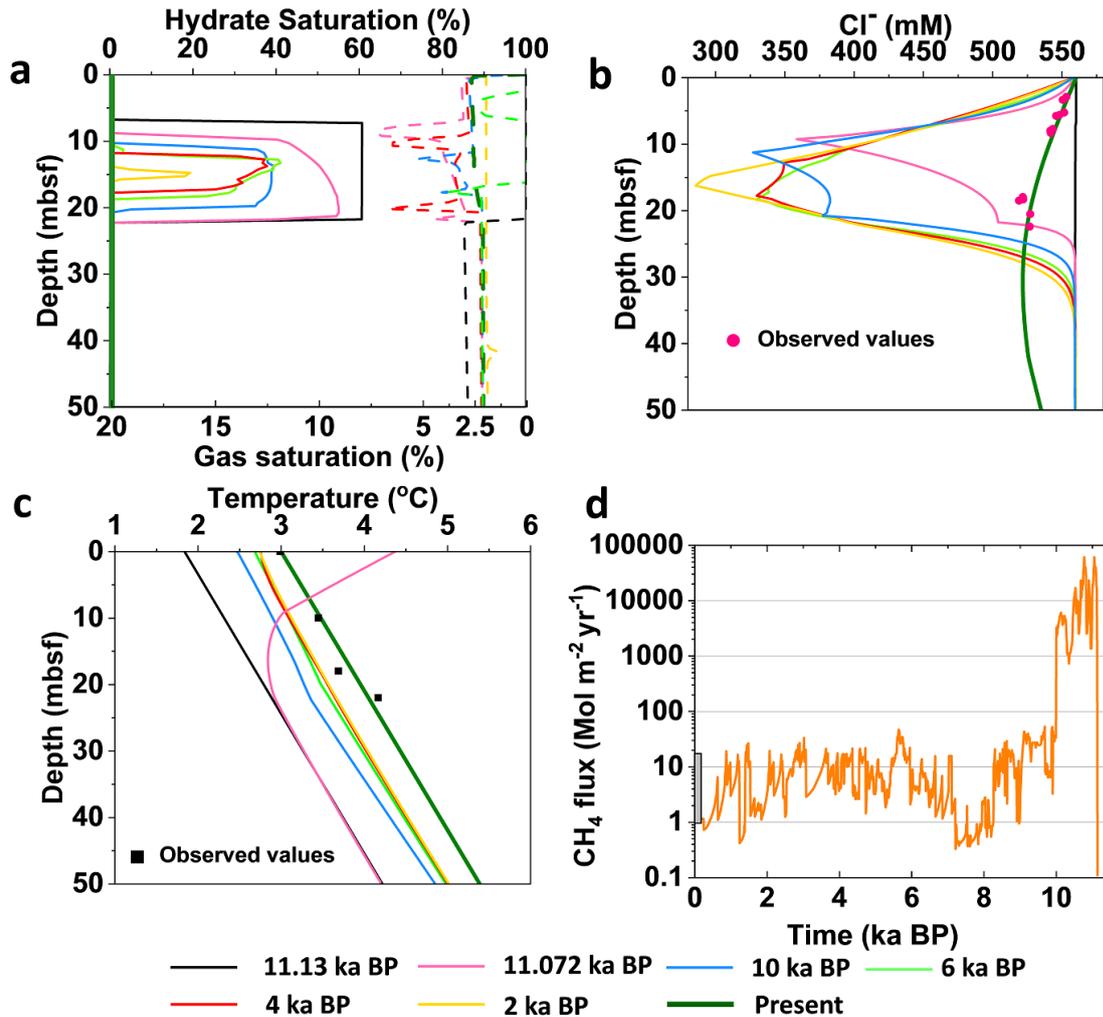
398 4.2. Response to BWT_{TF} variations and ICE-6G_C RSL change (Scenario 2)

399 In an alternate scenario 2, we tested the impact of BWT_{TF} and ICE-6G_C RSL on the hydrate
 400 layer. The ICE-6G_C model predicts RSL rise over the past ~8000 years at the upper
 401 continental slope off west Svalbard. The rise in BWT during this period could destabilize

402 hydrate while the RSL rise will have a stabilizing effect; the combined effect of these
403 contrasting factors has not previously been considered.

404 In this case, we initialized the model at 11.138 ka (rather than 8.6 ka). Initialization at 8.6 ka
405 resulted in a very shallow hydrate layer (0.7–1.2 m) due to the low hydrostatic pressure (~380
406 m water depth), and the dissociation of this hydrate over 8.6–0.04 ka resulted in a chloride
407 anomaly at relatively shallow depths, whereas the observed chloride anomaly is deeper (10–
408 20 mbsf).

409 The colder BWT at 11.138 ka favoured a deeper initial hydrate layer (7–22 mbsf). We used
410 an initial hydrate saturation of 60% because this resulted in the chloride anomaly that best
411 matched the observations (Figure 5a). During the initial 1100 years of the simulation,
412 intermittent bottom water warming led to enhanced hydrate dissociation with substantial
413 methane and freshwater release (Figures 5a–c and Supplementary Figure 4). At ~10 ka,
414 hydrate saturation was reduced to 40%. Over the period 10–8 ka, there was a slight RSL fall,
415 resulting in a hydrostatic pressure decrease, and this was accompanied by bottom water
416 cooling (2.7–2 °C, Figures 3a, and 3c). Dissociation slowed due to the bottom water cooling,
417 but the sediment column remained outside the hydrate stability. Since 4 ka, the rise in RSL is
418 predicted to have been small (Figure 3c). The remaining hydrate layer completely dissociated
419 at 40 years BP due to warm bottom water conditions (>2.5 °C). The resulting chloride
420 anomaly started to diffuse away, and the final chloride values are consistent with the
421 observations (Figure 5b). The methane flow (Figure 5d) is predicted to have been highest
422 during the earliest Holocene (11.2–10 ka), continuing through to the present day at a reduced
423 rate. In this scenario, the BWT rise was primarily responsible for hydrate destabilization
424 despite the RSL rise.



425

426 **Figure 5.** Simulation results for pressure-temperature driven hydrate dissociation forced by
 427 ICE-6G_C RSL and BWT_{TF} over the past 11000 years. (a) Evolution of the hydrate layer
 428 thickness and saturation over the simulation period. Solid lines mark hydrate saturation and
 429 dashed lines mark gas saturation. Hydrate completely dissociated 40 years BP. (b) Dissolved
 430 chloride concentration. (c) Temperature profile in the sediment column. Black squares
 431 indicate temperatures measured in drill holes at 391 m water depth. The steep rise in
 432 temperature towards the seafloor at 11.072 ka reflects the sharp increase in BWT (Figure 3a)
 433 imposed at the top boundary cell. Temperature profiles at selected times are shown for
 434 clarity. (d) Modeled methane emission rates from sediments into the overlying ocean are
 435 compared to the area-averaged range of present-day methane gas fluxes measured at active
 436 seeps (grey vertical bar) in the study area (Sahling et al., 2014; Wallmann et al., 2018; Ferré
 437 et al., 2020). The methane flow rate was highest during the early Holocene.

438 4.3. Response to BWT_C variations with UiT and ICE-6G_C RSL change (Scenarios 3 439 and 4)

440 In two additional scenarios, we applied BWT_C with UiT RSL (Scenario 3) and BWT_C with
 441 ICE-6G_C RSL (Scenario 4) to the hydrate layer. To investigate the development of hydrate
 442 in each of these cases, we initialized the model when the BWT_C attained a minimum value of
 443 1.4 °C at 3.16 ka (Figure 3b) using a hydrate saturation of 5% and we ran the model to
 444 observe changes in the hydrate layer. Simulation results for scenarios 3 and 4 are shown in
 445 Figures 6a-d and 6e-h, respectively.

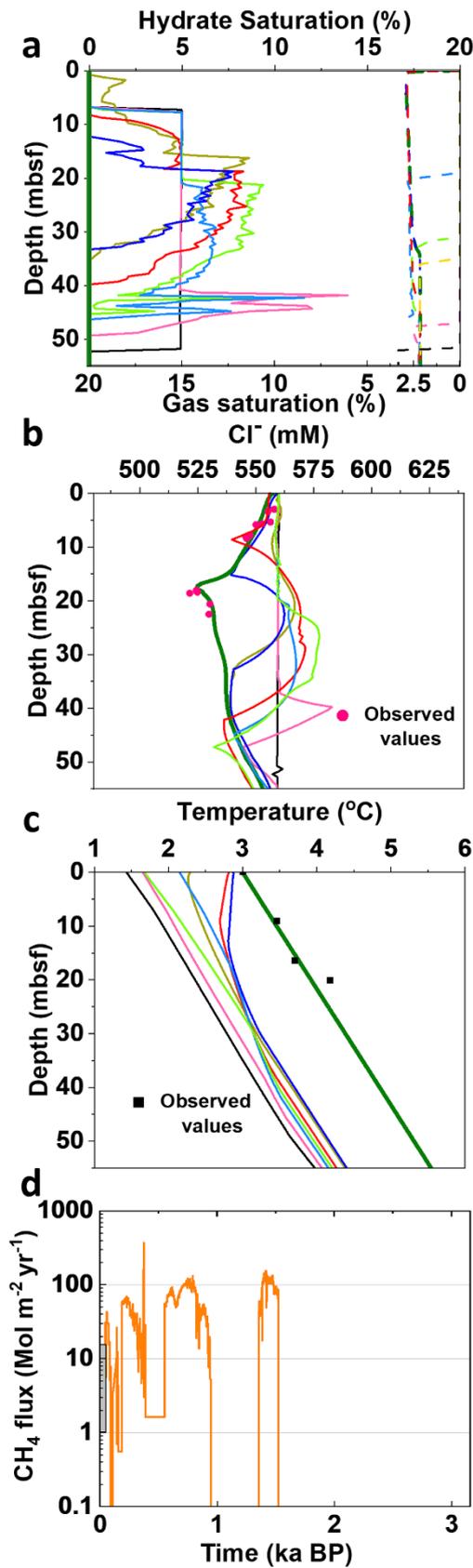
446 In scenario 3 (BWT_C+UiT RSL), hydrate dissociation was triggered by depressurization and
447 BWT rise, while in scenario 4 (BWT_C+ICE-6G_C RSL), dissociation was exclusively
448 induced by bottom water warming. Initially, gas flowed through the hydrate stability zone
449 (Figures 6a and 6e) and was converted into hydrates, increasing hydrate saturation by ~10–
450 15% as well as chloride concentration (Figures 6b and 6f). The BWT_C record documents two
451 prominent warming pulses (BWT>2.5 °C) at ~1.5 ka and ~0.7 ka (Figure 3b). During these
452 warming episodes, the shallowest hydrate dissociated in both models, and gas reached the
453 seafloor. Hydrate completely dissociated at 53 and 49 years BP in scenarios 3 and 4,
454 respectively, and the final chloride anomalies matched the observed values at the end of the
455 simulations (Figures 6b and 6f). The temperature profiles migrated toward the steady
456 geotherm once the hydrate was exhausted (Figures 6c and 6g). Chloride sensitivity was
457 gauged by varying the initial hydrate saturation. The final chloride values for the model with
458 hydrate saturation more than 5% did not fit the observed chloride profile.

459 We also compared the response of the hydrate layer to the BWT reconstruction of Marín-
460 Moreno et al. (2013) (refer to section 3.2). In this case, we initialized the model with 5%
461 hydrate saturation (Supplementary Figures 5a and 5d) at 2.16 ka and forced it using BWT_M
462 (Figure 3b) and both sets of RSL values (UiT and ICE-6G_C) (Figure 3c). In these cases,
463 hydrate dissociated completely at 49 and 46 years BP for BWT_M+ UiT RSL and BWT_M+
464 ICE-6G_C RSL, respectively, and final chloride values matched the present-day observations
465 (Supplementary Figure 5). Although there are differences between BWT_C and BWT_M over
466 the past 2 ka, model outputs at the end of the simulations did not show any major differences.

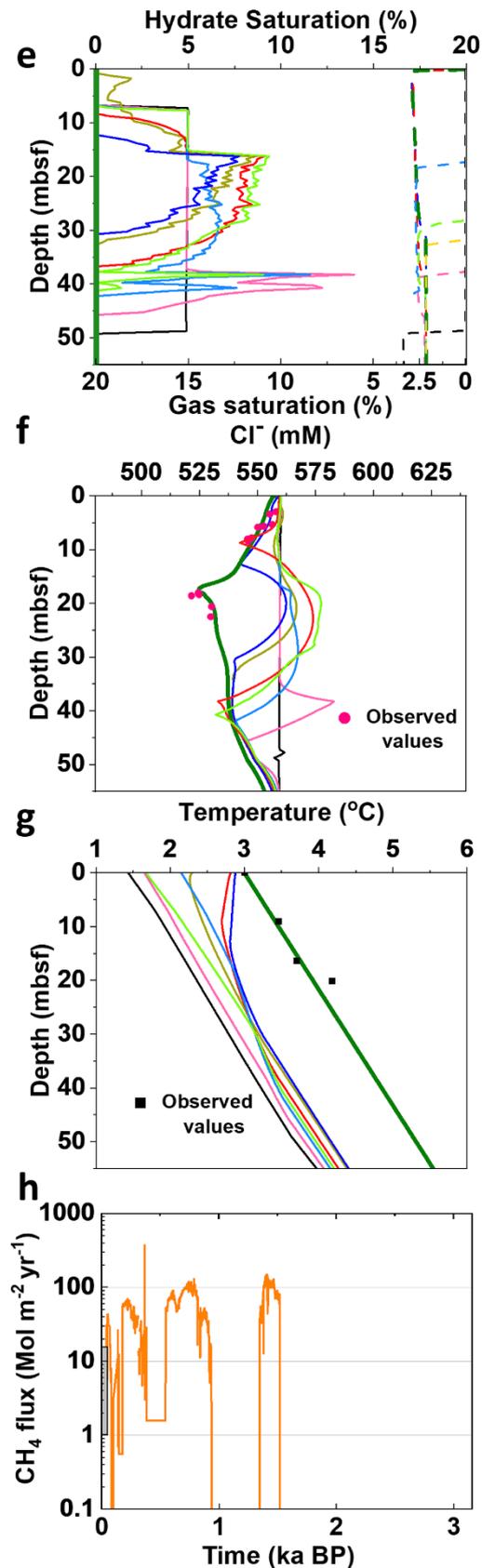
467 The experiments described above were initialized in the late Holocene. To check whether
468 some hydrates could have survived warming episodes earlier in the Holocene, we initialized
469 the model with a high hydrate saturation of 60%, as used in Wallmann et al. (2018), at 11.2
470 ka (Supplementary Figures 6 and 7). In the BWT_C and ICE-6G_C scenario, hydrate
471 dissociation was completed at 10.6 ka (Supplementary Figure 6a) and was driven by seabed
472 warming and minor RSL fall (~2 m). In the BWT_C and UiT scenario, the hydrate layer
473 completely dissociated at 9.6 ka (Supplementary Figure 7a) in response to early Holocene
474 warming despite ~10 m RSL rise. In both cases, the chloride concentration was restored in
475 the next thousand years (Supplementary Figures 6b and 7b). The equivalent early Holocene
476 warming model by Wallmann et al. (2018) also showed complete hydrate dissociation by
477 10.7 ka; however, they did not consider any RSL change.

478

Scenario 3 (a-d)



Scenario 4 (e-h)



— 3.2 ka BP — 2.4 ka BP — 2.1 ka BP — 1.8 ka BP
 — 1.5 ka BP — 1 ka BP — 0.8 ka BP — Present

480 **Figure 6.** Simulation results show methane hydrate dissociation forced by $BWT_C + UiT$ RSL
481 (a–d) and $BWT_C + ICE-6G_C$ RSL (e–h) over the past 3160 years (Figures 3b and 3c). (a, e)
482 Evolution of the hydrate layer thickness and saturation over the simulation period. Solid lines
483 mark hydrate saturation and dashed lines mark gas saturation. Hydrate dissociated at the base,
484 and gas migrated through the hydrate stability zone. Hydrate completely dissociated at 53 and
485 49 years BP in scenarios 3 and 4, respectively. (b, f) Change in dissolved chloride
486 concentration. Hydrate dissociation caused pore water freshening. The chloride concentration
487 started to diffuse after the completion of hydrate dissociation, and the final profiles match the
488 present-day observations (dots). (c, g) The temperature profile in the sediment column. Black
489 squares indicate temperatures measured in drill holes at 391 m water depth. (d, h) Methane
490 emission rate from sediments into the overlying ocean (orange curve) and area-averaged
491 range of present-day methane gas fluxes (grey vertical bar) in the study area (Sahling et al.,
492 2014; Wallmann et al., 2018; Ferré et al., 2020).

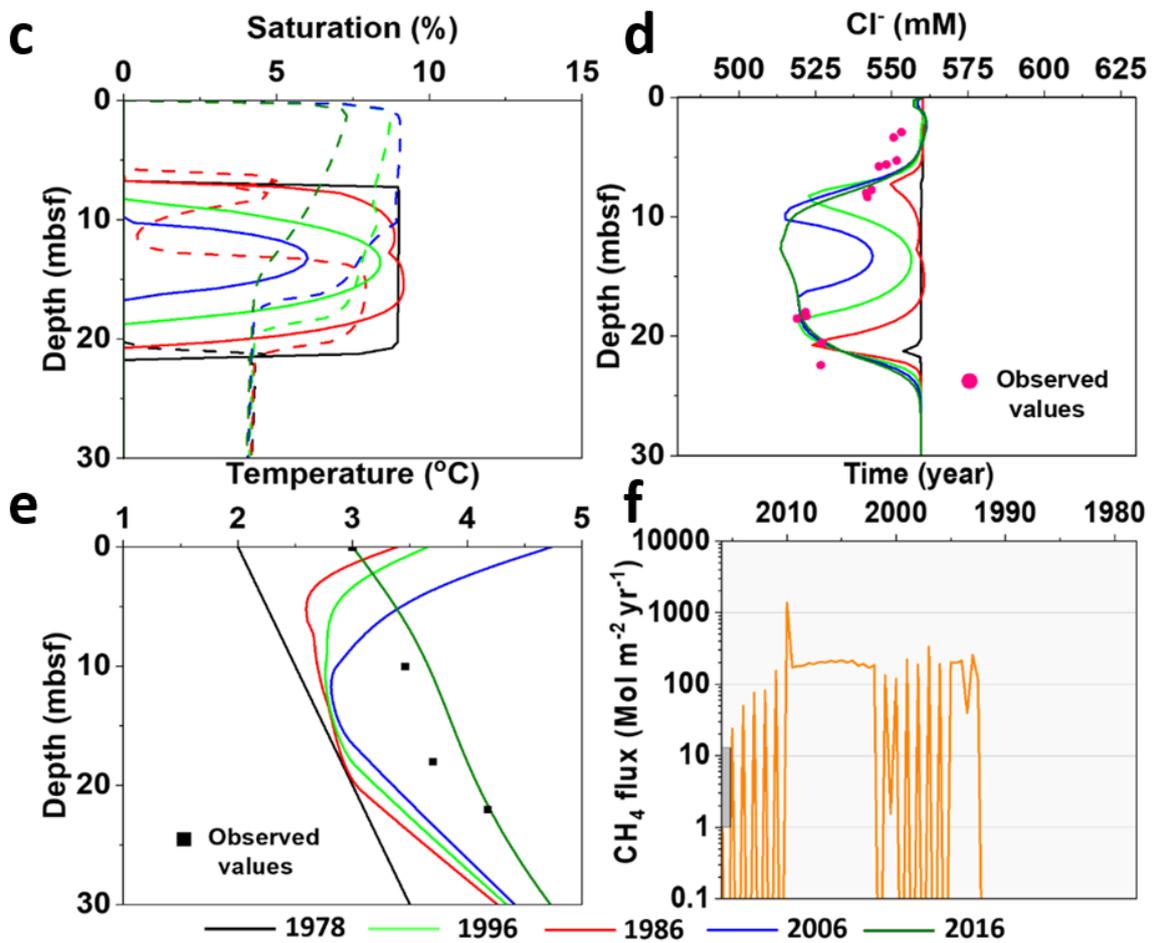
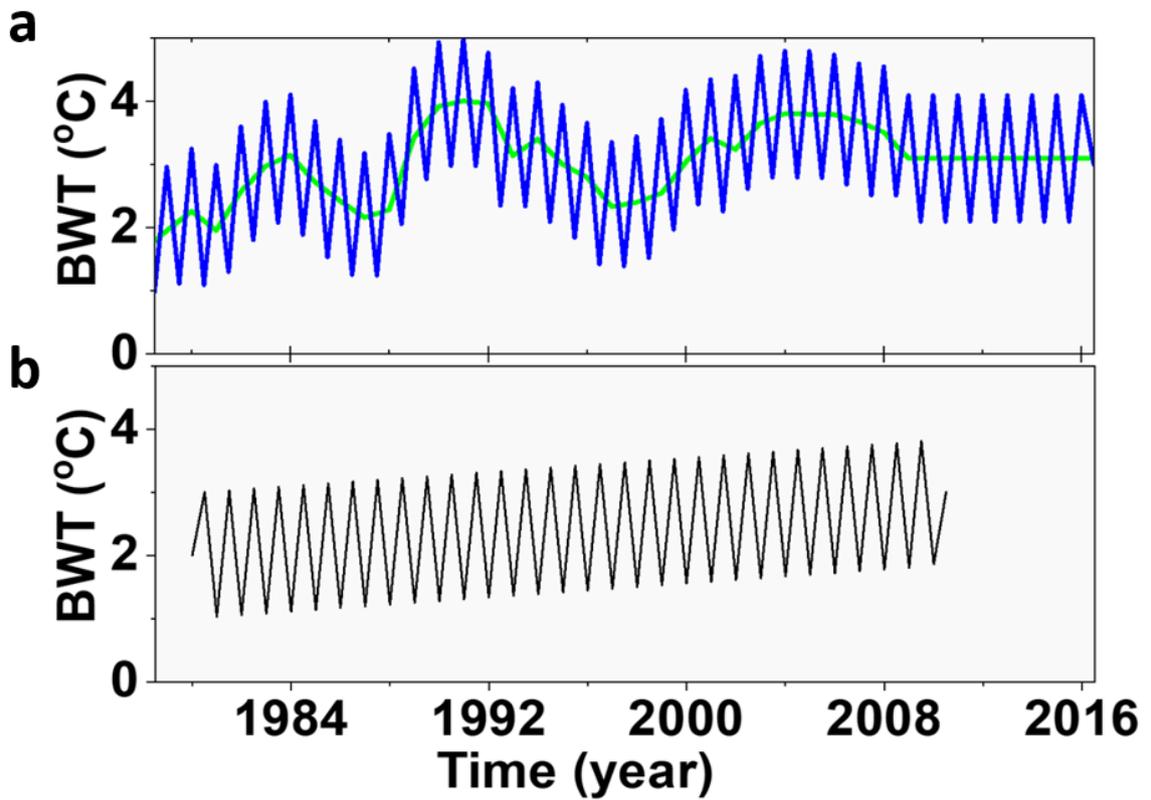
493

494 **4.4. Response to final 38 years of BWT changes**

495 It is evident from the previous long-time series runs that complete hydrate dissociation at this
496 site was possible well before the accelerated bottom water warming during 1978–2008
497 (Westbrook et al., 2009). For example, dissociation of 12% hydrate was complete at 80 years
498 BP in scenario 1 (BWT_{TF} and UiT RSL, Figure 4a), while in scenarios 3 and 4, covering the
499 past ~3200 years (BWT_C and UiT RSL, Figure 6a and BWT_C and $ICE-6G_C$ RSL, Figure
500 6e), 5% hydrate completely dissociated around 50 years BP.

501 In such cases, we examine whether hydrate reformation and subsequent dissociation were
502 possible over the last 38 years in our model runs (1978–2016). Several factors provided
503 suitable conditions for hydrate reformation during this period. During the previous episode of
504 hydrate dissociation, pore water freshening would have enhanced hydrate stability. Similarly,
505 gas from previously dissociated hydrate and deeper sources would have been available to be
506 converted into hydrate during cold BWT conditions. Importantly, any new hydrate formation
507 will have erased any earlier chloride depletion anomaly because chloride is excluded from the
508 hydrate during the reformation.

509



511 **Figure 7.** BWT variations at 400 m water depth off west Svalbard and simulation results
512 showing the behavior of hydrate forced by BWT variations over 1978–2016 in a sediment
513 column with permeability 10^{-13} m^2 . (a) The instrumental record of BWT during 1978–2016
514 (Green curve, Ferré et al. 2012) with seasonal fluctuations of $\pm 1^\circ\text{C}$ (Blue) following Riedel et
515 al. (2018). (b) The 1980–2010 linear temperature rise from $2\text{--}3^\circ\text{C}$ was adopted by Wallmann
516 et al. (2018) with seasonal fluctuations of $\pm 1^\circ\text{C}$. (c) Methane hydrate saturation over the
517 simulation period. The initial hydrate saturation was 9% in 1978. Solid lines mark hydrate
518 saturation and dashed lines mark gas saturation. Hydrate saturation decreased with
519 progressive thermal dissociation. Hydrate completely dissociated 2 years before the end of
520 the model run. (d) Dissolved chloride concentration. The final chloride profile provides a
521 good fit to the observations. (e) The temperature profile in the sediment column. Black
522 squares indicate temperatures measured in drill holes at 391 m water depth. (f) Modeled
523 methane fluxes induced by methane hydrate dissociation compared to the area-averaged
524 range of present-day methane gas fluxes measured at active seep sites (grey vertical bar)
525 (Ferré et al., 2020). In the experiment, methane appeared at the seabed 14 years after the
526 onset of warming.

527 Our 38-year run (1978–2016), driven by the BWT time series shown in Figure 7a, finds that
528 an initial hydrate layer with 9% hydrate saturation completely dissociates and produces a
529 chloride profile that matches the observations (Figures 7c and 7d). Since the RSL change in
530 past 4 decades is negligibly small, we did not include RSL change in this model. The net 1°C
531 BWT (Figures 7a and 7e) rise led to hydrate dissociation, and the resulting methane flowed
532 through the hydrate stability zone (Figure 7c). Methane appeared at the seabed 14 years after
533 the onset of warming when we considered a permeability of 10^{-13} m^2 (Figure 7f). We
534 observed periodicity in methane flow (Figure 7f) that is attributed to seasonally-induced
535 hydrate formation and dissociation at a shallow depth. The seasonal fluctuations in BWT
536 (Figure 7a) strongly affected the top 10 m of the sediment column, where we find repeated
537 episodes of shallow hydrate reformation and dissociation during colder and warmer
538 conditions, respectively.

539 Previously, Wallmann et al. (2018) ran a 30-year model (1980–2010) driven by a linear
540 temperature increase from 2 to 3°C that included seasonal temperature variations $\pm 1^\circ\text{C}$
541 (Figure 7b). The initial hydrate layer in their model was at 10–30 mbsf, and hydrate
542 saturation was varied in the range of 6–8%. Hydrate did not completely dissociate by the end
543 of their model run, and the chloride anomaly could not be reproduced. Hence, they concluded
544 that the thermal dissociation of hydrates over the past three decades was unlikely. Significant
545 time and energy are required to dissociate 6–8 % hydrate. Although measured down-core
546 temperature gradients lie in $45\text{--}50^\circ\text{C}/\text{km}$ (Riedel et al., 2018), their 30-year model assumed a
547 geothermal gradient of $45^\circ\text{C}/\text{km}$.

548 Our model was run for an additional 8 years compared to Wallmann et al. (2018), so there
549 was sufficient time for hydrate dissociation. The heat flux ($10 \times 10^{-2} \text{ W m}^{-2}$) applied at the base
550 of our model resulted in a geothermal gradient of $50^\circ\text{C}/\text{km}$ in the sediment column, which,
551 in turn, resulted in a thinner initial hydrate layer (14 m thick, 7–21 mbsf) than the 20 m
552 thickness considered by Wallmann et al. (2018). Our thinner hydrate layer completely
553 dissociated over 38 years. However, when we considered a lower geothermal gradient
554 ($45^\circ\text{C}/\text{km}$) by applying a heat flux of $9 \times 10^{-2} \text{ W m}^{-2}$, the initial hydrate layer was thicker
555 (22 m thick, 7–29 mbsf). It did not completely dissociate after 38 years.

556 We also tested lower (10^{-15} m^2) and higher permeability (10^{-11} m^2) scenarios (Supplementary
557 Figure 8) and different irreducible gas saturation (S_{irg}) values ranging between 2–10% (Ma et
558 al., 2020). We found that hydrate dissociated completely in all cases, and the methane flow

559 rate was very sensitive to the permeability and $S_{i_{rg}}$ values (Thatcher et al., 2013). As
560 expected, a reduction in permeability caused a delay in the arrival of the methane at the
561 seabed (Supplementary Figure 9a), while an increase in the $S_{i_{rg}}$ value caused a delay in the
562 arrival of the methane at the seabed and a decline in the maximum methane flux
563 (Supplementary Figure 9b).

564 The simulations described in this section emphasize the role of recent BWT rise as a potential
565 cause for hydrate dissociation and pore water freshening. They also illustrate the role of
566 interannual temperature changes in controlling changes in seafloor methane flux in response
567 to near seafloor hydrate dissociation-formation stages.

568 **5. Discussion**

569 We have examined the impact of RSL change coupled with BWT variations on hydrate west
570 of Svalbard at 391 m water depth, where recent drilling confirmed pore water freshening due
571 to hydrate dissociation. Table 1 summarizes model outputs from scenarios where the final
572 chloride anomaly and subsurface temperature profile closely match the present-day
573 observations.

574 The previous modelling by various authors either emphasized the role of temperature rise or
575 RSL fall in explaining hydrate destabilization. For example, Thatcher et al. (2013), Marín-
576 Moreno et al. (2013), and Marín-Moreno et al. (2015) considered the effects of BWT on
577 hydrates, while Wallmann et al. (2018) attributed hydrate dissociation to RSL fall at constant
578 BWT. Extending these studies, we have examined the fate of hydrates in response to the
579 combined effects of a plausible set of BWT and RSL values.

580 We have demonstrated plausible scenarios of hydrate dissociation based on RSL
581 reconstruction and BWT changes that show prominent millennial to seasonal changes over
582 the past 8000 years at different time scales (decadal to seasonal). All of the scenarios
583 considered here predict thinning of the hydrate layer from the base during the Holocene,
584 releasing methane that can get trapped in shallow sediments under cold BWT conditions
585 ($<2.5^{\circ}\text{C}$). Sharp rises in BWT dissociate the shallow hydrates and enhance gas flow at the
586 seabed. Periods of high methane flow at the seabed could lead to seep-carbonate
587 accumulation at the feather edge of hydrate stability. The precipitated carbonate can influence
588 benthic ecosystems since it provides hardground shelter for them (Wilson et al., 2007).
589 Thicker carbonate accumulations of varying ages could be related to multiple episodes of
590 hydrate dissociation (Berndt et al., 2014). Our modelling shows a range of possible past
591 methane emission scenarios. Future work should focus on mineralogical and isotopic studies
592 of authigenic carbonates (e.g., Liang et al., 2017) to unravel fluid sources, long-term seepage
593 dynamics, and the fate of Arctic methane hydrates.

594 Our results based on coupled BWT_{TF} and UiT RSL values (Scenario 1) challenge the
595 hypothesis that a fall in RSL was the sole cause of hydrate dissociation over 8–1 ka
596 (Wallmann et al., 2018). The depressurization model proposed by Wallmann et al. (2018)
597 only considered a constant BWT over 8–0.1 ka. By using the UiT RSL in combination with
598 the more realistic BWT_{TF} curve, we find that pulses of warm bottom water temperatures
599 ($>2.5^{\circ}\text{C}$) also contributed to hydrate dissociation together with depressurization. The initial
600 hydrate saturation was significantly higher (60%) in the Wallmann et al. (2018) model
601 compared to the initial saturation (12%) in our BWT_{TF} and UiT RSL scenario. Previously
602 inferred hydrate saturations in the top 50 m of glaciomarine sediments on the upper
603 continental slope of Svalbard, in water depths of 480–866 m, range between 5–22% (Chabert

604 et al., 2011 and Goswami et al., 2016). While the initial hydrate saturations considered here
605 are within this range in most cases, we infer a higher value (60%) in the model forced by
606 BWT_{TF} and ICE-6G_C RSL (Scenario 2). We needed such a high initial hydrate saturation to
607 generate the chloride anomaly that mimics present-day observations. In scenario 2, hydrates
608 of such high saturation slowly dissociated over a prolonged period (11138 years). The model
609 shows a prolonged record of methane emission at the seabed, with methane flux
610 progressively decreasing with time (Figure 5d).

611 Based on their model (UiT RSL), Wallmann et al. (2018) concluded that methane fluxes from
612 hydrate dissociation ceased at 1 ka upon completion of depressurization-induced dissociation,
613 and the present-day gas seeps are not related to warming-induced hydrate dissociation. This
614 inference disagrees with the findings of Veloso-Alarcon et al. (2019) and Ferré et al. (2020),
615 who proposed that the present-day methane flux is likely controlled by temperature-driven
616 seasonal hydrate dynamics. In our simulation using the BWT_{TF} and UiT RSL, complete
617 hydrate dissociation takes longer than in Wallmann et al. (2018) mainly because some time is
618 required between the onset of hydrate dissociation and methane reaching the seabed and it
619 depends on effective permeability, porosity, irreducible gas saturation, hydrate saturation,
620 distance between the top of the hydrate layer and the seafloor (Thatcher et al., 2013). For this
621 scenario, the modeled methane flux at present (year 2016) lies at the lower end of the range
622 of observed values ($1 \text{ mol m}^{-2} \text{ year}^{-1}$, compared with $1\text{--}13 \text{ mol m}^{-2} \text{ year}^{-1}$, Ferré et al., 2020),
623 suggesting that other possible explanations for the methane seeps should be investigated.

624 Despite the possibility of several episodes of hydrate formation and dissociation in the past,
625 the latest episode of hydrate formation and subsequent warming (1978–2016) stands out as a
626 strong possibility that can explain the observed pore water freshening and variations in
627 methane flux at the seabed. The observed gas flare activity shows large variations that are
628 strongly influenced by seasonal hydrate dynamics (Berndt et al., 2014; Veloso-Alarcon et al.,
629 2019; Ferré et al., 2020). This pattern is well explained by our 38-year simulation (1978–
630 2016), where the initial 9% hydrate saturation completely dissociated, and hydrate dynamics
631 are clearly influenced by decadal- and seasonal-scale BWT fluctuations (Figure 7f). There is
632 seismic evidence of gas (Chabert et al., 2011; Sarkar et al., 2012) in the top 100 m of the
633 sediment column at the upper continental slope (water depths 350–1200 m). In all the cases
634 we considered (Table 1), our modeling predicts that hydrates have completely dissociated
635 prior to 1978. However, the combination of the delayed upwards migration of the resulting
636 methane through vertical fractures and permeable strata (Sarkar et al., 2012) and colder
637 BWTs, such as the remarkable cooling (10-year) during the late 1970s (Westbrook et al.,
638 2009), likely created favourable conditions for hydrate re-formation. Since no hydrates were
639 recovered by the drilling, ocean bottom warming must have completely dissociated these
640 newly-formed hydrates to explain the observed pore-water freshening. Therefore, there is
641 insufficient evidence to reject the (*recent*) warming-induced hydrate dissociation hypothesis
642 (Westbrook et al., 2009) in explaining the present-day methane seeps off west Svalbard.

643 Our 38-year model provides valuable insights into the response of methane hydrates to recent
644 warming at the feather edge of hydrate stability in this area. Moreover, it evidences the
645 importance of reducing uncertainty in the predictions of methane emissions, particularly on
646 the continental shelf. Since the shelf regions are at shallower water depths (100–200 m),
647 seafloor methane emissions have a greater chance to reach the atmosphere. Even if methane
648 does not reach the atmosphere in these areas, methane dissolution and its oxidation in the
649 ocean can cause local changes in pH, thereby affecting the local marine biodiversity

650 (Valentine et al., 2001; Riebesell, 2008). Although gas emission activity could substantially
651 subside at the feather edge of hydrate stability during bottom water cooling on the upper
652 continental slope (~390 m), deeper gas can still be deflected by the impermeable hydrate
653 layer and migrate upslope towards the shelf through permeable beds (Sarkar et al., 2012;
654 Veloso-Alarcon et al., 2019). The beds outcrop at the shelf, and methane is released directly
655 into the ocean. Veloso-Alarcon et al. (2019) reported enhanced methane flux on the
656 continental shelf (100–200 m water depths) compared to the upper slope during cooler BWT
657 conditions since the shelf is outside the gas hydrate stability. Strong methane emission has a
658 greater chance of reaching the sea surface and atmosphere, contributing to further warming
659 (Westbrook et al., 2009).

Model forcing and duration of simulation (2016 represents ‘Present’)	Time taken to completely dissociate hydrate (years)	Initial hydrate saturation, initial hydrate layer thickness	Correlation between observed chloride anomaly and final chloride values obtained from the model	Methane flux at the end of the simulation (Observed area-averaged methane flux is 1–13 Mol m ⁻² y ⁻¹)
Scenario 1: BWT _{TF} + UiT RSL (8.6–0 ka)	8520	12%, 30 m	0.9	1
Scenario 2: BWT _{TF} + ICE-6G_C RSL (11.138–0 ka)	11098	60%, 15 m	0.8	1
Scenario 3: BWT _C + UiT RSL (3.16–0 ka)	3107	5%, 45 m	0.9	2.25
Scenario 4: BWT _C + ICE-6G_C RSL (3.16–0 ka)	3111	5%, 42 m	0.9	2.2
Scenario 5: BWT from (1978–2016) (Ferré et al., 2012)	36	9%, 14 m	0.94	7

Table 1. Summary table showing the modeling scenarios and quantification of how well the final results (chloride values, and methane flux) matched the observations.

660 **6. Conclusions**

661 We have re-examined the fate of Arctic marine hydrates in response to BWT and RSL
662 changes at the landward limit of hydrate stability off west Svalbard. We have considered an
663 alternative ice history model (ICE-6G_C) that suggests RSL rise and contradicts the RSL fall
664 predicted by the UiT model over the past 8000 years used in a previous modelling study in
665 the area (Wallmann et al., 2018). We demonstrate that with the available constraints, multiple
666 episodes of hydrate formation and dissociation are still possible based on a comprehensive
667 scenario testing approach that has never been attempted before for such a long-time scale.

668 In the simulation that used the UiT scenario, we found that although RSL fall may have
669 triggered hydrate dissociation, decadal-scale BWT fluctuations (BWT_{TF} and BWT_C) played a
670 critical role in controlling hydrate dynamics. Specifically, in these scenarios the eventual fate
671 of methane released from depressurization-induced hydrate dissociation depended on BWT.
672 Warmer seabed temperatures (e.g., $>2.5^\circ\text{C}$) prevented the shallow hydrate reformation,
673 allowing gas to flow through the hydrate stability zone and into the ocean. On the other hand,
674 colder bottom water temperatures facilitated shallow hydrate reformation. In the simulations
675 that used the ICE-6G_C scenario, hydrate dissociation was primarily caused by warming of
676 bottom waters, which overrode the stabilising effect of RSL rise.

677 We have not been able to determine the actual RSL and BWT scenario through which the
678 hydrate system evolved since several scenarios produce outputs that match the present-day
679 observations. However, we demonstrate that earlier hydrates completely dissociated before
680 the late 1970s' bottom water cooling episode, irrespective of the assumed scenario. During
681 this cooling episode, hydrate could have reformed and then dissociated in response to BWT
682 rise over the period 1978–2016. The gas produced by dissociation takes time to migrate to the
683 seabed and can explain the modern methane seeps. Seasonal warming and cooling cycles
684 impact hydrate dynamics and can explain the large observed variations in methane flux. The
685 pore water freshening observed in the marine sediments is indicative of the latest episode of
686 hydrate destabilization caused by the latest episode of seabed warming.

687 If bottom water warming continues at an accelerated rate, we can expect downslope
688 migration of the landward limit of hydrate stability. The results presented here bring into
689 focus the necessity to closely monitor the effect of warming Atlantic waters on hydrates
690 because continued warming may degrade more Arctic marine hydrates in the impending
691 future at current or enhanced rates.

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699 **Open research**

700 The bathymetric image shown in Figure 1 was derived from the Norwegian Hydrographic
701 Service (NHS) data and the bathymetric data acquired during the JR211 cruise. The JR211
702 bathymetric data can be obtained from www.bodc.ac.uk by contacting enquiries@bodc.ac.uk.
703 The NHS bathymetric data can be accessed from
704 <https://dybdedata.kartverket.no/DybdedataInnsyn/>. The hydrate model results were generated

705 using Tough+Hydrate v1.5, which can be obtained from <https://tough.lbl.gov/licensing->
706 [download/tough-licensing-download/](https://tough.lbl.gov/licensing-download/).

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